Local moment approach to the multi - orbital single impurity Anderson and Hubbard models

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PIPT/Les Houches Summer School on Quantum Magnetism Les Houches, 6th - 23rd June 2006

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Outline

- 1. Kondo effect
- 2. Single impurity Anderson model
 - single orbital Anderson model
 - mean field
 - local moment approach (LMA)
 - multiorbital Anderson model and generalization of LMA

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- 3. Hubbard model and DMFT
- 4. LMA + DMFT in multiorbital Hubbard model

Kondo effect

 Example: resistivity minimum in metals with rare earth impurities





 Universal – in rare earth impurities, quantum dots, carbon nanotubes

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Kondo effect

- below the temperature characteristic for a given material Kondo temperature T_K – (crossover)
- spin fluctuations
 - spin-flip processes



screening of local moments



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Single impurity Anderson model (SIAM)

$$H_{\rm SIAM} = \sum_{\mathbf{k},\sigma} \epsilon_{\mathbf{k}} c^{\dagger}_{\mathbf{k}\sigma} c_{\mathbf{k}\sigma} + \sum_{\mathbf{k},\sigma} V_{\mathbf{k}} \left(d^{\dagger}_{\sigma} c_{\mathbf{k}\sigma} + c^{\dagger}_{\mathbf{k}\sigma} d_{\sigma} \right) + \sum_{\sigma} \left(\epsilon_{d} - U n_{d\bar{\sigma}} \right) n_{d\sigma}$$

Hybridization function

$$\Delta(\omega) = \sum_{\mathbf{k}} \frac{|V_{\mathbf{k}}|^2}{\omega - \epsilon_{\mathbf{k}}} \approx \Delta$$

SIAM model describes

- charge fluctuations
- spin fluctuations
- screening of local magnetic moments







Mean field (Hartree-Fock) approximation

$$U n_{d\bar{\sigma}} n_{d\sigma} pprox U < n_{d\bar{\sigma}} > n_{d\sigma} + U n_{d\bar{\sigma}} < n_{d\sigma} > + U < n_{d\bar{\sigma}} > < n_{d\sigma} >$$

Energies depend on fillings $< n_{\uparrow} >$ and $< n_{\downarrow} >$

$$E_{\sigma} = \epsilon_d + U < n_{\bar{\sigma}} >$$

Non-vanishing local magnetic moment

$$\mu = < n_{\uparrow} > - < n_{\downarrow} >$$

Total density of states \longrightarrow



Figure: (Anderson 1961)

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Energy scales

- $\epsilon_d, U \gg \Delta \gg T$ atomic limit
- ▶ $U \ll \Delta$ non-magnetic regime
- ► $\epsilon_d, U \gg \Delta, T > T_K$ local moment regime
- ► $\epsilon_d, U \gg \Delta, T < T_K$ strong coupling regime (Kondo effect)
- $\Delta \ge |\epsilon_d|$ mixed valence regime



Local moment approach (Logan) – LMA

Unrestricted Hartree-Fock solution dependent explicitly on local moment μ:

$$G_{\sigma}^{HF}(\mu,\omega) \xrightarrow{\mathbf{G}_{\sigma}^{\mathbf{A}}(|\mu|,\omega)} \bigoplus_{\mathbf{G}_{\sigma}^{\mathbf{B}}(|\mu|,\omega)} \bigoplus_{\mathbf{G}^{\mathbf{B}}(|\mu|,\omega)} \bigoplus_{\mathbf{G}^$$

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• RPA perturbation series around the UHF solution: $G_{\sigma}^{A,BHF}$



Local moment approach (LMA)

 Using RPA we obtain two selfenergies Σ^A_σ and Σ^B_σ and two LMA Green functions:

$$G^{A}_{\sigma}(\Sigma^{A})^{LMA}, \quad G^{B}_{\sigma}(\Sigma^{B})^{LMA}$$

Symmetry restoration

$$G_{\sigma}(\omega)^{LMA} = \frac{1}{2} \left(G_{\sigma}^{A}(\omega)^{LMA} + G_{\sigma}^{B}(\omega)^{LMA} \right)$$

Minimization of the ground state energy with respect to the local moment value |µ| and n

$$E_{\text{phys.}} = \min_{|\mu|,n} E(|\mu|, n).$$

Local moment approach in SIAM

- Reproduces Hubbard bands (charge fluctuations)
- Describes well Kondo resonance near the Fermi level (spin fluctuations), reproducing its exponential scale
- Conservative (there exists Luttinger Ward functional for this theory)
- Numericaly efficient





