

Multiferroics for spintronics

Manuel Bibes²

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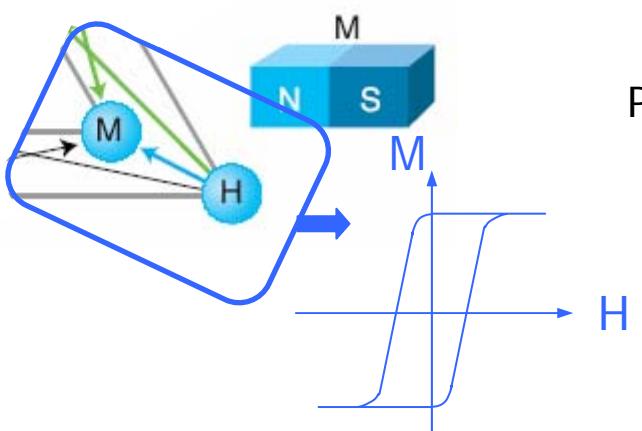
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Ferromagnetic materials

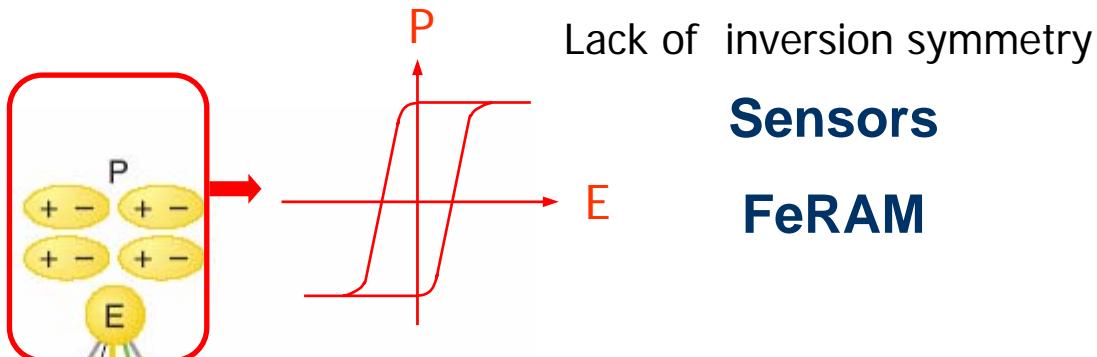


Partially filled d/f shells

Magnetic Data Storage
Spintronics

Switchable parameter : M
External stimulus : H

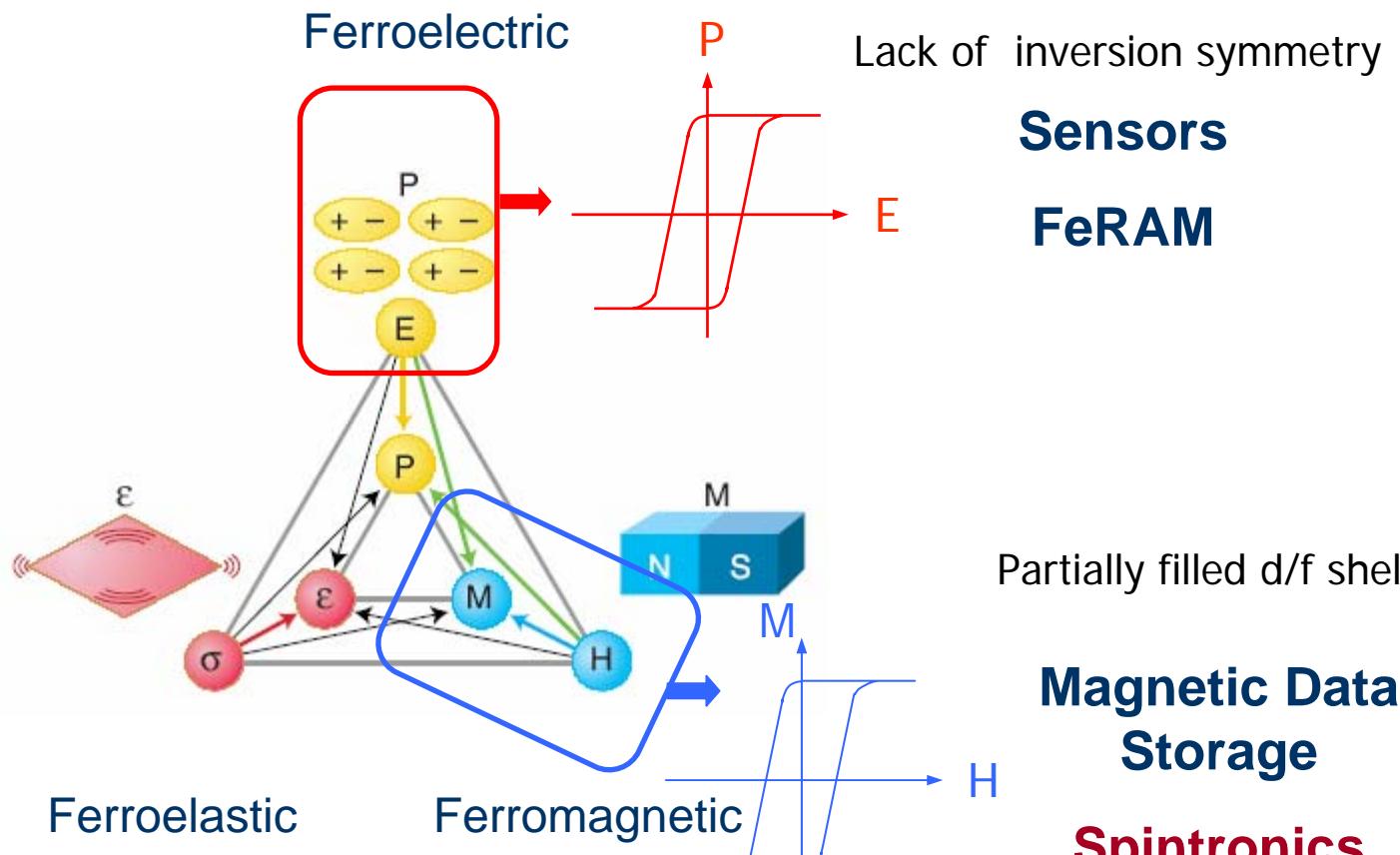
Ferroelectric materials



Sensors
FeRAM

Switchable parameter : P
External stimulus : E

Multiferroic materials

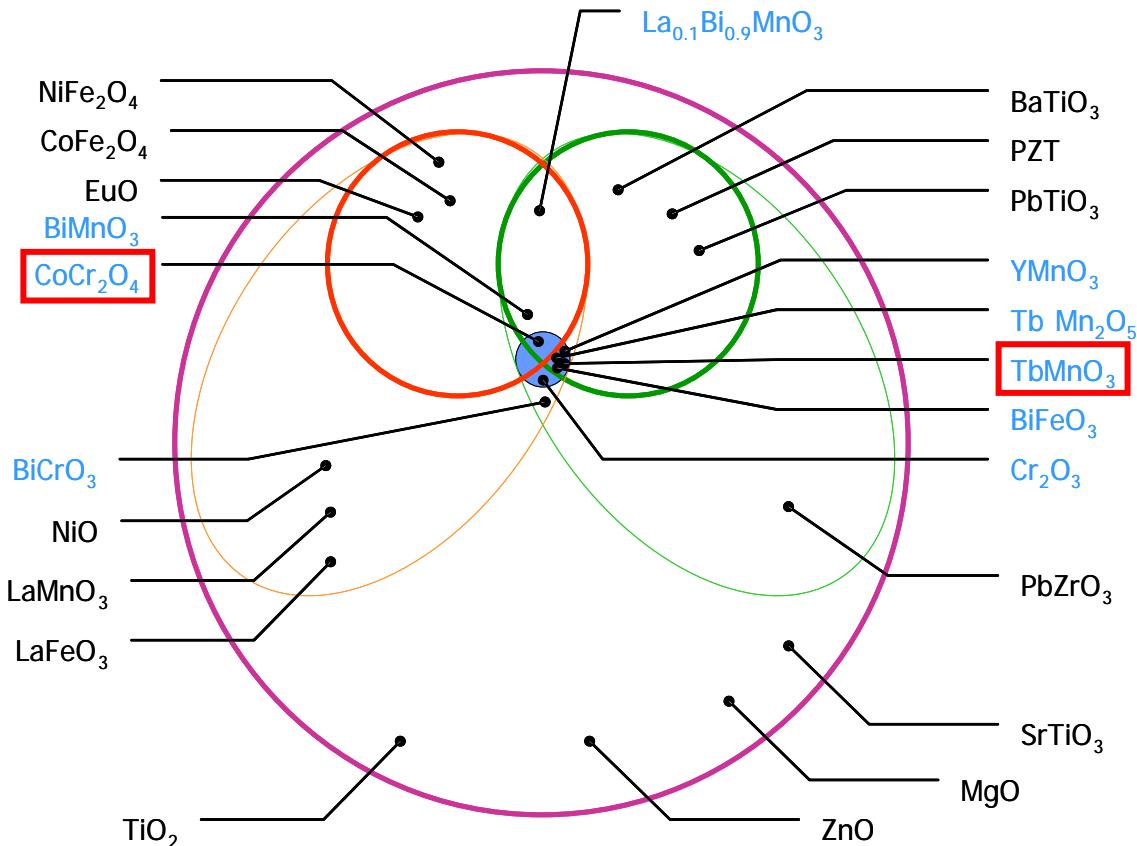


From Spaldin and Fiebig, Science 309, 391 (2005)

**Switchable parameter : M, P
External stimulus : H, E**

- Multiple-state memories
- M switching by E
- P switching by H

Insulating oxides : a classification



Magnetically polarizable
Ferro(i)magnetic

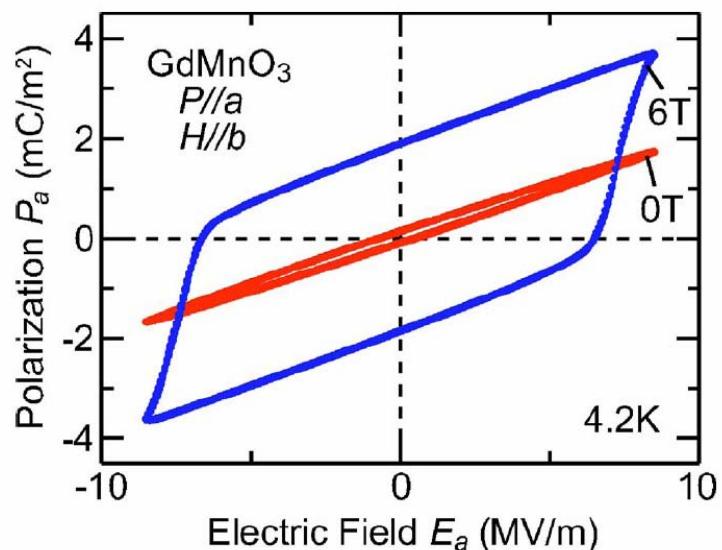
Electrically polarizable
Ferro(i)electric

Magnetoelectric

Derived from Eerenstein, Mathur and Scott, Nature 442, 759 (2006)

(Gd,Dy,Tb)MnO₃ : spiral magnets

Spiral ordering breaks inversion symmetry
 → (improper) ferroelectricity appears
 (but **no** finite magnetization)

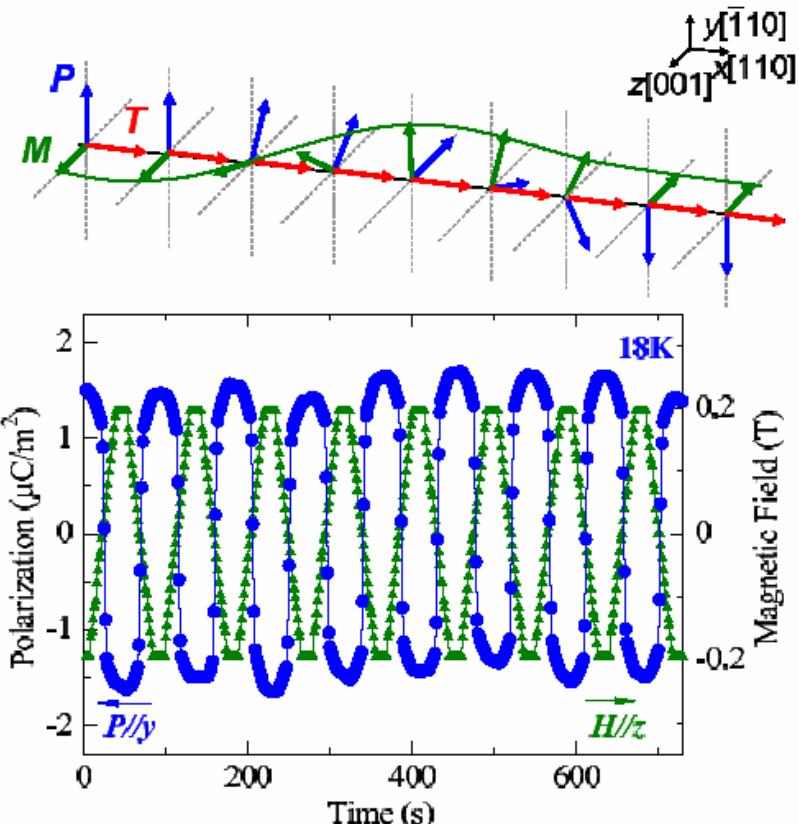


Apply a magnetic field : change magnetic state
 → switch polarization on/off

Kimura et al, PRB 71, 224425 (2005)

