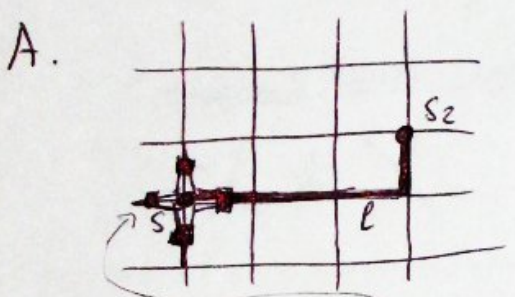


# Lecture 2

## FAQ regarding the previous lecture

Q. Why does the path operator create a pair of charges?



$$W_e = \prod_{j \in \ell} \sigma_j^z$$

$$|\psi_{s,s_2}\rangle = W_e |\psi_0\rangle$$

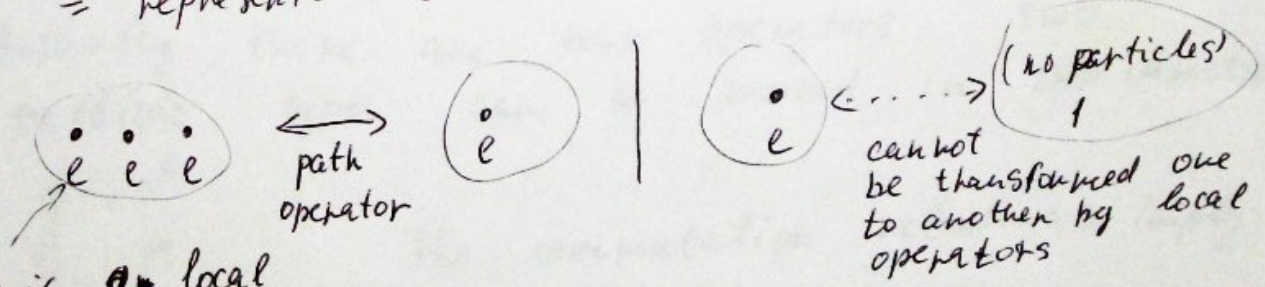
$$A_{s_1} = \prod_{i \in \text{Star}(s_1)} \sigma_i^x$$

$$W_e A_{s_1} = -A_{s_1} W_e$$

$$A_{s_1} |\psi_{s,s_2}\rangle = A_{s_1} W_e |\psi_0\rangle = -W_e A_{s_1} |\psi_0\rangle = -W_e |\psi_0\rangle = -|\psi_{s,s_2}\rangle$$

Q. What are these labels:  $1, e, m, \epsilon$ ?

A. These are superselection sectors = equivalence classes of excited states w.r.t. local operators = representations of the operator algebra.



this is a local excited state.

cannot be transformed one to another by local operators

~~These are superselection sectors~~

Q. How does one read a braid diagram? |2-2

A.  $e \circ e = \begin{array}{c} \diagup \\ \diagdown \end{array}$  exchange  $m \circ e = \begin{array}{c} | \\ \diagdown \\ | \\ \diagup \\ | \end{array}$  full turn

bosons

$$\begin{array}{c} \diagup \\ \diagdown \end{array} = \begin{array}{c} \diagdown \\ \diagup \end{array} = \begin{array}{c} ) \\ ( \end{array}$$

fermions

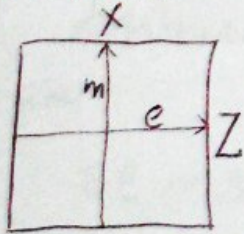
$$\begin{array}{c} \diagup \\ \diagdown \end{array} = \begin{array}{c} \diagdown \\ \diagup \end{array} = - \begin{array}{c} ) \\ ( \end{array}$$

anyons with nontrivial self-statistics (not present in our model)

$$\begin{array}{c} \diagup \\ \diagdown \end{array} \neq \begin{array}{c} \diagdown \\ \diagup \end{array}$$

Einarsson's argument (1990)

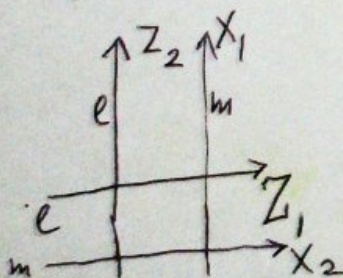
The degeneracy on the torus follows from nontrivial statistics



$$Z^{-1} X^{-1} Z X = \text{circle with } m \text{ and } e = -1$$

These are two noncommuting operators acting on the ground space  $\mathcal{L} \Rightarrow \dim \mathcal{L} > 1$

Actually, there are four operators: two particle types can be moved in two directions.



The commutation relations imply that  $\dim \mathcal{L} = 4$

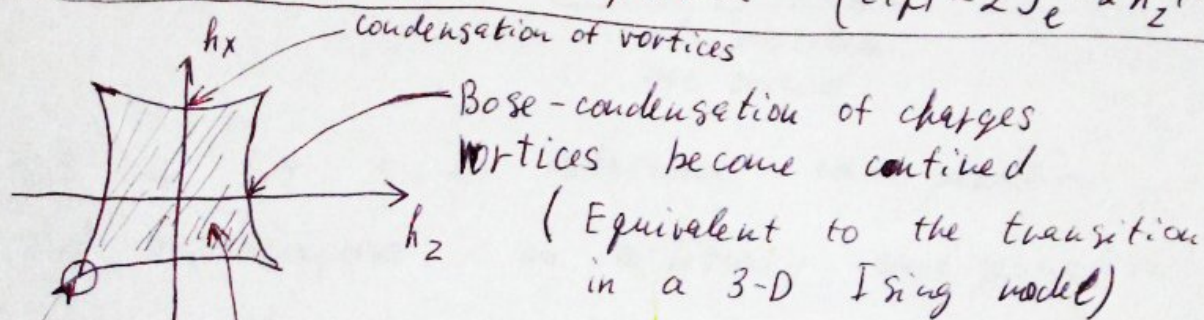
# Perturbation analysis

12-3

$$H = -J_e \sum_{\text{vertices}} A_s - J_m \sum_{\text{plaquettes}} B_p - \underbrace{\sum_{\text{edges}} (h_x \sigma_j^x + h_z \sigma_j^z)}_{\text{perturbation}}$$

