

# Gravitational microlensing and vacuum fluctuations

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Probing the Mystery:  
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## Setting up the question

Raychaudhuri equation:

$$\frac{d\theta}{d\lambda} = \ell^a \nabla_a \theta = -\frac{1}{2}\theta^2 - \sigma_{ab}\sigma^{ab} + \omega_{ab}\omega^{ab} - 8\pi G T_{ab}\ell^a\ell^b$$

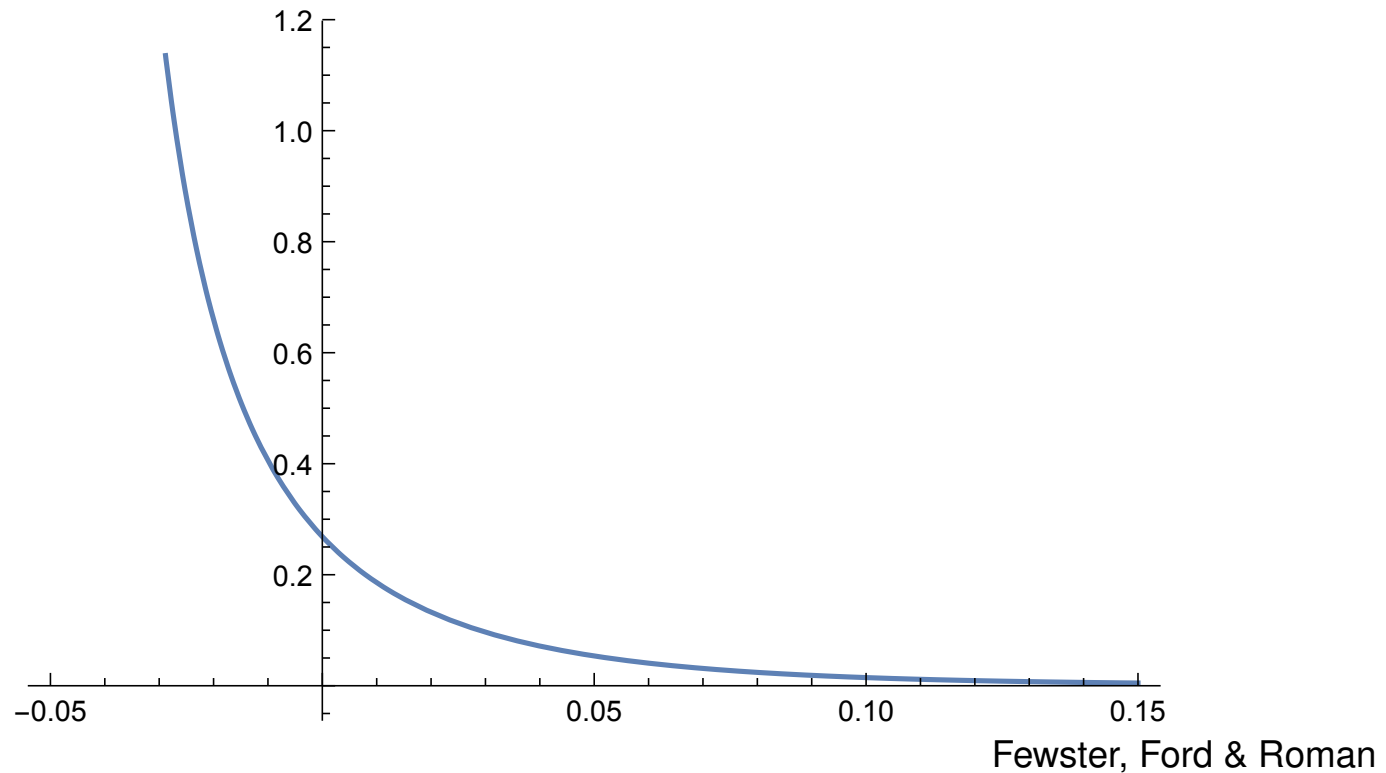
Vacuum fluctuations of the stress-energy tensor  
induce fluctuations of expansion  $\Rightarrow$  (micro)lensing

**Can we observe this effect?**

**First guess:** obviously not

Gravitational fluctuations occur at the Planck scale – much too small

**But:** distribution of fluctuations is *very* non-Gaussian



with a strict lower bound and a long subexponential tail

⇒ naive intuition is unreliable

More precisely:

Smear stress-energy tensor over length scale  $L$  [more later...]

$$\text{Let } x = \frac{16\pi^2}{\hbar} L^4 T_{00} = 16\pi^2 \frac{L^4}{\ell_p^2} G T_{00}$$

“typical” quantum fluctuation has  $x \sim 1$

Fewster, Ford, Roman:

Compute  $\langle T T T \dots \rangle$

Are these the moments of any probability distribution?

Two dimensions: yes, unique, exactly calculable

Four dimensions: yes, not quite unique, approximately calculable

For electromagnetism:

Strict lower bound  $x_0 \approx -.0472$

$$\text{Prob}(x) \sim c_0 (x - x_0)^{-2} e^{-a(x-x_0)^{1/3}}$$

with  $c_0 \approx .955$ ,  $a \approx .963$

## Vacuum fluctuations and gambler's ruin

$$\text{Raychaudhuri: } \frac{d\theta}{d\lambda} = -\frac{1}{2}\theta^2 - 8\pi G T_{ab}\ell^a\ell^b + \dots$$

Vacuum fluctuations of  $T_{ab}\ell^a\ell^b$  are usually negative (defocusing)

But lower bound, long positive tail (focusing)

### **“Gambler's ruin”**

Whatever the odds, if you bet long enough against a house with unlimited resources, you always lose in the end.