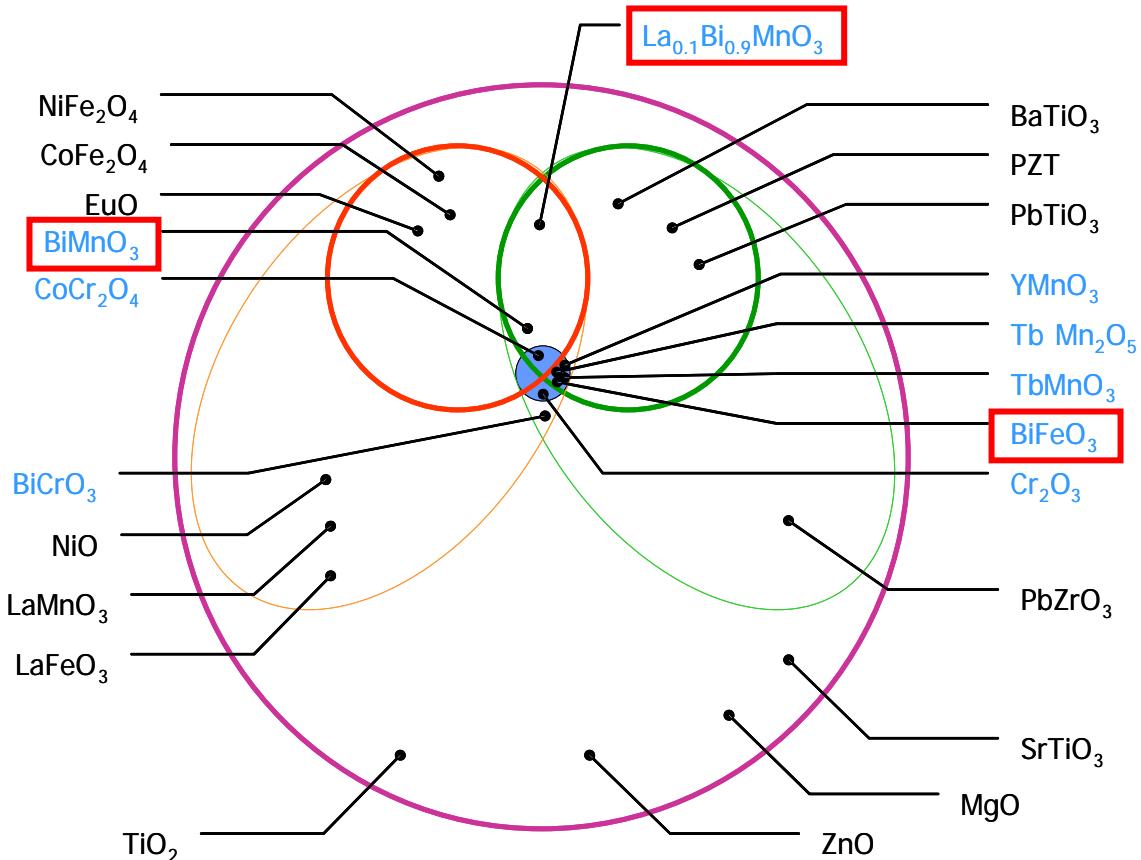
CoCr₂O₄ : conical magnet

Conical ordering breaks inversion symmetry
 → (improper) ferroelectricity appears
 (**with** a finite magnetization)



Flip polarization direction by a magnetic field
 Yamasaki et al, PRL 96, 207204 (2006)

Insulating oxides : a classification



Magnetically polarizable
Ferro(i)magnetic

Electrically polarizable
Ferro(i)electric

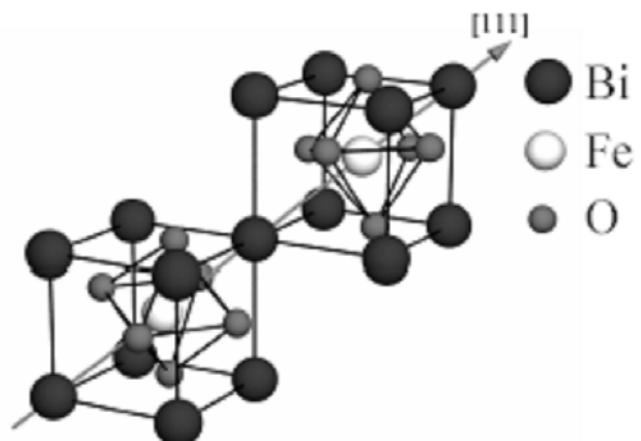
Magnetoelectric

Derived from Eerenstein, Mathur and Scott, Nature 442, 759 (2006)



A room-temperature antiferromagnetic ferroelectric

Rhombohedral Perovskite (R3c)
 $a=3.96\text{\AA}$, $\alpha=89.4^\circ$



Neaton et al. Phys. Rev. B, 71, 014113 (2005)

Ferroelectric

$T_C=1100\text{K}$

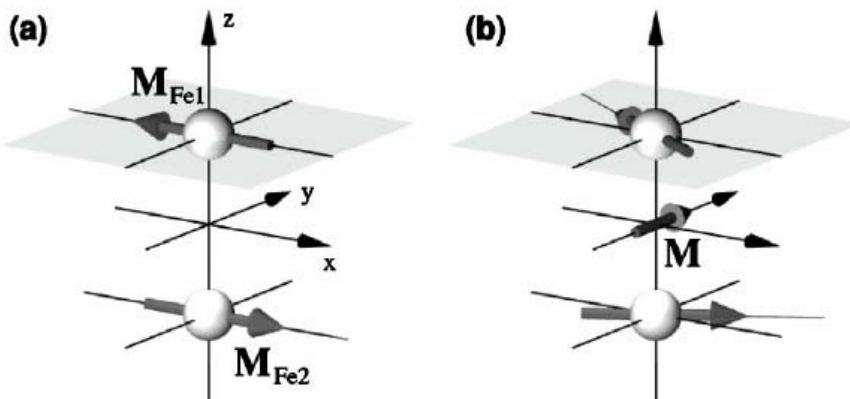
$P_S=6\mu\text{C}/\text{cm}^2$ (1970 paper, bad crystal)

J. R. Teague et al., Solid State Commun., 8, 1073 (1970)

$P_S=60\mu\text{C}/\text{cm}^2$ (2007 paper, good crystal)

D. Lebeugle, M. Viret et al, PRB in press

Condmat/0706.0404



Ederer et Spaldin, Phys. Rev. B, 71, 060401 (R)
(2005)

G-type Antiferromagnetic

$T_N=640\text{K}$

Canted spins \rightarrow weak FM

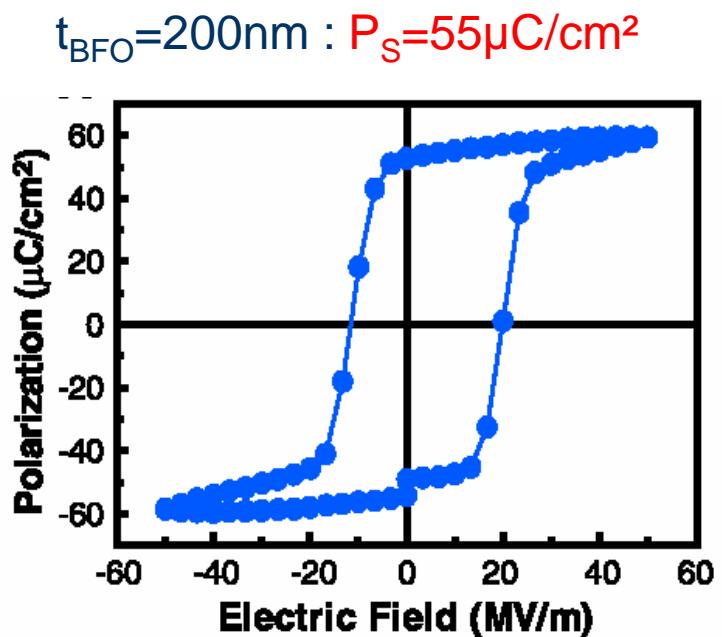
$M_S=0.01\mu_B/\text{f.u.}$

Incommensurate cycloidal modulation

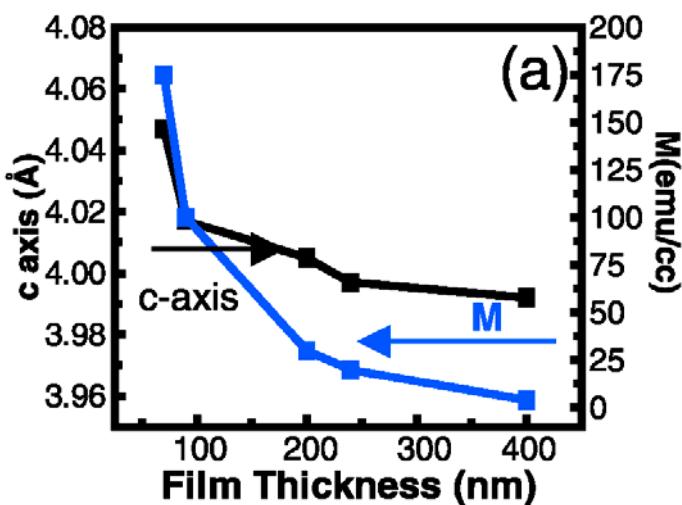
P. Fisher et al., J. Phys. C, 13, 1931 (1980)

Popov et al. in Magnetoelectric Interaction Phenomena in Crystals (NATO Science Series, 2004) p. 277

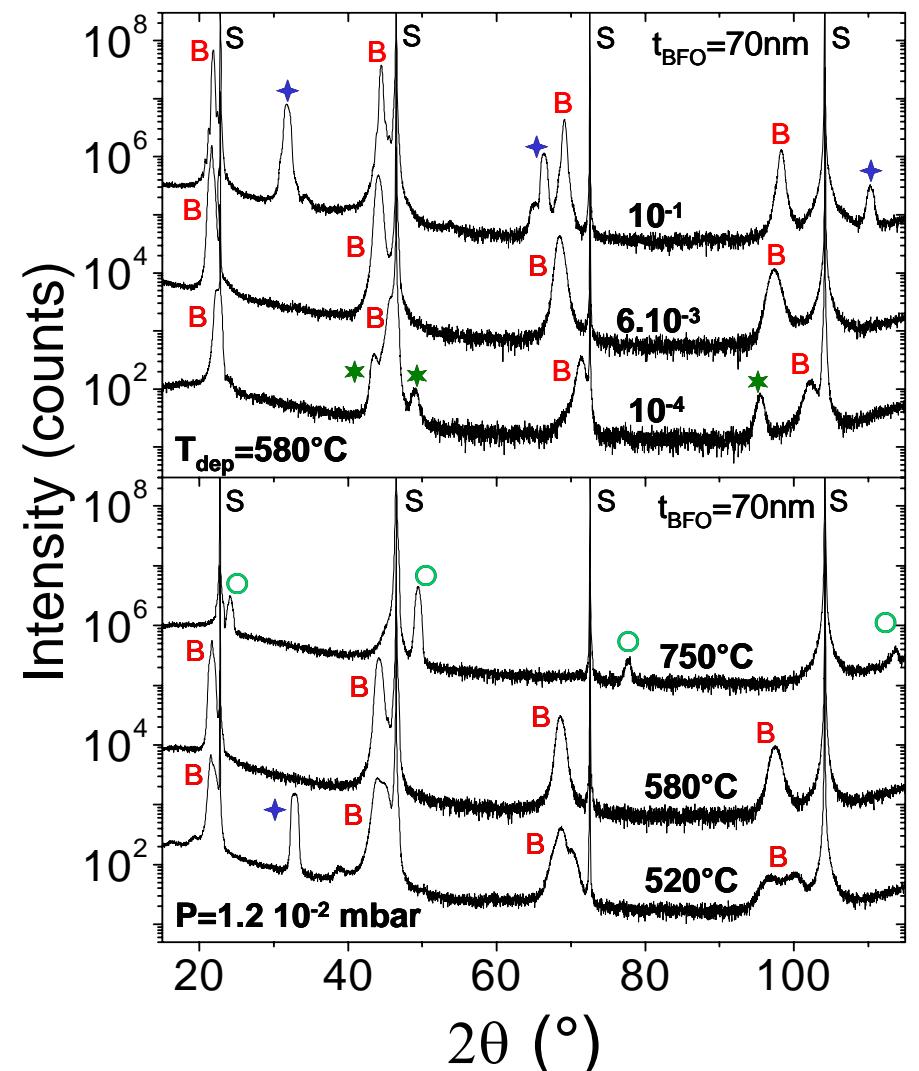
Wang et al., Science, 299, 1719 (2003) : BFO//STO (001)



$t_{\text{BFO}}=70\text{nm} : M_s=150\text{emu}/\text{cm}^3 (\sim 1\mu_B/\text{f.u.})$
 $t_{\text{BFO}}=400\text{nm} : M_s=5\text{emu}/\text{cm}^3 (0.03\mu_B/\text{f.u.})$



Claim of enhanced polarization and magnetization compared to bulk



Pulsed laser deposition on SrTiO₃ (001)

Target : Bi_{1.15}FeO₃

Growth conditions : $T = 580^\circ\text{C}$,
 $P_{O_2} = 6.10^{-3} \text{ mbar}$

Bi much more volatile than Fe :

