#### Displacing entanglement back and forth between the micro and macro domains

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Why do we not observe macro quantum systems?

Existence of a quantum/classical border

Decoherence

Measurement precision

Lacking suited entanglement witness



What is a macroscopic quantum state?

The number of particles is not a sufficient criteria



I. Usmani, C. Clausen, F. Bussieres, N. Sangouard, M. Afzelius, and N. Gisin, Nature Photonics. 6, 234 (2012)

#### What is a macroscopic quantum state?

The number of particles is not a sufficient criteria

Example: entanglement involving 100 000 photons



T.S. Iskhakov, I.N. Agafonov, M.V. Chekhova, and G. Leuchs, PRL 109, 150502 (2012)

#### The Schroedinger cat provides an example



From  $|1\rangle_{\rm atom}|0\rangle_{\rm photon} + |0\rangle_{\rm atom}|1\rangle_{\rm photon}$ 

to  $|1\rangle_{\rm atom} |{\rm Alive}\rangle_{\rm cat} + |0\rangle_{\rm atom} |{\rm Dead}\rangle_{\rm cat}$ 



#### Creating Entanglement at a beamsplitter

Consider a separable state

 $\rho = \sum_{i} p_{i} \ \rho_{a}^{i} \otimes \rho_{b}^{i} = \sum_{k} \bar{p}_{k} |\psi_{a}^{k}\rangle \langle\psi_{a}^{k}| \otimes |\psi_{b}^{k}\rangle \langle\psi_{b}^{k}|$ for each component  $U_{\rm bs}^{-1} |\psi_{a}^{k}\rangle |\psi_{b}^{k}\rangle = |\Psi_{i}\rangle |0\rangle$ 

Since the only pure state leading to a product state after a beamsplitter is the coherent state

All non-classical states lead to entanglement  $\int d^2 \alpha \ P(\alpha) |\alpha\rangle \langle \alpha | \otimes | 0 \rangle \langle 0 |, \ P(\alpha) \ge 0$ 



dip, reported in FIG. 2, has a visibility limited by tains, leaving aside the reflectivity of the beamsplitter and the photon statis tate that describes the only. As such, after the displacement, the detection of sharing a single photon idler photon heralds the generation of an entangled st of the form path entangled state, glement, can be seen as  $\frac{1}{\sqrt{2}} \left( \mathcal{D}_a(\alpha) |1\rangle_A |0\rangle_B + |\alpha\rangle_A |1\rangle_B \right).$ d feature of the heralded  $|\alpha\rangle_A$  results from  $\sqrt{3}$  he displacement of the vacuum. X(2) follows a Poissonian photon number distribution v plifying the mode A by mean photon number  $|\alpha|^2$  equal to the variance.  $\lambda(\alpha)$  corresponding to  $\alpha$ displacement also increases the mean photon num of the single photon state 1, but it preserves e. The latter is obtained in intense local oscillator non-gaussian character. Specifically the state  $\mathcal{D}_a(\alpha)$ tter 18, 19. The physics is characterised by a photon number distribution v d on an interference proa mean photon number  $|\alpha|^2 + 1$  and a variance 3 ocal oscillator need to be The state (1) thus describes entanglement betwee ed in practice by producmicroscopically populated mode B and a mode A wh of a difference frequency mean population can be adjusted by tuning the inten itical nonlinear crystal to of the local oscillator. Remarkably, for large  $|\alpha|^2$ involves a superposition of two components  $\alpha \rangle_A$ creation but stimulated

# Definitions of a macroscopic quantum state $\Phi_0 + \Phi_1$ involving N particles 1\_ Sensitive to decoherence mechanisms

W. Dur, C. Simon, and J.I. Cirac, Phys. Rev. Lett. 89, 210402 (2002)

2\_ Local distinguishability between  $\Phi_0$  and  $\Phi_1$ 

J.I. Korbakken et al., Phys. Rev. A 75, 042106 (2007)

3\_ Large number of one particle operators to go from  $\Phi_0$  to  $\Phi_1$ 

F. Marquardt et al., Phys. Rev. A 78, 012109 (2008)

4\_  $\Phi_0 + \Phi_1$  significant advantage for interferometric applications over  $\Phi_0$  and  $\Phi_1$ G. Bjork and P. Mana, J. of Opt. B 6, 429 (2004) NaN

NaN

#### Displaced single-photon entanglement Micro-macro entanglement?

\_ large number of photons



# Displaced single-photon entanglement

#### Micro-macro entanglement ?

\_ large number of photons

\_ local unitary



# Displaced single-photon entanglement

#### Micro-macro entanglement ?

- \_ large number of photons
- \_ local unitary



\_ 1ebit

## Displaced single-photon entanglement

#### Micro-macro entanglement ?



\_ local unitary

\_ 1ebit

\_ entanglement decreases fast under dephasing process

 $\mathcal{N} = (1 - 3|\alpha|^2 \delta \phi)/2$ 



#### Displaced single-photon entanglement Micro-macro entanglement?



#### Displaced single-photon entanglement Micro-macro entanglement ?





Proposal of a macro measure

based on the distinguishability with a «classical» detector

No need for a microscopic resolution to distinguish  $|{\rm Alive}\rangle_{\rm cat}$  and  $|{\rm Dead}\rangle_{\rm cat}$ 



Pointer state  $\rho_{\hat{x}} = \int dx \ G(x) \ |x\rangle \langle x|$ G(x) Gaussian with spread  $\sigma$ 



P. Sekatski, N. Sangouard, and N. Gisin, in preparation



#### Proposal of a macro measure

based on the distinguishability with a «classical» detector



Macroscopicity as a function of Pguess

(1)  $|0\rangle, |\alpha\rangle$ (2)  $\mathcal{D}(\alpha)|+\rangle, \ \mathcal{D}(\alpha)|-\rangle$ 

P. Sekatski, N. Sangouard, and N. Gisin, in preparation

#### Proposal of a macro measure

based on the distinguishability with a «classical» detector



Any phase fluctuation can be seen as a weak measurement of the photon number

States that are macro with respect our criteria are inevitably very sensitive to phase decoherence

P. Sekatski, N. Sangouard, and N. Gisin, in preparation

## Detecting displaced single-photon entanglement

Detecting macro entanglement is a difficult task

- \_ decoherence inevitably increases the Hilbert space dimension
- \_ requires high resolution detection













The decoherence problem and the requirement on the precision of the measurement are two facets of the same problem

## Conclusion

Proposal for a macro measure based on the distinguishability of superposition components with a «classical» measurement

Displaced single-photon entanglement as an example

Use of a well established entanglement measure

measurements precision <----> decoherence

Useful for opto-mechanics? phase estimation in interferometric measurement?