$h_z \neq 0 \Rightarrow$  electric charges can hop from site to site

$\Rightarrow$  nontrivial dispersion ( $\varepsilon(q) \approx 2J_e - 2h_z(\cos q_x + \cos q_y)$ )



trichritical point (not studied) topological phase

Within the topological phase, the ground state degeneracy is preserved with exponential precision:

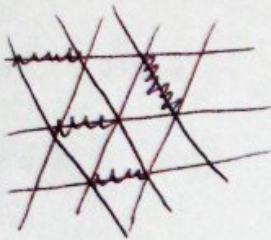
$$\delta E \sim \Delta \underbrace{\exp(-L/\xi)}$$

due to virtual quasiparticle tunneling:

the only way to act upon the ground state is to wind a quasiparticle around the torus

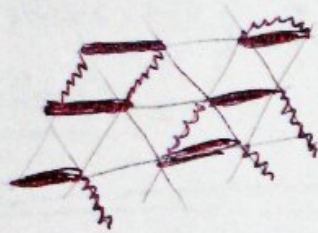
# Analogy to the dimer model on the triangular lattice

L2-4



Dimer liquid =  
= uniform superposition of all  
dimer covers  
all vertices  
are covered

Let us fix some reference configuration  
and superimpose an arbitrary configuration



$$x_j = 0 \quad \text{---} \quad \text{or} \quad \text{---}$$

$$x_j = 1 \quad \text{---} \quad \text{or} \quad \text{---}$$

$$u_s(z) = \sum_{j \in \text{star}(s)} x_j = 0 \pmod{2}$$

$$|\Psi_{\text{dimer}}\rangle = \sum_{\text{Some loop conf.}} |\bar{x}\rangle$$

Now, consider

$$|\Psi_0\rangle = \sum_{\substack{\bar{x}: u_s(\bar{x})=0 \\ \text{for all } s}} |\bar{x}\rangle$$

This is the toric code state in the dual basis

$$x=0 = \frac{|0\rangle + |1\rangle}{\sqrt{2}} = \frac{|\uparrow\rangle + |\downarrow\rangle}{\sqrt{2}}$$

$$x=1 = \frac{|0\rangle - |1\rangle}{\sqrt{2}}$$

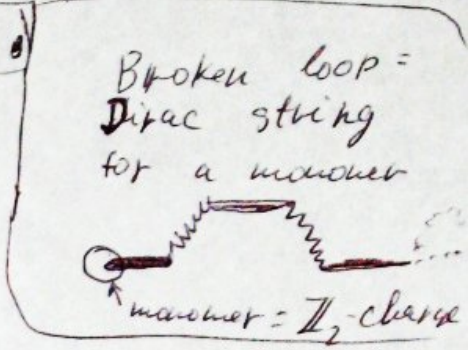
$$u_s(\bar{x}) = 0 \Leftrightarrow A_s |\Psi\rangle = |\Psi\rangle$$

~~The basis = Broken loop~~

Properties of the dimer state are similar, but not identical, to the toric code state.

Not so on the square lattice

Loops have direction  $\Rightarrow$  integer charges  $\Rightarrow \mathbb{Z}_2 \rightarrow U(1)$



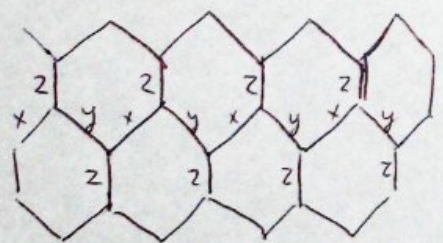
It is known that the  $U(1)$ -theory (= electrodynamics) is unstable in 2D.

Vortices have arbitrary value in  $U(1) \Rightarrow$  no gap.

Moreover, the vortices spontaneously condense.

Reason:  $S_{\text{monopole}} < \infty$  (Polyakov)

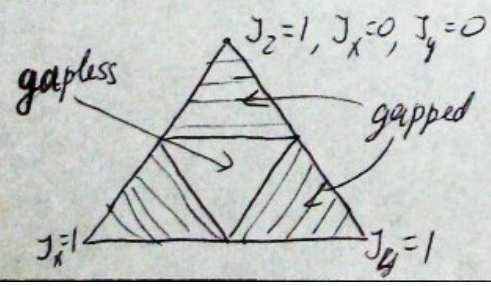
The honeycomb lattice model



Spins on the sites

$$H = -J_x \sum_{x\text{-links}} \sigma_j^x \sigma_k^x - J_y \sum_{y\text{-links}} \sigma_j^y \sigma_k^y - J_z \sum_{z\text{-links}} \sigma_j^z \sigma_k^z$$

Answer: The signs of  $J_x, J_y, J_z$  don't matter



The gapped phases are in the same univ. class as the toric code.

But: subtle breaking of the translational symmetry.