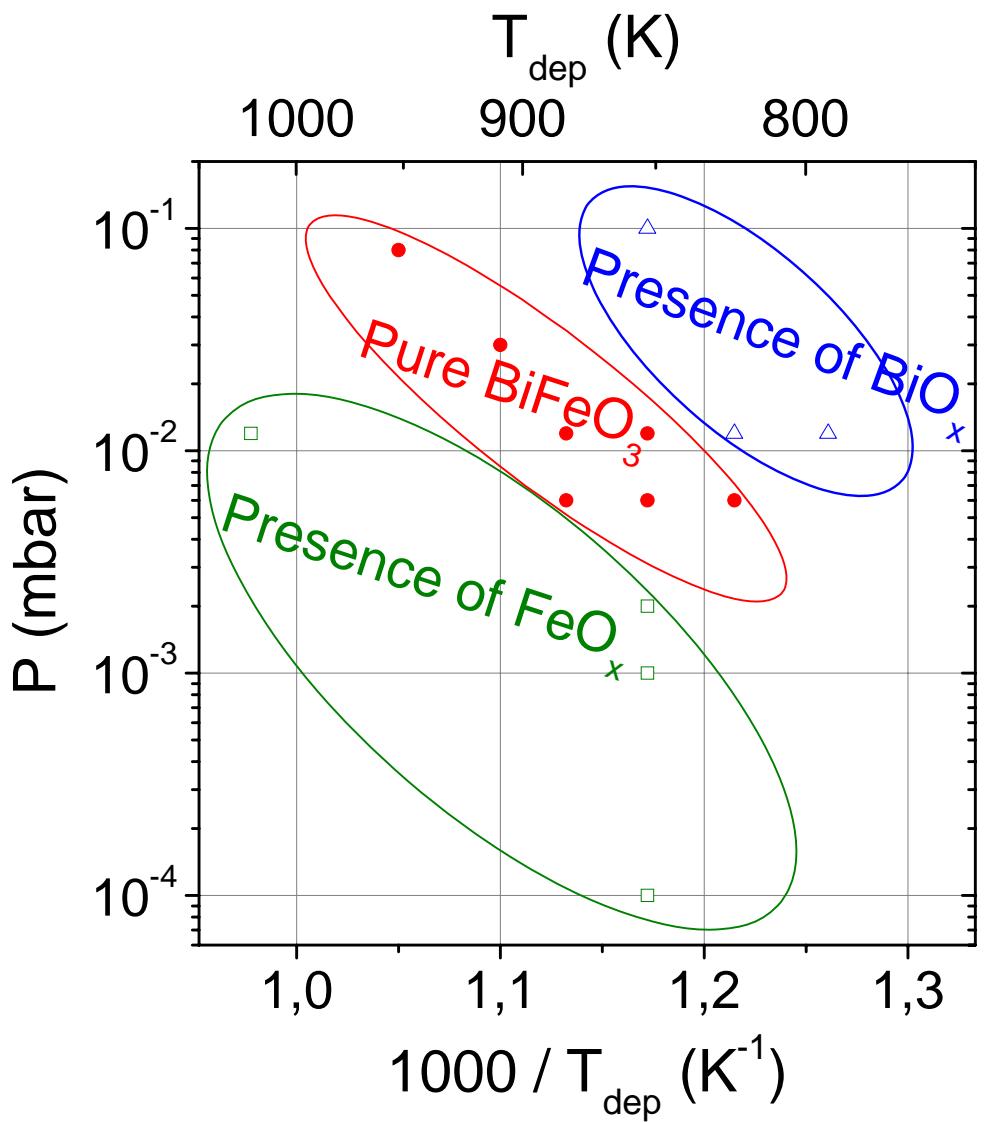
- low P or high T (Bi evaporates more than Fe)

\Rightarrow **Fe oxides**

- high P or low T (excess Bi)

\Rightarrow **Bi oxides**

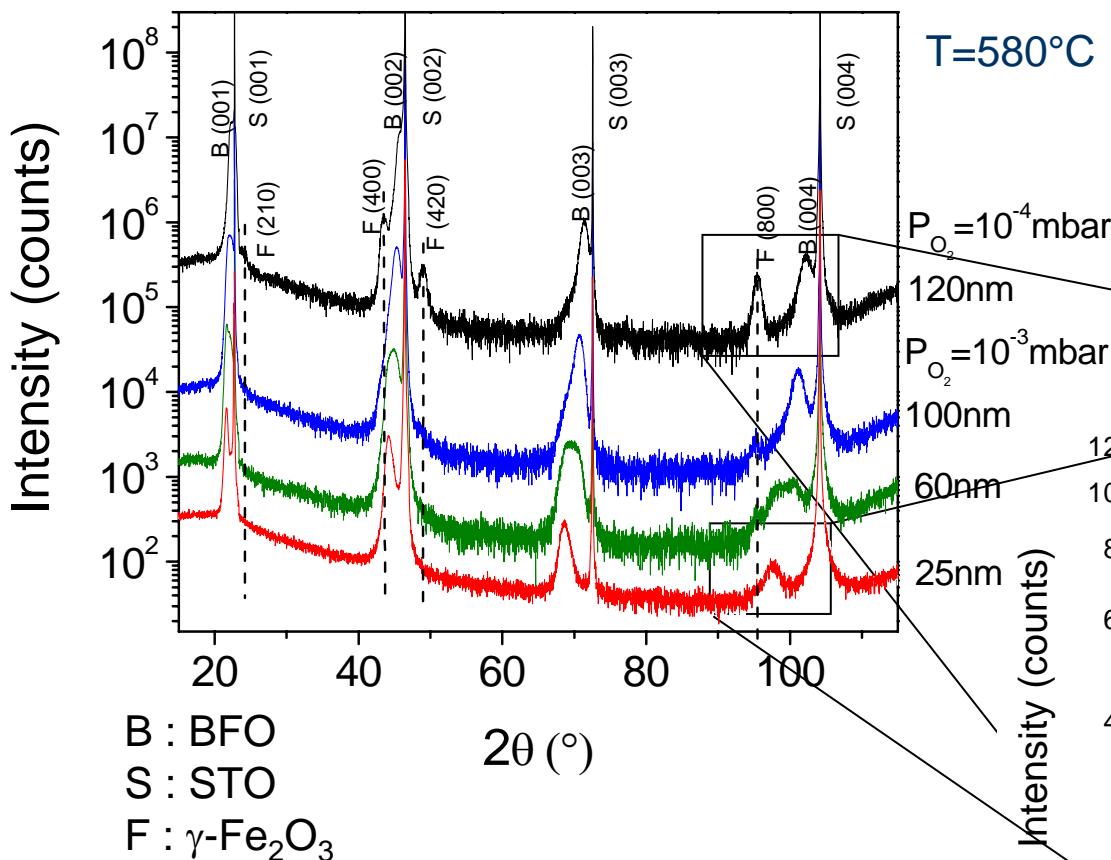
Phase diagram



Pure BiFeO_3 obtained in a very
narrow region around $T_{\text{dep}}=580^\circ\text{C}$
and $P_{\text{O}_2}=6.10^{-3}\text{mbar}$

Appl. Phys. Lett., 87, 072508 (2005)

X-ray diffraction



$T = 580^{\circ}\text{C}$

$P_{\text{O}_2} = 10^{-4} \text{ mbar}$:

120nm

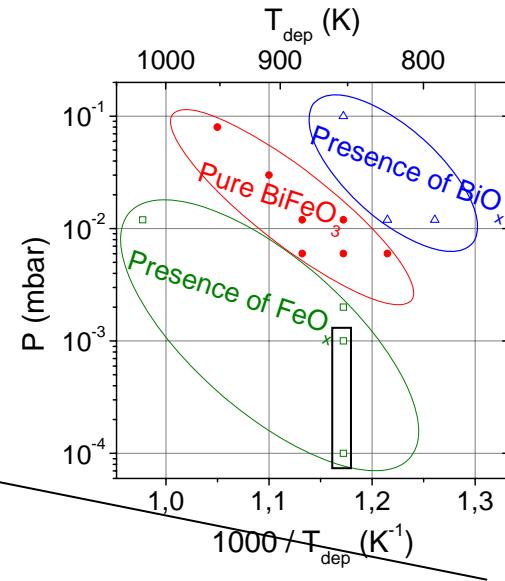
$P_{\text{O}_2} = 10^{-3} \text{ mbar}$:

100nm

60nm

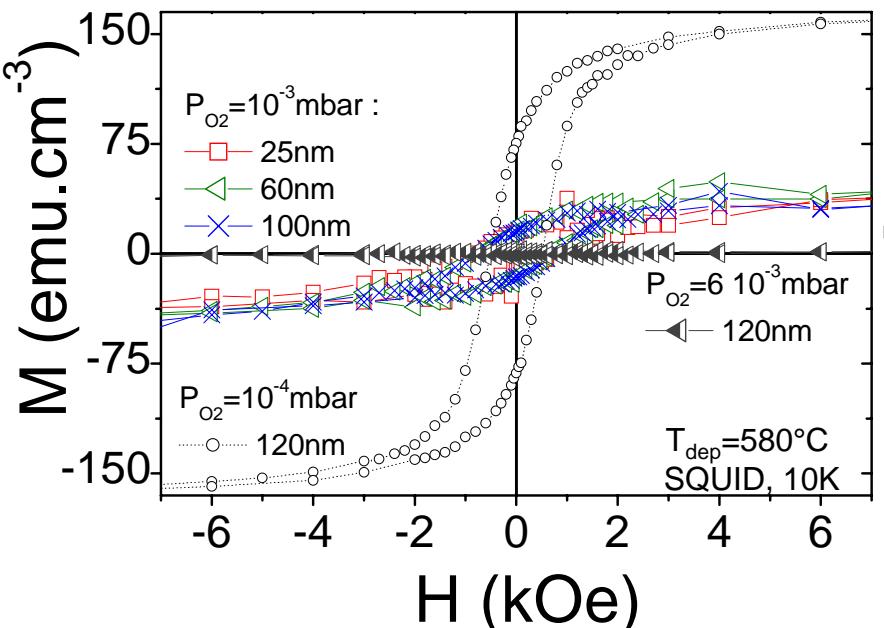
25nm

25nm
 $P = 10^{-3} \text{ mbar}$



Quantification of γ -Fe₂O₃ by XRD : **not easily detectable**

SQUID measurements

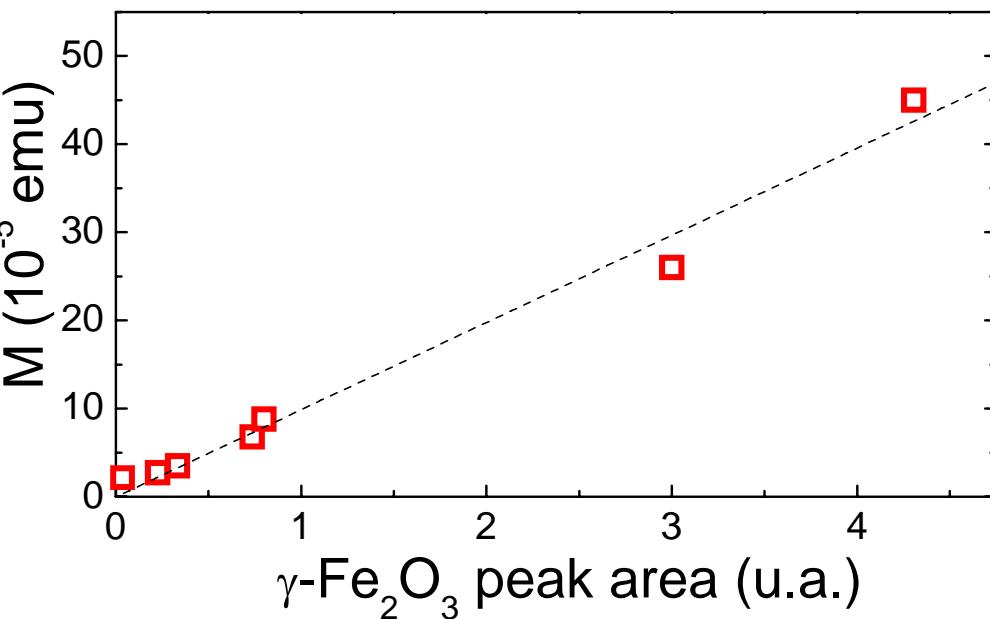


BFO + $\gamma\text{-Fe}_2\text{O}_3$: up to $M_s \sim 150 \text{ emu/cm}^3$
 $\gamma\text{-Fe}_2\text{O}_3$ ferrimagnetic ($M_s = 430 \text{ emu/cm}^3$)
Pure BFO : $M_s \sim 2 \text{ emu/cm}^3$

Phys. Rev. B, 74, 020101R (2006)

Similar conclusions :

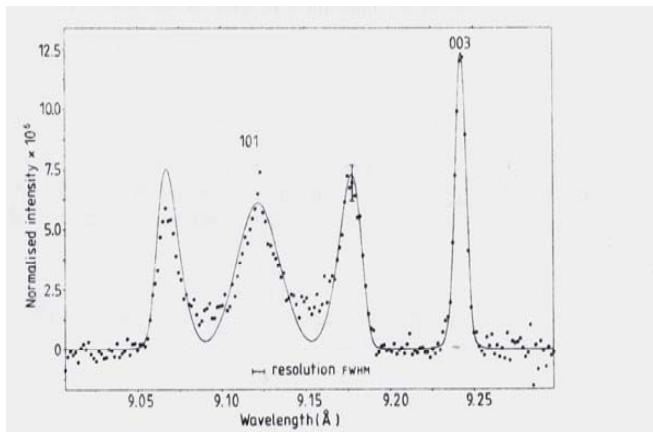
- Eerenstein et al. Science, 307 1203a (2005)
 Wang et al., Science, 307, 1203b (2005)



Magnetic moment comes from $\gamma\text{-Fe}_2\text{O}_3$

BFO weak bulk-like ferromagnetic moment

Neutron diffraction



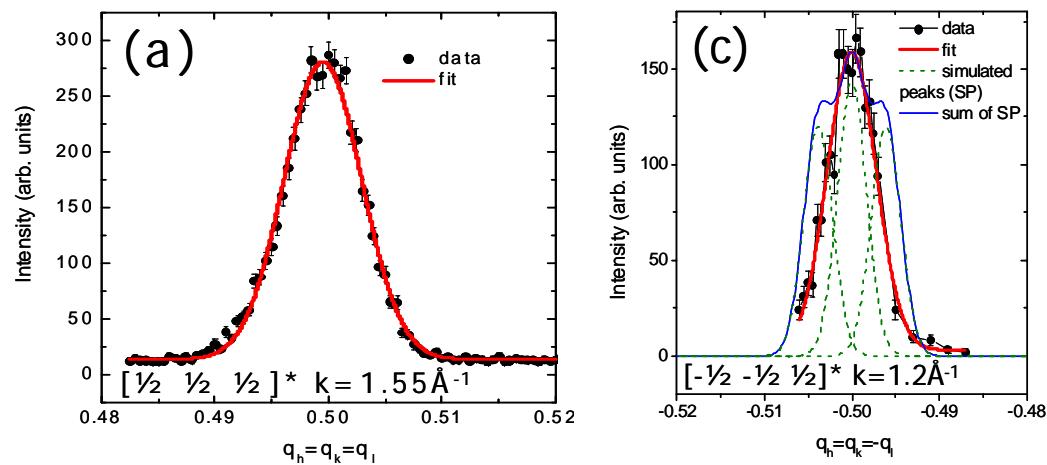
In R3c bulk BFO :

$[003]^*_H$ single peak \rightarrow **G-type antiferromagnet**
 $[101]^*_H$ peak with satellites \rightarrow **cycloidal modulation**
 \Rightarrow **Averaging to zero** of the linear ME effect