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Multiorbital Anderson model and LMA generalization

$$\begin{split} H_{\mathrm{SIAM}} &= \sum_{\mathbf{k},\sigma} \epsilon_{\mathbf{k}} c_{\mathbf{k}\sigma}^{\dagger} c_{\mathbf{k}\sigma} + \sum_{\mathbf{k},\sigma,\alpha} V_{\mathbf{k}\alpha} \left(d_{\alpha\sigma}^{\dagger} c_{\mathbf{k}\sigma} + c_{\mathbf{k}\sigma}^{\dagger} d_{\alpha\sigma} \right) \\ &+ \sum_{\alpha,\sigma} \left(\epsilon_{\alpha} - U_{\alpha} n_{\alpha,\bar{\sigma}} \right) n_{\alpha,\sigma} + \sum_{\sigma,\sigma'} \sum_{\alpha\neq\beta} \left(U_{\alpha\beta}' - J \delta_{\sigma\sigma'} \right) n_{\alpha\sigma} n_{\beta\sigma'} \end{split}$$

- RPA perturbation series around the UHF solutions, that are now dependent on local moments on the orbitals μ_α
- Symmetry restoration

$$G^{lphaeta}_{\sigma}(\omega) = rac{1}{2N_{
m orb}}\sum_{\pm|\mu_{lpha}|,\pm|\mu_{eta}|}G^{lphaeta}_{\sigma}(\mu_{lpha}\mu_{eta};\omega)$$

• Energy minimization with respect to μ_{α} and n_{α}

$$E_{\mathrm{phys.}} = \min_{\mu_{\alpha}, n_{\alpha}} E(\mu_{\alpha}, n_{\alpha}).$$

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Multiorbital SIAM and LMA

Spectral functions

$$A^{lphalpha}_{\sigma}(\omega)=-rac{1}{\pi}{\it Im}G^{lphalpha}_{\sigma}$$

- ► The influence of Hund's rule J ≠ 0 suppresses Kondo peak
- Differences in spectral functions between different orbials
- Possible investigation of transport in multiorbital quantum dots



Hubbard model

interaction U (repulsion)
 between electrons on each site

$$H_{
m Hubbard} = \sum_{ij\sigma} t_{ij} c^{\dagger}_{i\sigma} c_{j\sigma} + U \sum_{i} n_{i\uparrow} n_{i\downarrow}$$

 solution (approx.) within the scheme of dynamical mean field theory



Dynamical mean field theory (DMFT)

- Hubbard model is mapped onto a model of a single impurity in mean field G₀(iω_n)
- this mean field is frequency dependent (dynamical)



DMFT Green function



$$G(\mathbf{k}, i\omega_n)^{-1} = i\omega_n + \mu - \epsilon_{\mathbf{k}} - \Sigma^{LOC}(i\omega_n)$$

self-consistency condition

$$\Sigma^{LOC}(i\omega_n) = \mathcal{G}_0(i\omega_n)^{-1} + G^{LOC}(i\omega_n)^{-1}$$

oraz $\mathcal{G}_0 = \mathcal{G}_0[G^{LOC}]$

Dynamical mean field theory (DMFT)

- Thermodynamicaly consistent and conservative
- Exact in the non-trivial limit of $d = \infty$
- ► Frequency dependent mean field (Weiss field G₀(iω_n)) describes spin and charge fluctuations on sites
- Exact in the atomic (t_{i,j} = 0) and non-interacting (U = 0) limits
- Requires an impurity solver to solve self-consistent equations the solver must work for arbitrary hybridization function (e.g. NRG, QMC, IPT, NCA) – LMA can serve as an impurity solver

$\mathsf{DMFT} + \mathsf{LMA}$

 Mott-Hubbard metal - insulator transition



- Suppression of Kondo peak when J ≠ 0 (Hund's rule)
- Orbital selective metal insulator transition



(Bethe lattice at half-filling)

Summary

Local moment approach (LMA):

- Conservative, numerically efficient method to treat multiorbital Anderson and Hubbard models
- Can be used to invesigate:
 - Mott- Hubbard metal insulator transition in multiorbital Hubbard model
 - transport in multilevel quantum dots that are described by Anderson impurity model