## Back-of-the-envelope estimate:

$$\text{Let } \bar{\theta} = \sqrt{\frac{|x_0|}{\pi}} \frac{\ell_p}{L^2}$$

- If  $\theta > \bar{\theta}$ , RHS of Raychaudhuri eqn  $< 0$   
 $\Rightarrow \theta$  is forced down to  $\bar{\theta}$
- If  $\theta < -\bar{\theta}$ , RHS of Raychaudhuri eqn  $< 0$   
 $\Rightarrow \theta$  is forced down to  $-\infty$   
 $\Rightarrow$  focusing, microlensing

To force runaway focusing, need

$$\Delta\theta = \frac{d\theta}{d\lambda} L < -2\bar{\theta}$$

$$\Rightarrow x > 4\sqrt{\pi|x_0|} \frac{L}{\ell_p} = x_1$$

$$\text{Prob}(x > x_1) \sim (L/\ell_p)^{-4/3} e^{-\alpha(L/\ell_p)^{1/3}} \quad (\alpha = \mathcal{O}(1))$$

Significant probability of observation for  $L$  as large as  $5 \times 10^5 \ell_p$

## Time scale for focusing

Solve Raychaudhuri equation after  $\theta < -\bar{\theta}$

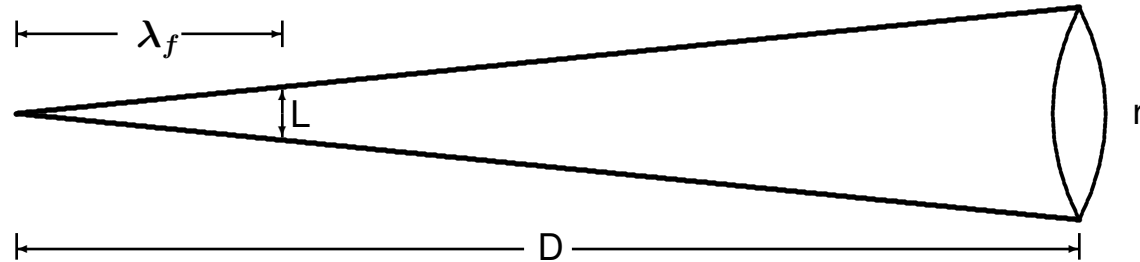
Find time to focusing

$$\lambda_f \sim \beta \frac{L^2}{\ell_p} \quad (\beta = \mathcal{O}(1))$$

Two scales:

- $L$  sets scale of fluctuations
- $\lambda_f$  sets scale at which these have potential observable effect

## Some numbers



For significant magnification, need

$$\frac{r}{D} \sim \frac{L}{\lambda_f} \sim \frac{\ell_p}{L} \Rightarrow L \sim \frac{D}{r} \ell_p, \lambda_f \sim \left(\frac{D}{r}\right)^2 \ell_p$$

For large  $L$  to be relevant, need  $D/r$  large — small, distant sources



## Unfortunately ...

For observability, need  $L/\ell_p \lesssim 5 \times 10^5$

$$\Rightarrow \lambda_f \lesssim 10^{-21} \text{ cm}$$

$$E \sim 10^{10} \text{ TeV}$$

## Conversely ...

Largest realistic  $L$ —white dwarfs in Large Magellanic Cloud

$$L \sim 10^{-21} \text{ m}$$

$$\Rightarrow \text{Prob}(x) \sim 10^{-20000}$$

$$t_{fluc} \sim 10^{-29} \text{ s}$$

## What about multiple fluctuations?

Say  $N$  fluctuations of size  $\delta(GT) \gtrsim \pm \ell_p^2 / L^4$  along path of length  $D$

$$\text{random walk} \Rightarrow \langle \theta \rangle \sim \sqrt{N} \ell_p^2 / L^3$$

$$\Rightarrow \text{nonlinear regime when } L^3 \sim D \ell_p^2$$

For cosmological source: significant for  $L \sim 10^{-15}$  m: 1 GeV photon

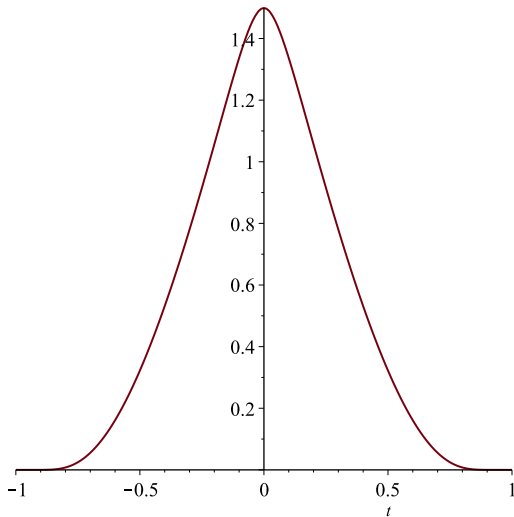
Maybe better: not really random walk (work in progress)

Is any further amplification possible? Maybe...

## Dependence on smearing

Results so far based on Lorentzian smearing of  $T$   
(Fewster, Ford, and Roman, 2012)

New results: probability distribution depends on smearing function;  
“fatter” tail for function with sharp cut-off  
(Fewster and Ford, 2015)



$$\text{Prob}(x > x_1) \sim (L/\ell_p)^{-2/3} e^{-\alpha(L/\ell_p)^{1/6}}$$

Fewster & Ford

Now plausible probabilities for  $L \sim 10^{-23}$  m

Need to better understand role of smearing