In pseudo-cubic notation : $[003]^*_H \rightarrow [\frac{1}{2} \frac{1}{2} \frac{1}{2}]^*_C$
 $[101]^*_H \rightarrow [-\frac{1}{2} -\frac{1}{2} \frac{1}{2}]^*_C$

from Sosnowska et al., J. Phys. C, 15, 4835 (1982)

BFO(240nm)//STO (001)



➤ **G-type antiferromagnetic order** as in bulk BFO

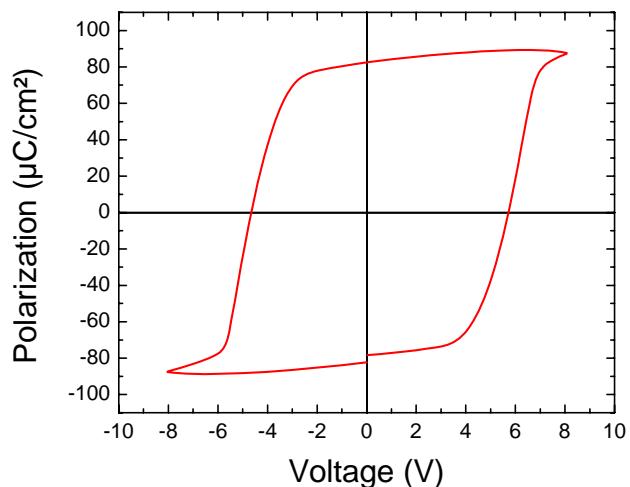
➤ **No cycloidal modulation**

contrary to bulk BFO

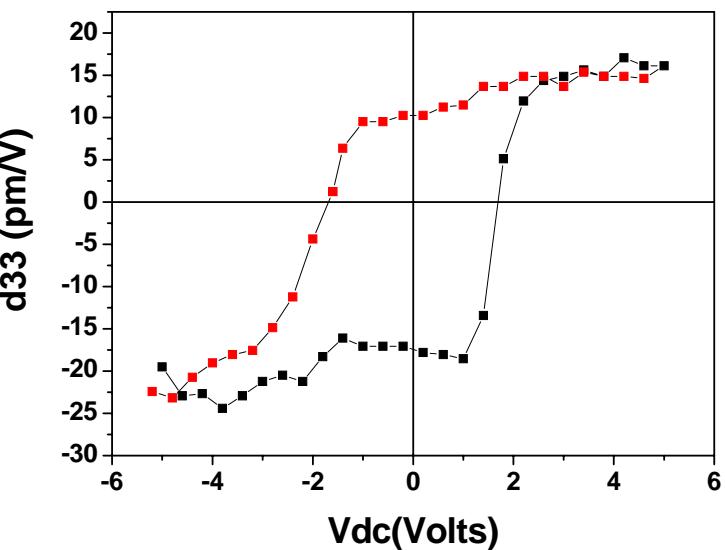
➤ Linear magnetoelectric effect
allowed

Phil. Mag. Lett. 87, 165 (2007)
(Special issue on multiferroic thin films
Eds. N.D. Mathur and MB)

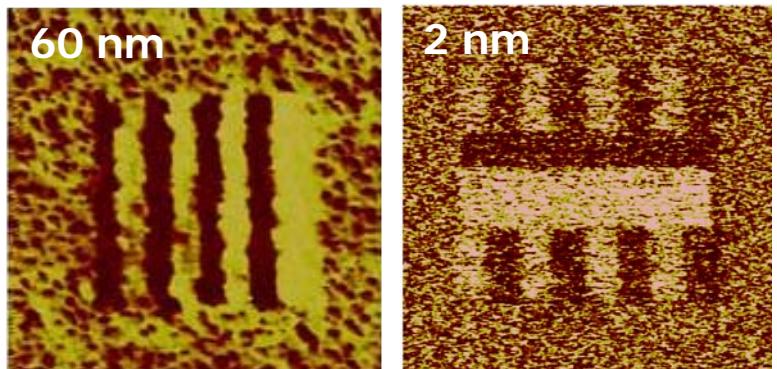
Polarization loops



Piezoelectric loops



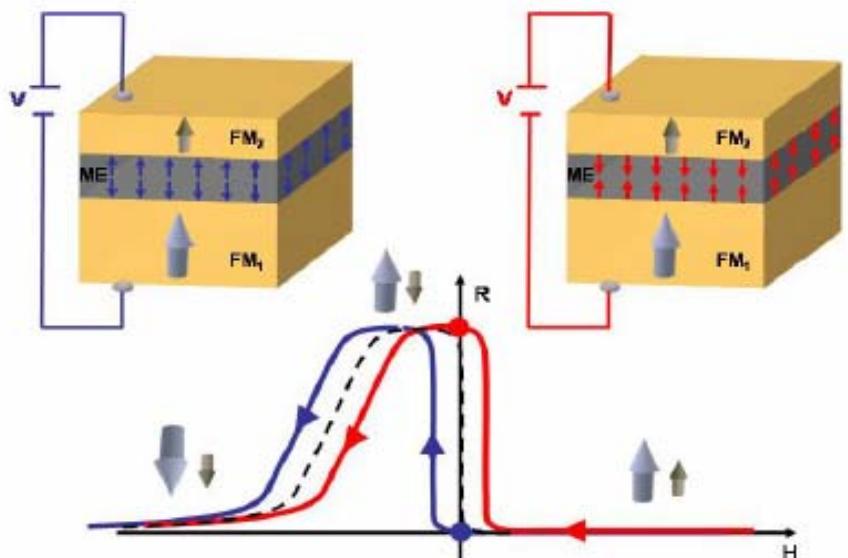
Piezoresponse force microscopy



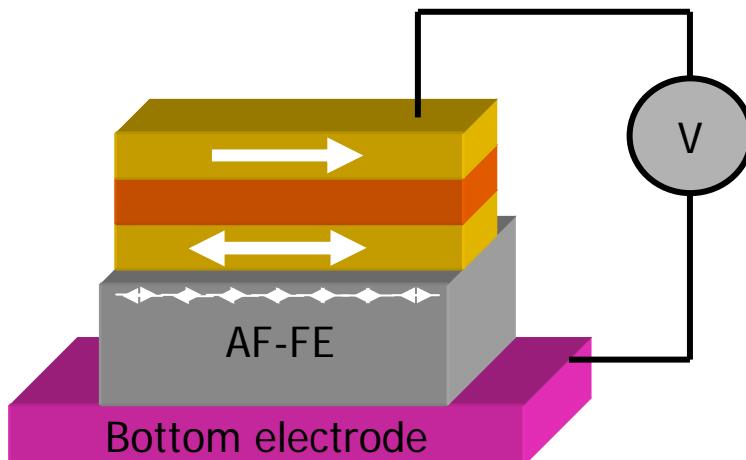
- BFO films are ferroelectric with a large polarization ($\sim 70 \mu\text{C}/\text{cm}^2$) and piezoelectric with a d_{33} coefficient of $\sim 20 \text{ pm/V}$
- Ferroelectricity is preserved down to 2 nm.

Jpn. J. of Appl. Phys., 45, L187 (2006)

Principle : voltage-controlled exchange bias



Binek et Doudin, JPCM, 17, L39 (2005)



Spin-valve on top of a ***multiferroic film*** (FE and AFM)

Magnetic tunnel junction with ***multiferroic barrier*** (FE and AFM)

Electric field control of the junction resistance state

Prerequisites :

1. observe robust exchange bias at room-temperature
2. spin-dependent tunneling through multiferroic barrier
3. switch exchange bias direction by E-field (magnetoelectric coupling)

Exchange bias with BiFeO_3 films

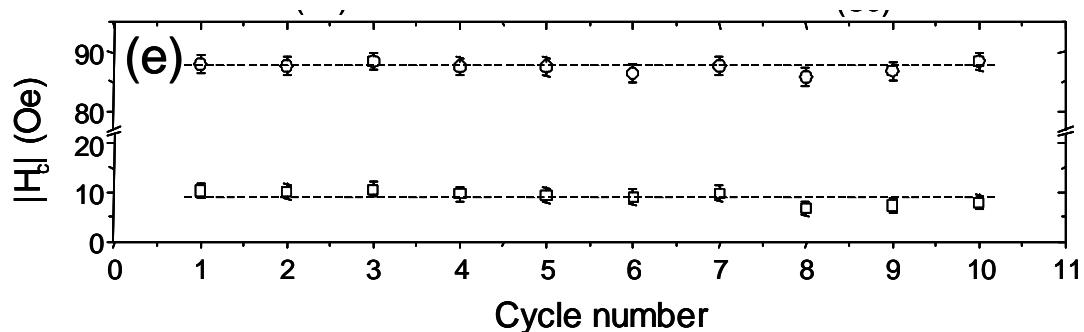
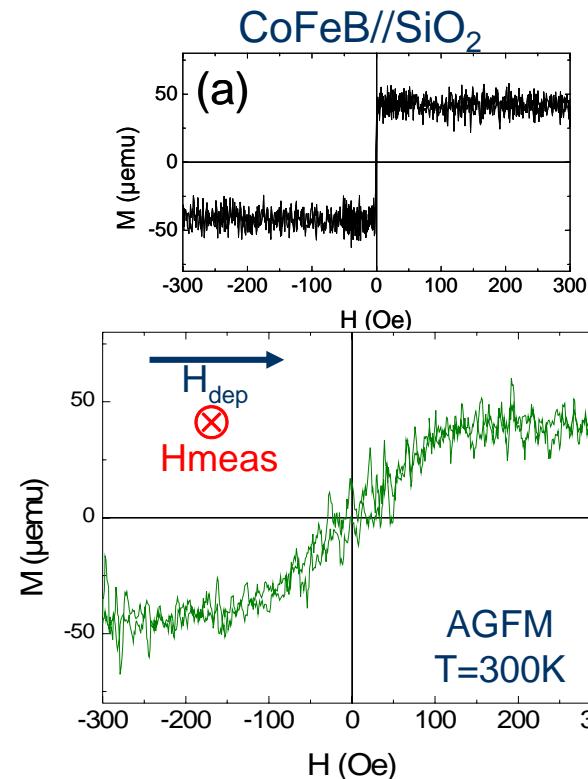
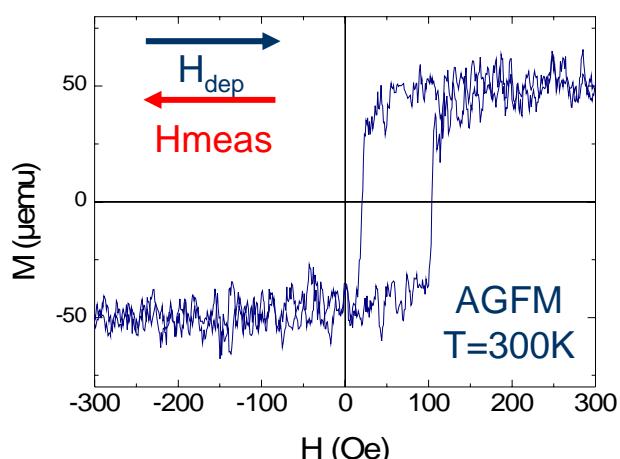
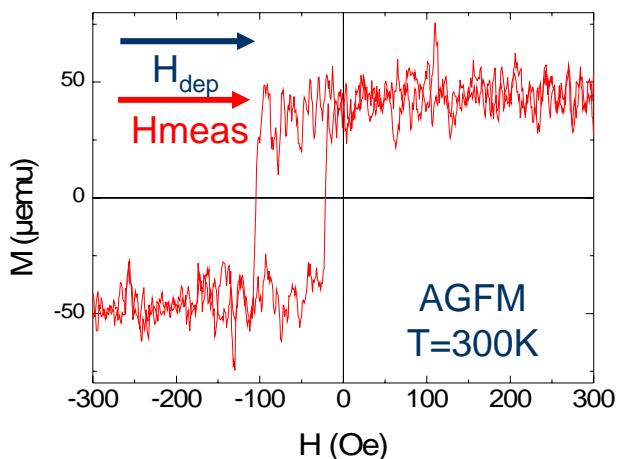
Brought to you by PTP (www.ptp.phas.ubc.ca)



CoFeB deposited by sputtering

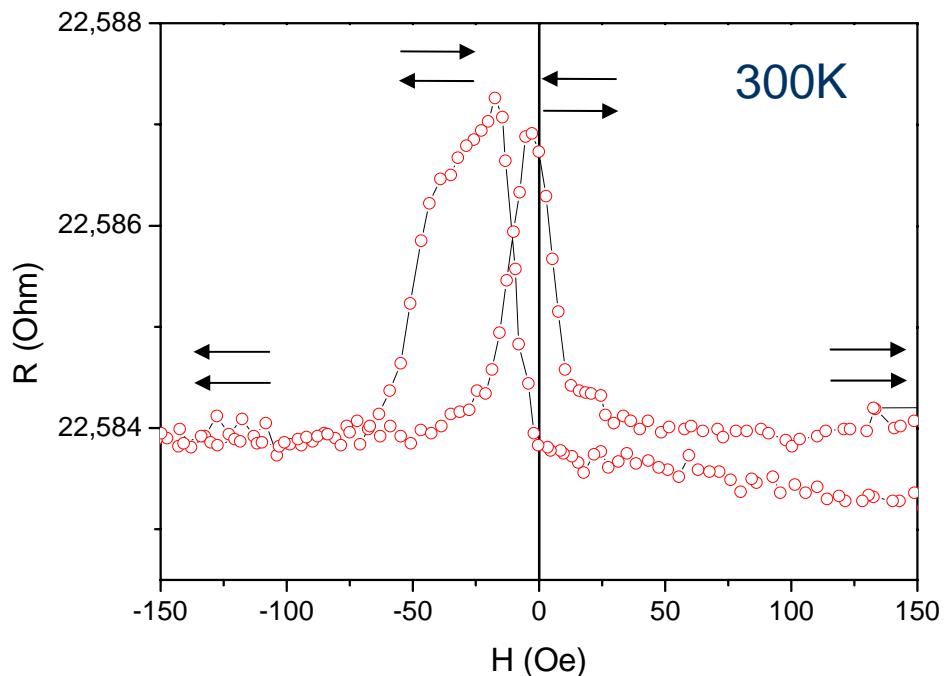
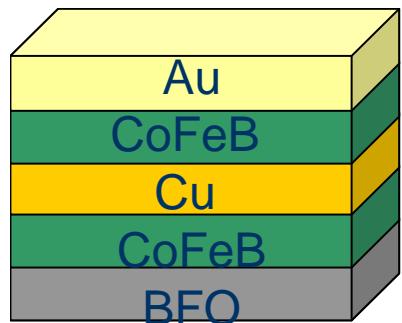
$H_c=42\text{Oe}$, $H_e=-62\text{Oe}$

Appl. Phys. Lett., 89, 242114 (2006)



Robust room temperature exchange bias between BFO and CoFeB

With NiFe : J. Dho et al., Adv. Mat., 18, 1445 (2006)

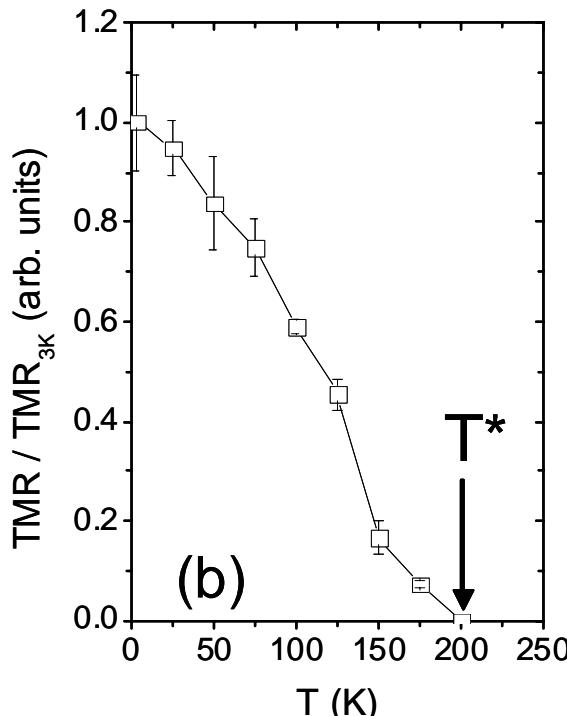
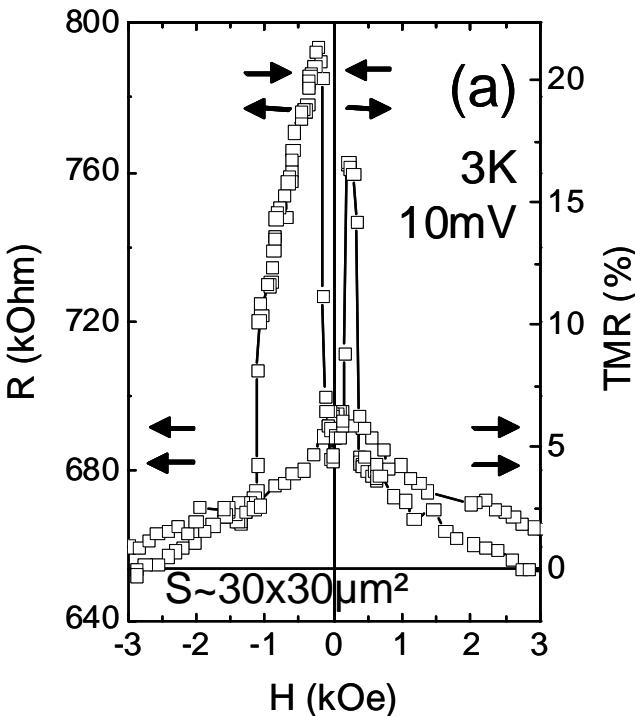
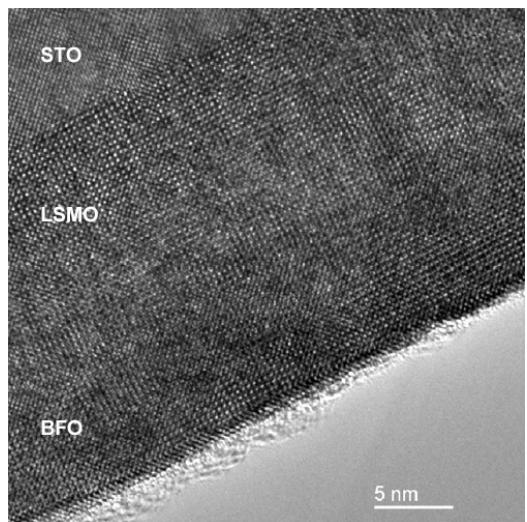


- **GMR** measured on top of BFO at RT
- **Exchange bias** has **shifted** the GMR curve

Tunnel junctions with BiFeO_3 barriers

Brought to you by PTP (www.ptp.phas.ubc.ca)

O_2 growth
pressure :
 $6 \cdot 10^{-3}$ mbar
 0.46 mbar



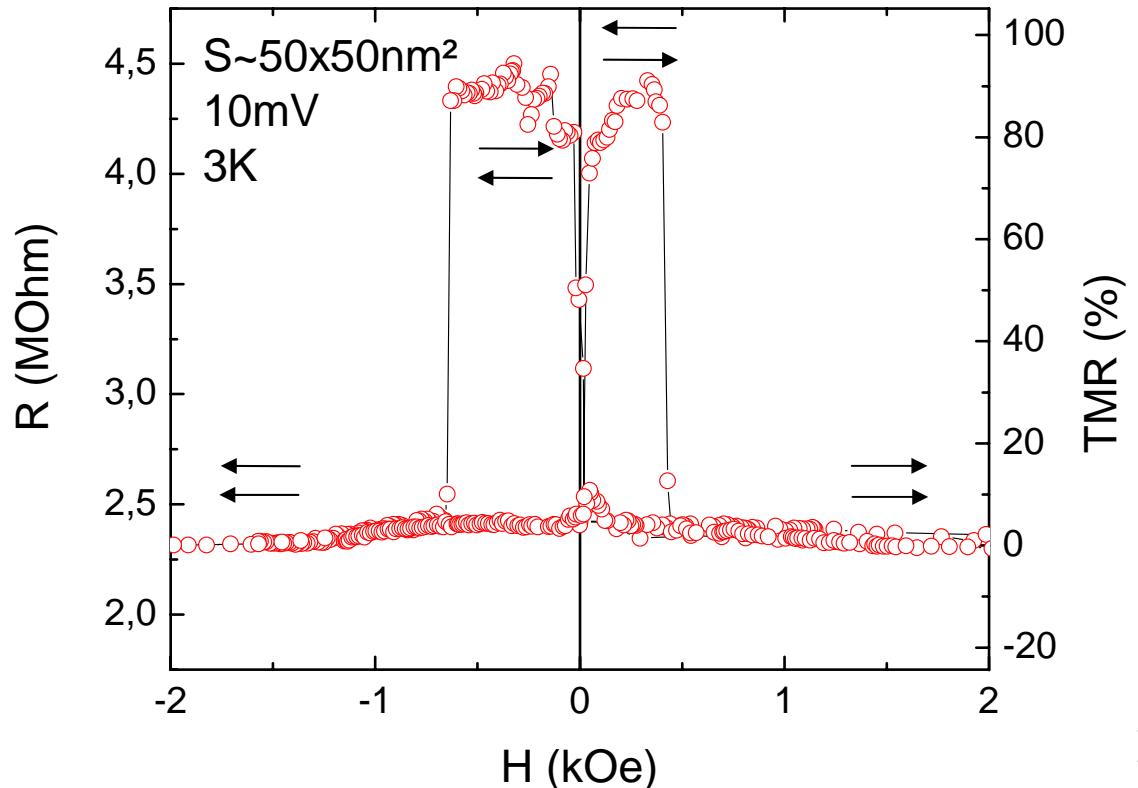
Appl. Phys. Lett., 89, 242114 (2006)

(exchange bias due to CoO, not to BFO)

Positive TMR up to ~30%

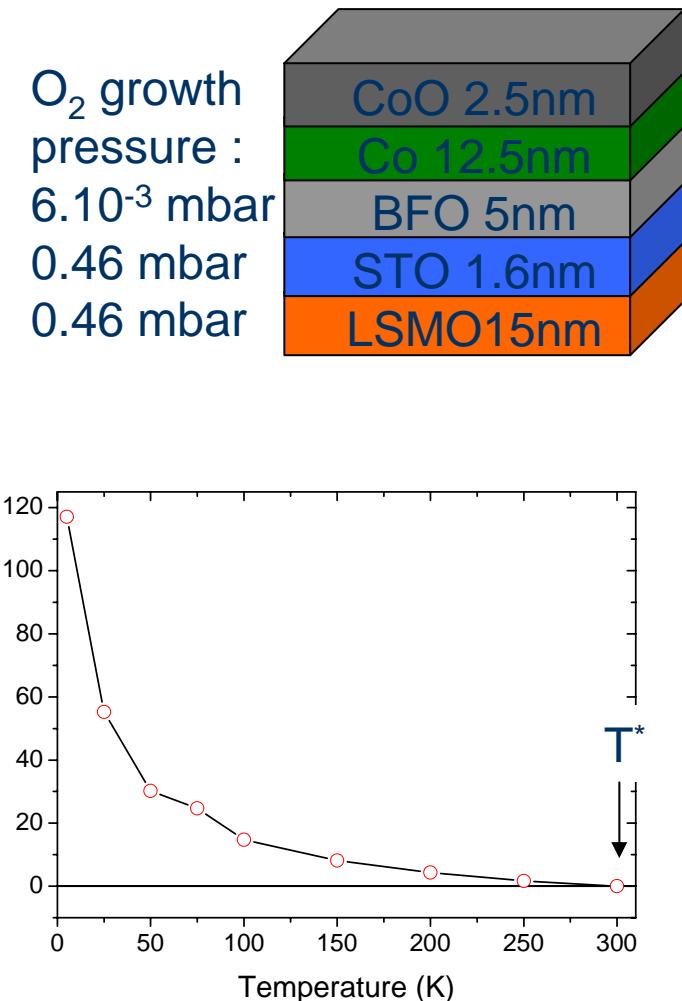
- rather large value for an AFM barrier
- positive spin polarization at Co/BFO interface

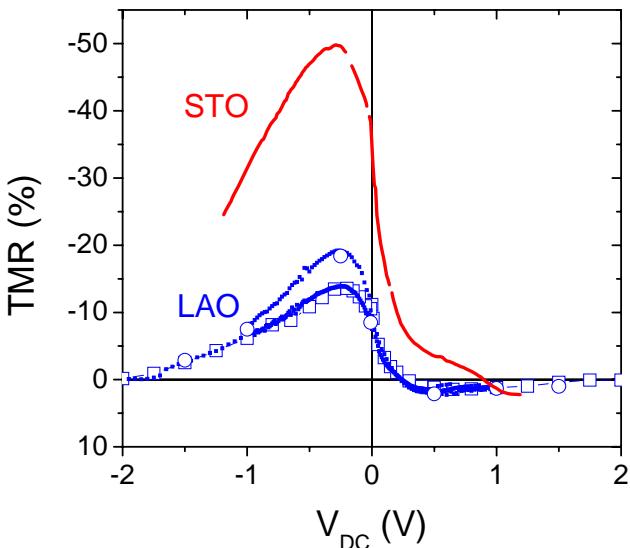
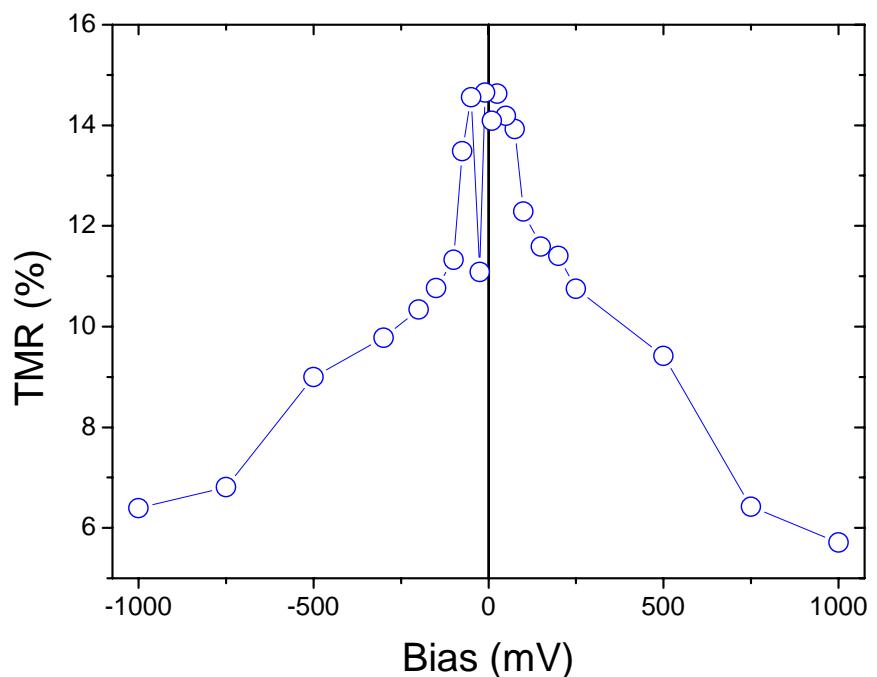
TMR vanishes around **200K**:
Local deoxygenation of LSMO



Positive TMR up to ~100%

- larger value than without STO
- polarization at Co/BFO interface still positive
- thanks to STO, better quality of the upper LSMO interface see [Appl. Phys. Lett., 82, 233 \(2003\)](#)



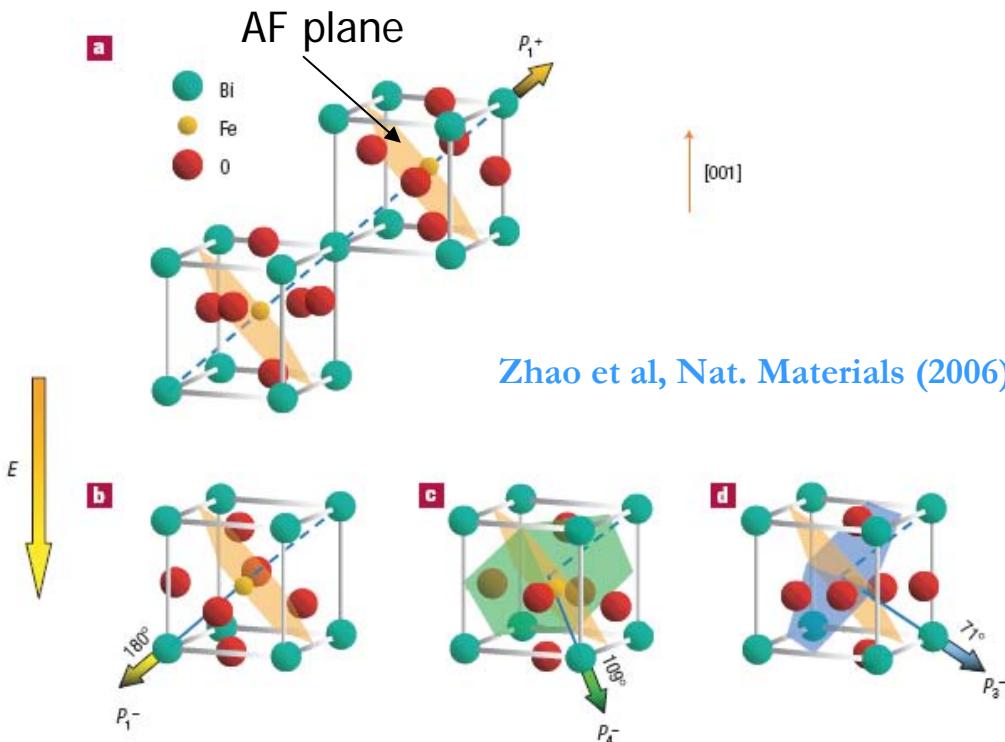
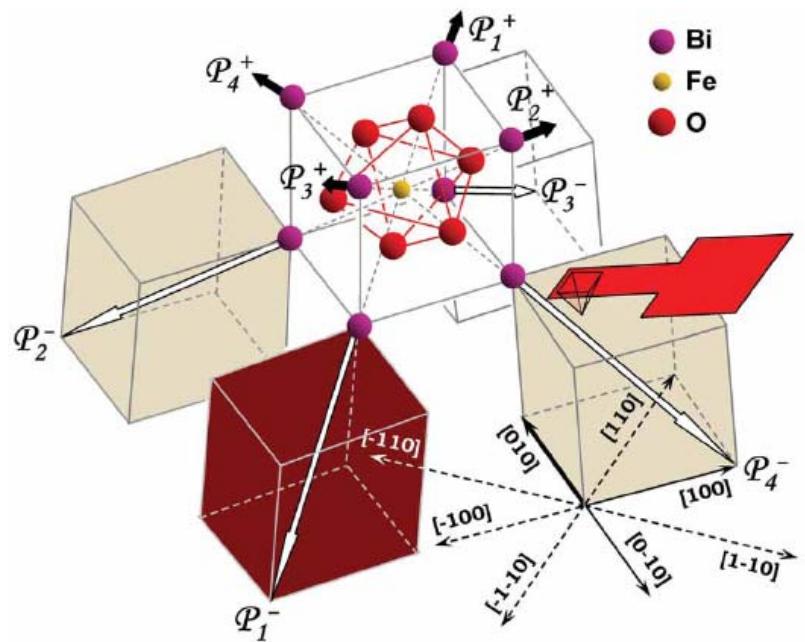


Bias dependence of TMR in
LSMO/STO/Co and LSMO/LAO/Co
junctions

Science 286, 507 (1999)
APL 87, 212501 (2005)

- TMR is positive and decreases symmetrically with bias voltage
- Behaviour is completely different from the case of LSMO/STO/Co and LSMO/LAO/Co junctions
- Possible reasons : oxygen vacancies at LSMO/BFO interface, symmetry filtering, etc

Domain structure and magnetoelectric switching



P along $<111>$ directions :
8 possible variants

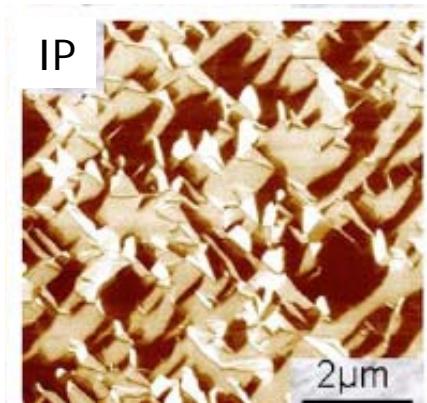
When an electric field is applied, P can be switched to different directions
Only some of them yield a rotation of the AF plane

Important to know and control the ferroelectric domain structure

Domain structure

Piezoresponse force microscopy (PFM)

(001) film
0.8° miscut



8 variants are present

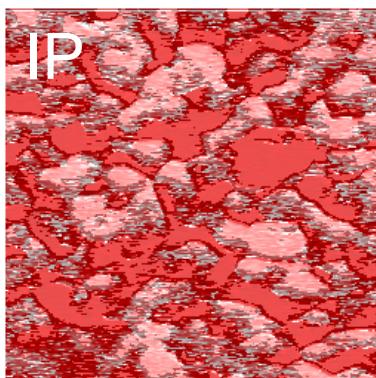
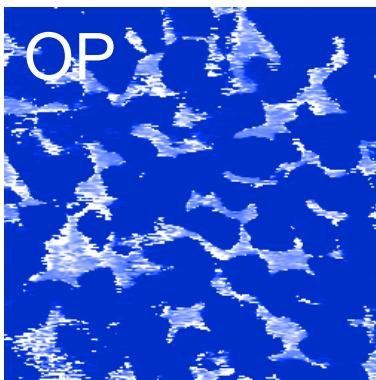
[Das et al., APL, 88, 242904 \(2006\)](#)

Domain structure can be controlled by playing with the substrate miscut angle and/or orientation.

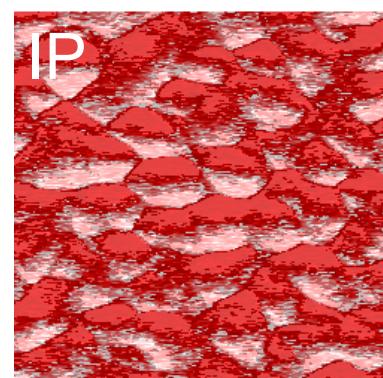
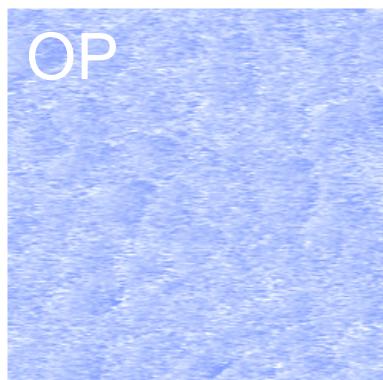
Ideally, one type of domain orientation would be required, with the appropriate polarization switching mechanism (109° ?)

(001) films

(111) films



8 variants are present



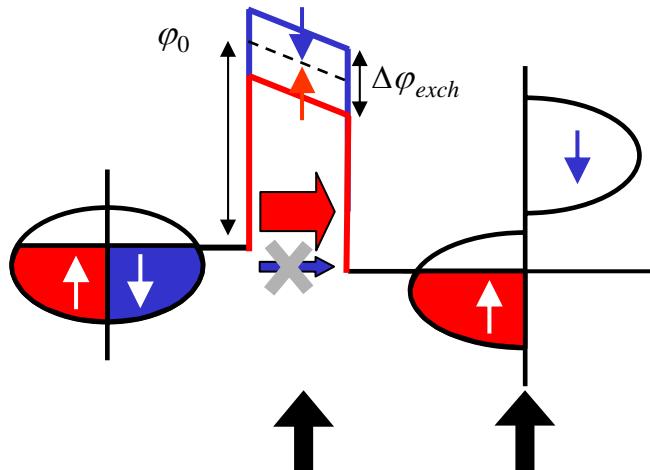
only 4 variants are present

(La,Bi)MnO₃

A ferromagnetic ferroelectric

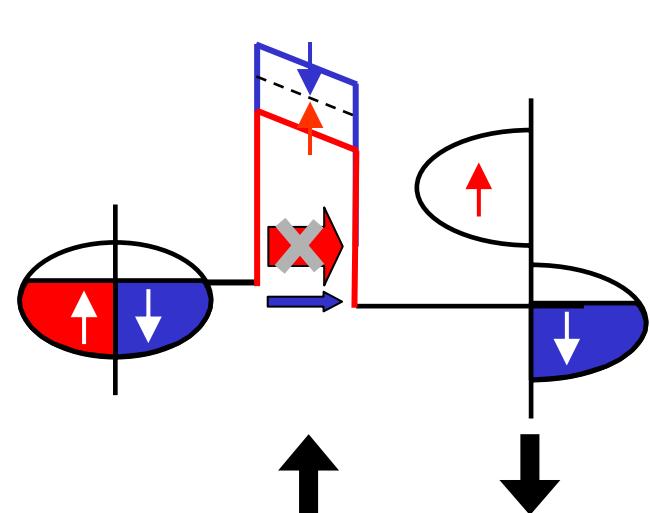
Ferromagnetic tunnel barriers: the spin-filter effect

Non Magn. Metal Ferro. Insulator Ferro. Half Metal



Parallel Config.

large I, low R



Antiparallel Config.

low I, large R

Since the tunnel transmission depends exponentially on the barrier height, a highly-spin polarized current is generated by the barrier.

$$J_{\uparrow\downarrow} \propto e^{-(\phi_0 \pm \frac{\Delta\phi_{exch}}{2})^{1/2}}$$

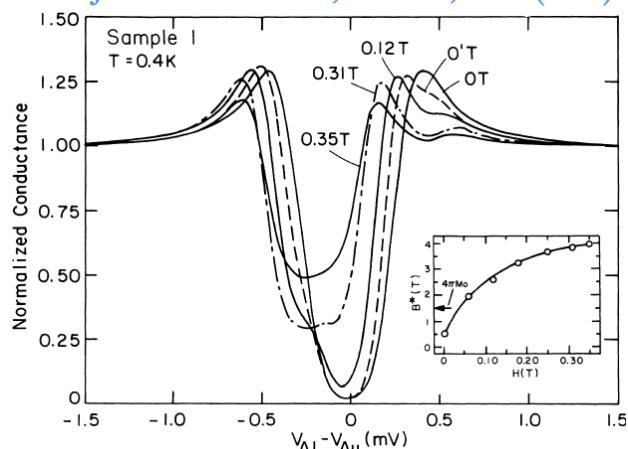
Expected TMR :

$$TMR = \frac{2P_E P_B}{1 - P_E P_B}$$

$$P_B = \frac{J_{\uparrow} - J_{\downarrow}}{J_{\uparrow} + J_{\downarrow}}$$

Spin-filters

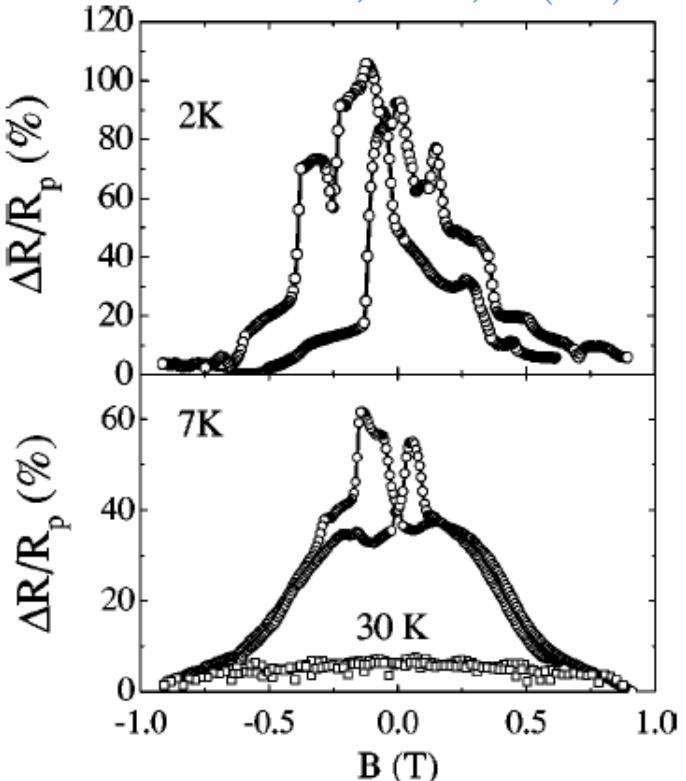
J.S. Moodera et al, PRB 42, 8235 (1990)



Au/EuS/Al spin-filter

Superconducting Al is used the spin-analyzer

P. LeClair et al, APL 80, 625 (2002)



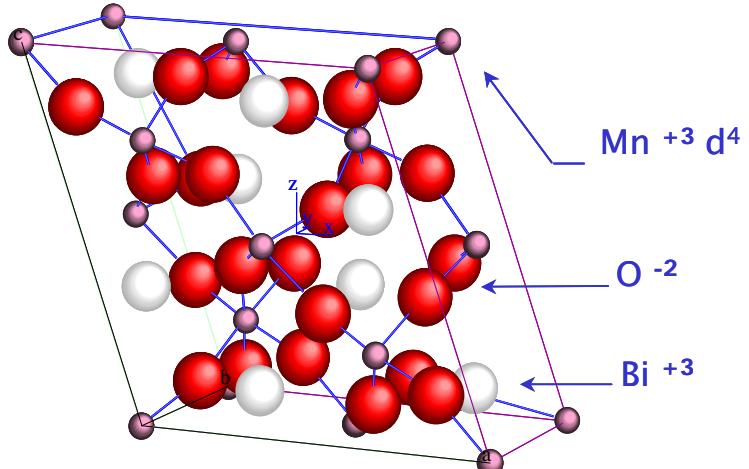
Au/EuS/Gd spin-filter

Ferromagnetic Gd is used as the spin-analyzer

- Early spin-filtering experiments focused on Eu chalcogenides
- Large spin-filtering efficiency (>90%) have been measured
- Limited by low T_c of EuX compounds

→ Ferromagnetic oxides

Crystal structure



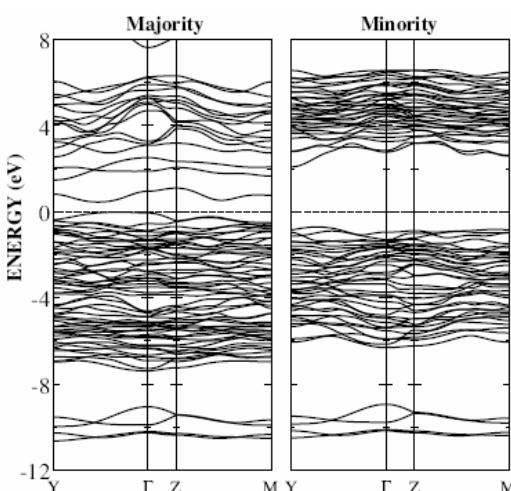
- Synthesis at high pressure
- Distorted perovskite structure (Monoclinic symmetry)
- Polar structure
- Ferroelectric (?) [Son et al, APL 84, 4971 \(2004\)](#)
- « Magnetodielectric effect » observed in bulk samples

[Kimura et al, PRB 67, 180401 \(2003\)](#)

Magnetic ordering

- Unusual orbital ordering resulting from Bi_{6s}-O_{2p} interaction [Moreira dos Santos et al, PRB 66 064425 \(2002\)](#)
- Ferromagnetic order
- T_C=105K, M_s=3.6 μ_B

Electronic structure



Ferromagnetic insulator with E_{g↑} = 0.5 eV and E_{g↓} = 2.6 eV

[Shishidou et al, JPCM 2005](#)

- Partial substitution of Bi by La : lower synthesis pressure, stabilization of perovskite phase [Troyanchuk et al, Low. Temp. Phys. 28, 569 \(2002\)](#).
- Growth of BiMnO₃ and La_{0.1}Bi_{0.9}MnO₃ thin films

Growth by PLD

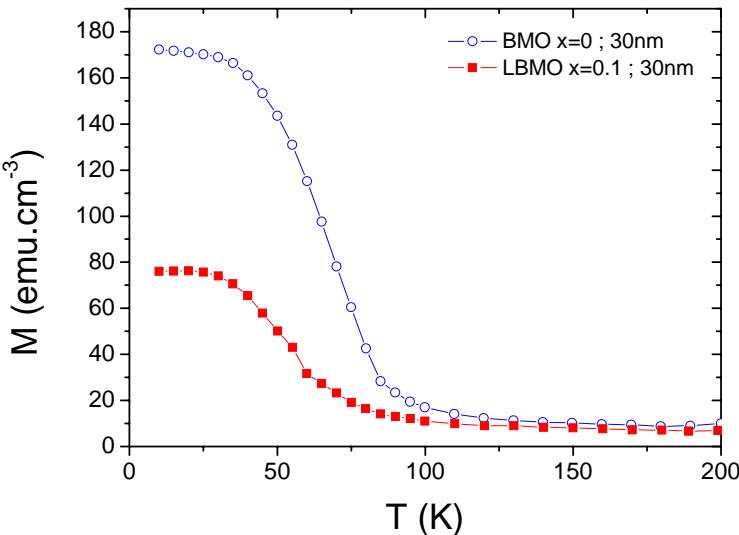
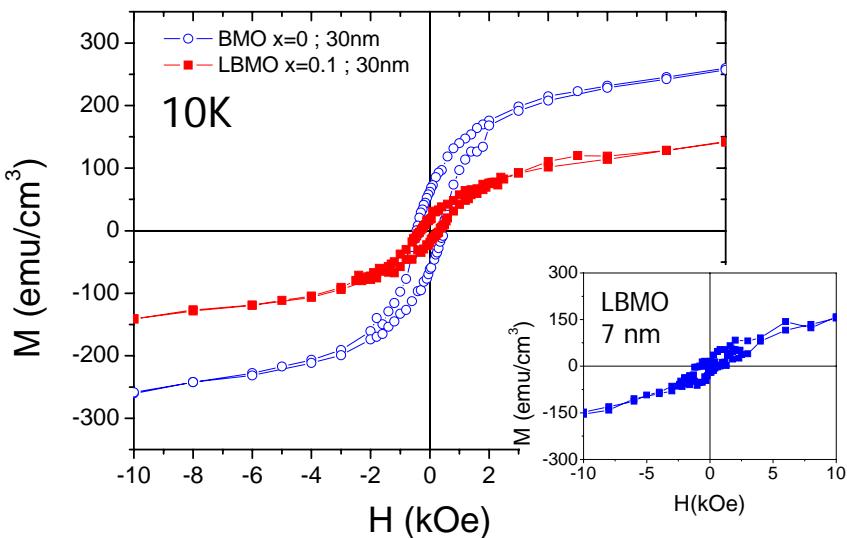
- KrF laser at 2 Hz
- O₂ Pressure: 10⁻¹ mbar.
- Fluence: 2 J/cm².
- T_{dep} from 575°C to 700°C

Single phase films only in a very narrow window
La-substitution helps to stabilize the perovskite phase

See for details : [PRB 75, 174417 \(2007\)](#)
[JAP 97 103909 \(2005\)](#)

- BMO films have a reduced moment compared to bulk
- T_C close to bulk (105K)
- Substitution by La further reduces max. moment and slightly decreases T_C (as in bulk)
- Low magnetic moment likely due to Bi vacancies
- Very thin films are also ferromagnetic

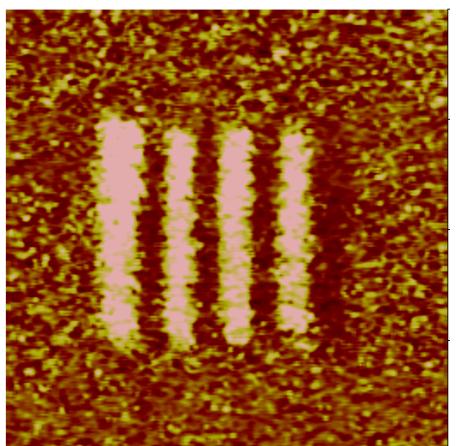
Magnetic properties



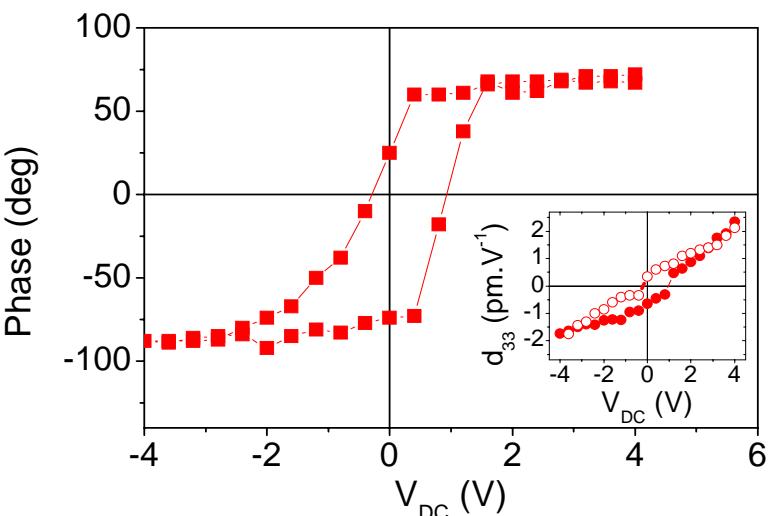
Ferroelectric properties

LBMO(30nm)/LSMO

PFM image after writing stripes at $\pm 4\text{V}$



180°
0°



Phase (deg)

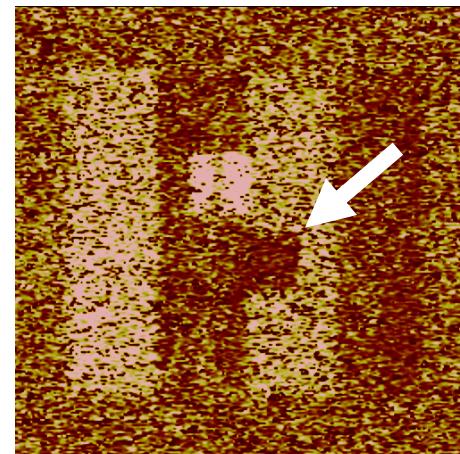
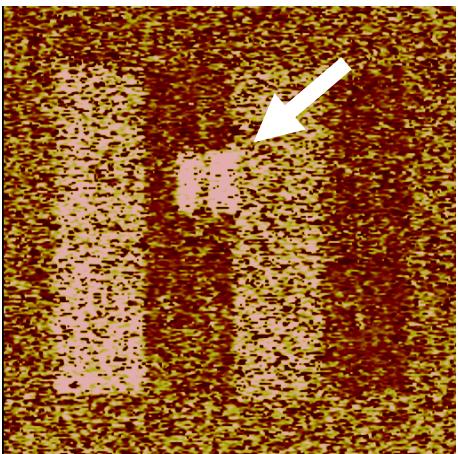
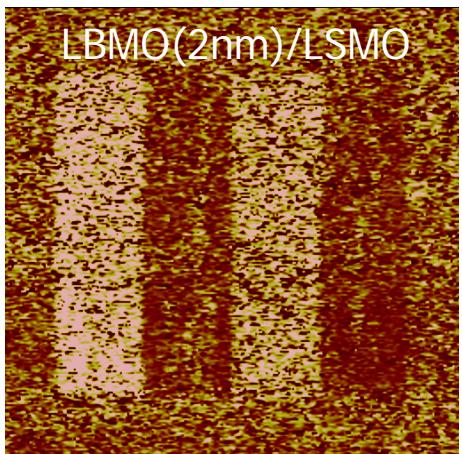
100
50
0
-50
-100

-4 -2 0 2 4 6
 V_{DC} (V)

d_{33} (pm.V⁻¹)
-2 -1 0 1 2
 V_{DC} (V)

LBMO(2nm)/LSMO

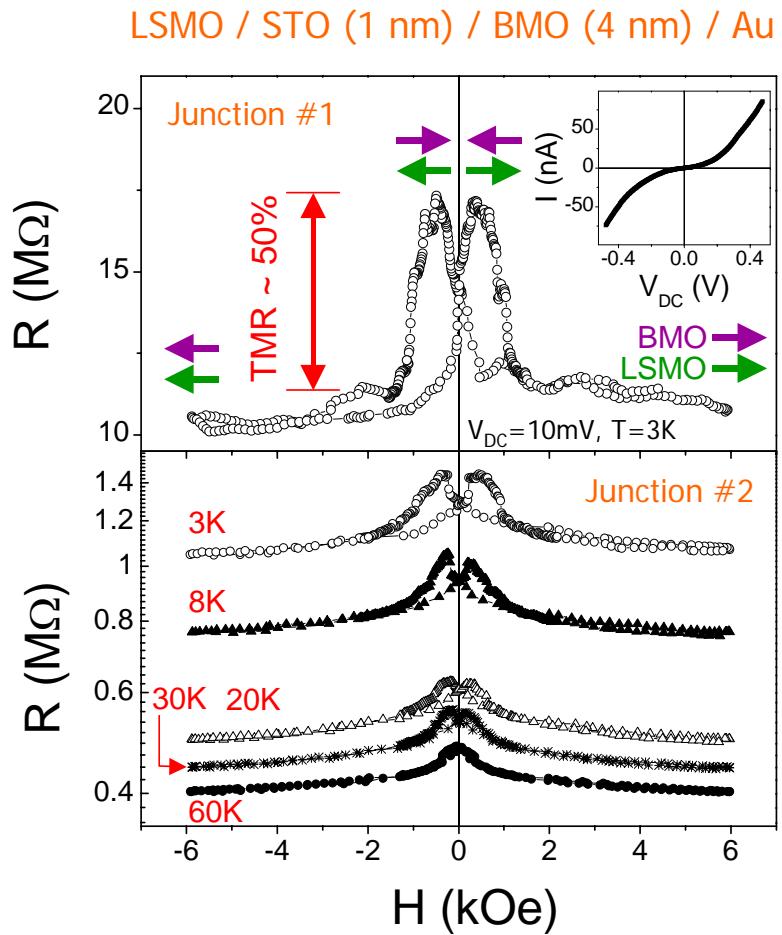
180°
0°



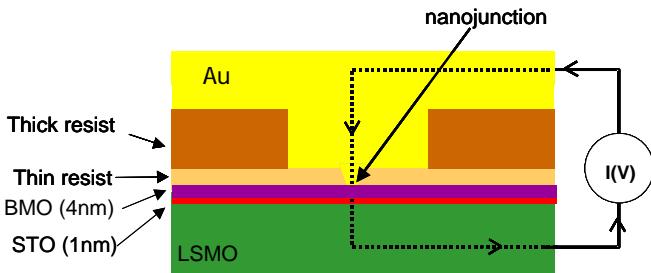
LBMO films as thin as 2 nm are still ferroelectric at room temperature

Nature Materials 6, 196 (2007)

Tunnel magnetoresistance



PRB 72 020406(R) (2005)



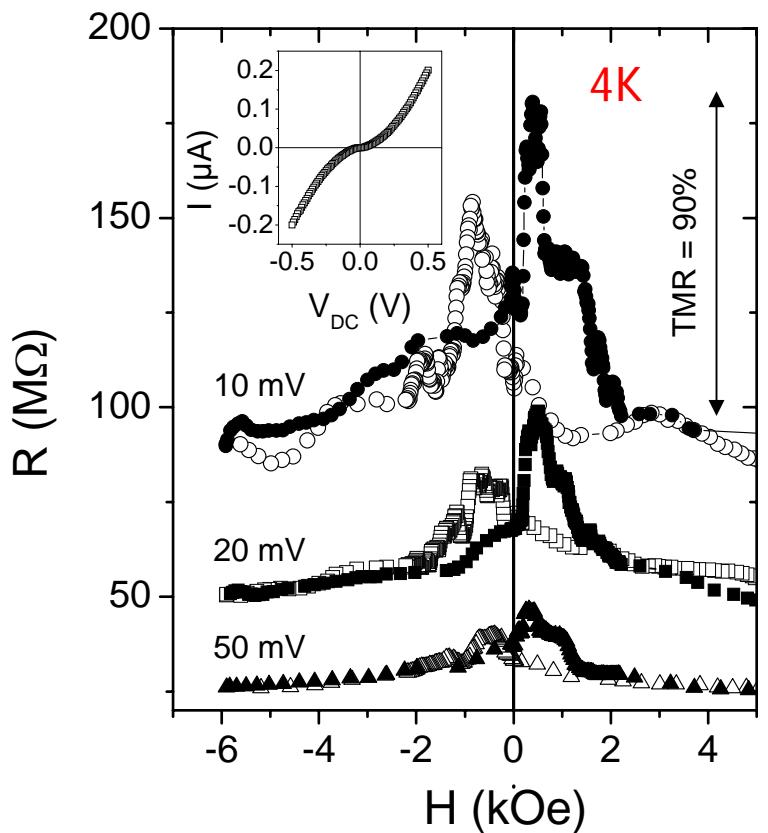
Junctions defined by nano-indentation lithography

Nanoletters 3, 1599 (2003)

- ➊ Large TMR at low temperature
- ➋ spin-filtering by the BMO layer
- ➌ Polarization induced by BMO : **22%**
- ➍ Spin-filter effect vanishes at $T < T_c$

Tunnel magnetoresistance

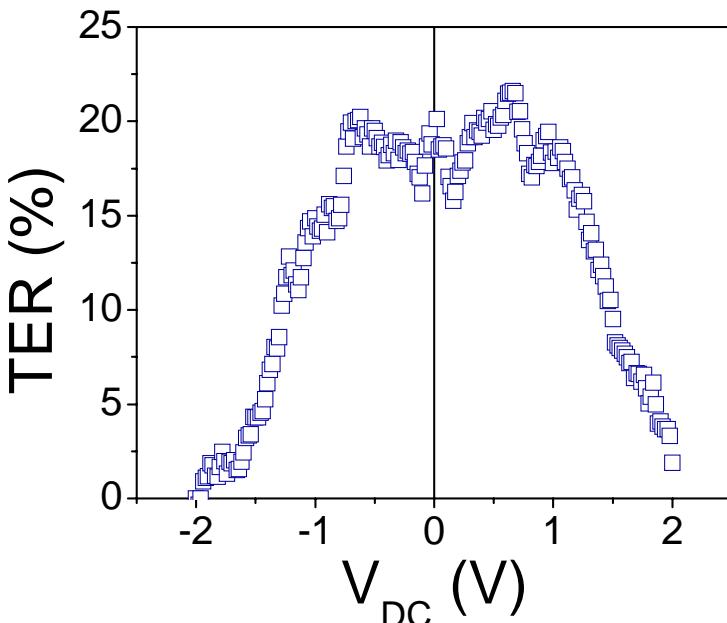
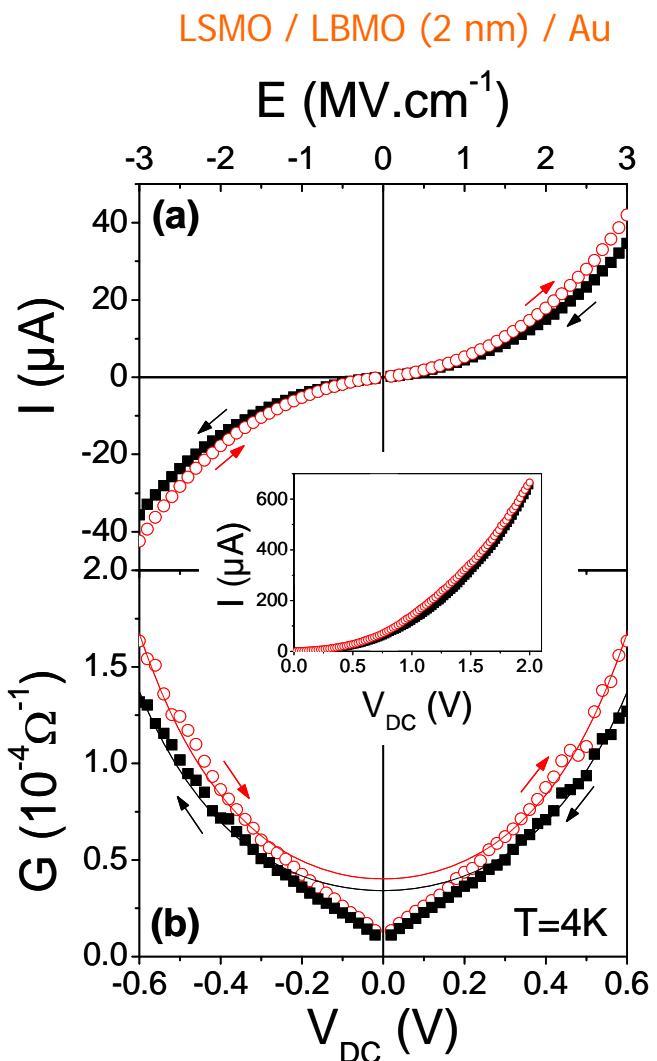
LSMO / STO (1 nm) / LBMO (4 nm) / Au



- Polarization induced by LBMO : 35%
- TMR decreases rapidly with bias
Magnons ? Defects ?

JAP 99, 08E504 (2006)

Tunnel electroresistance

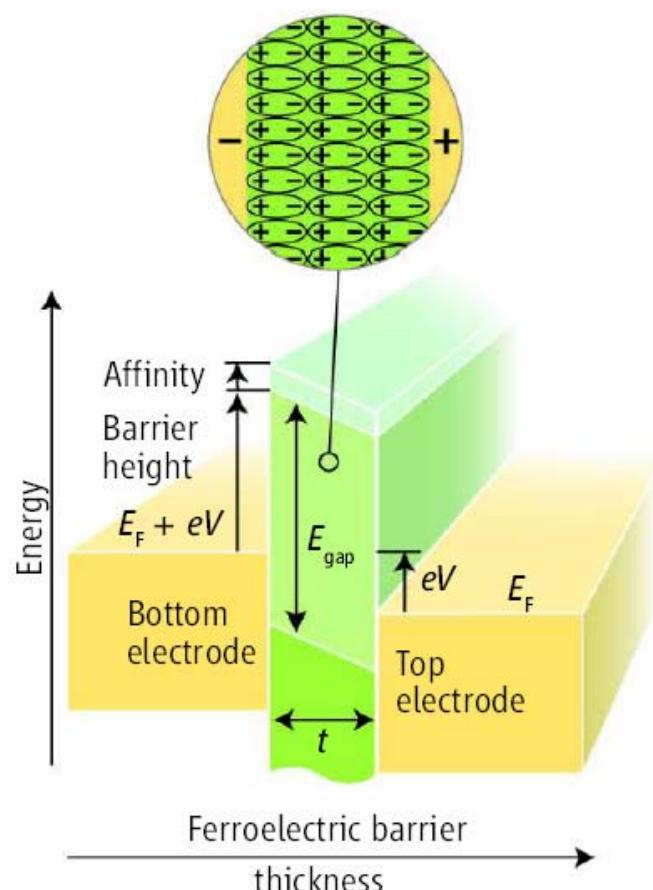


- A ~20% electroresistance effect is observed
- The shape of the $G(V)$ curve suggests direct tunneling

What is the origin of the TER ?

Nature Materials 6, 196 (2007)

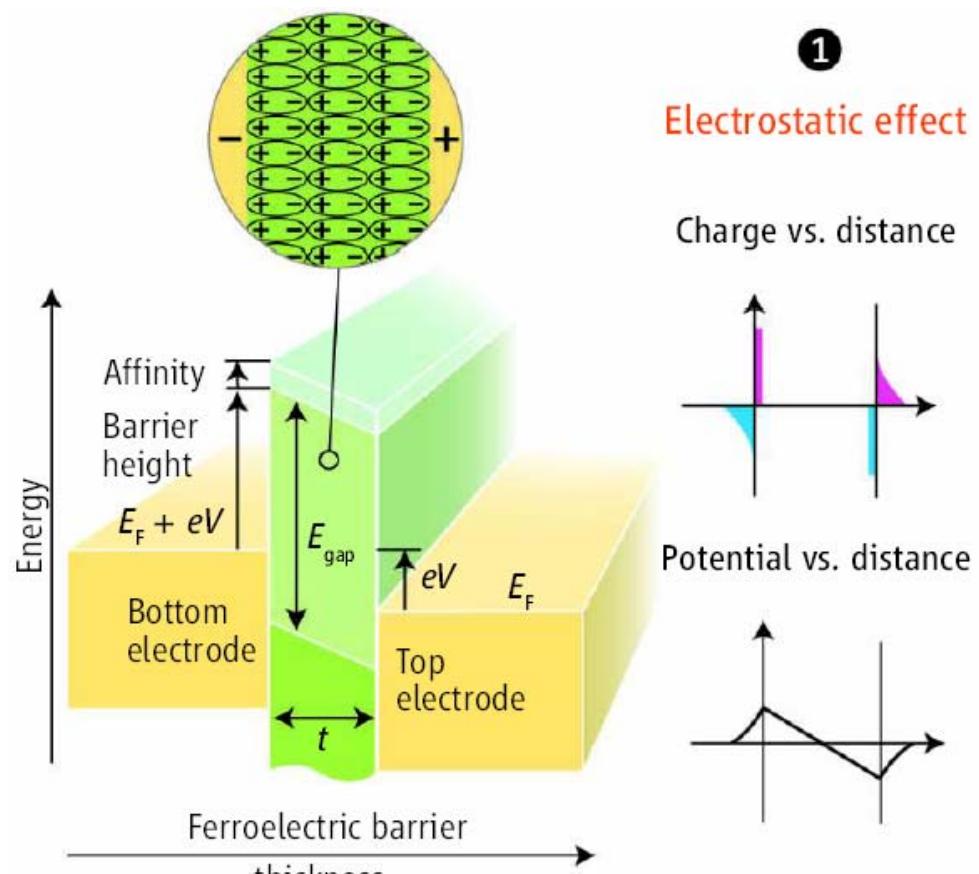
Tunneling through a ferroelectric tunnel barrier



From Tsymbal and Kohlstedt
Science 313, 181 (2006)

Several mechanisms leading to current modulation upon polarization reversal

Tunneling through a ferroelectric tunnel barrier

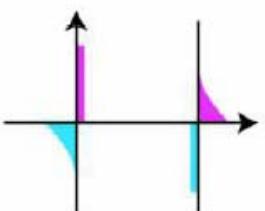


From Tsymbal and Kohlstedt
Science 313, 181 (2006)

①

Electrostatic effect

Charge vs. distance

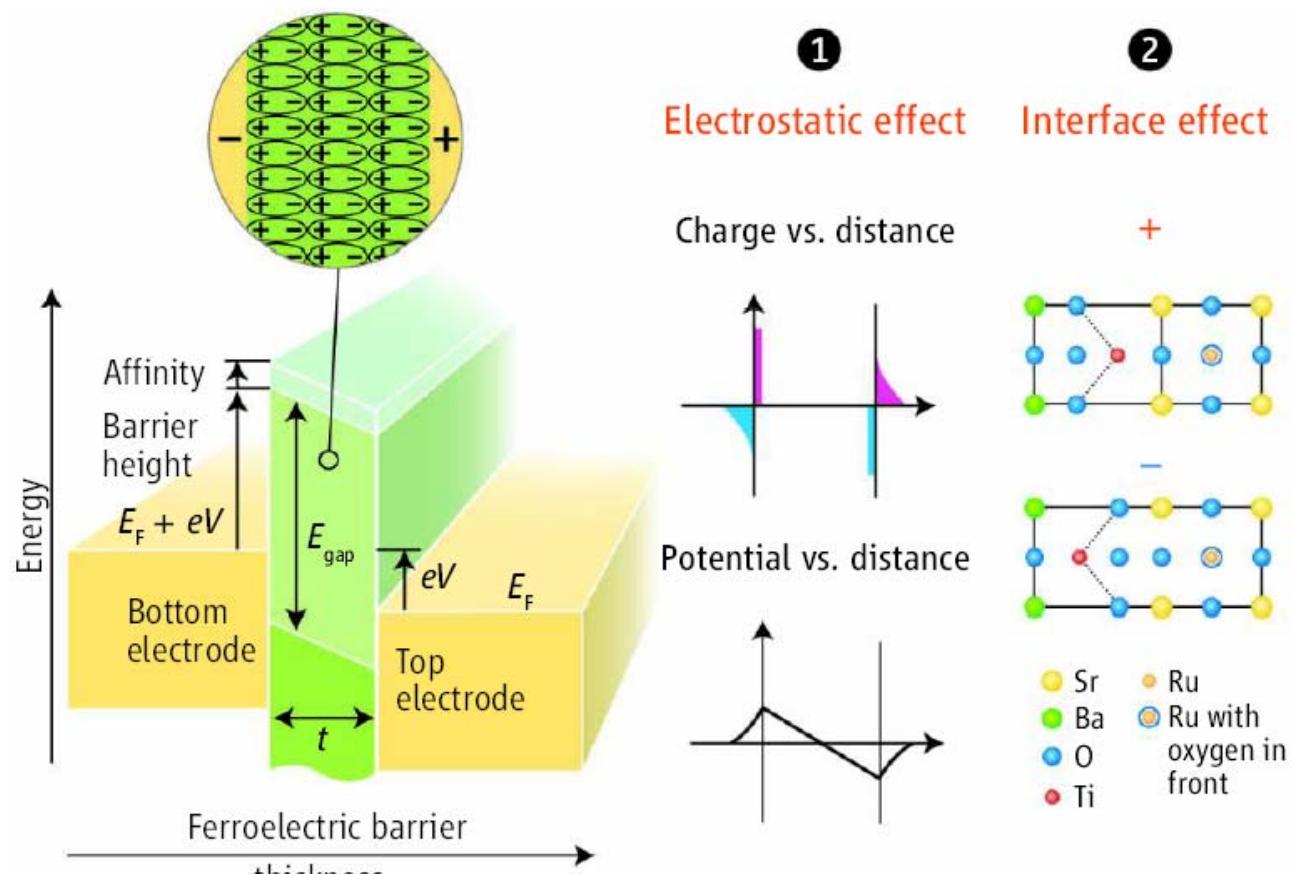


Potential vs. distance



Several mechanisms leading to current modulation upon polarization reversal

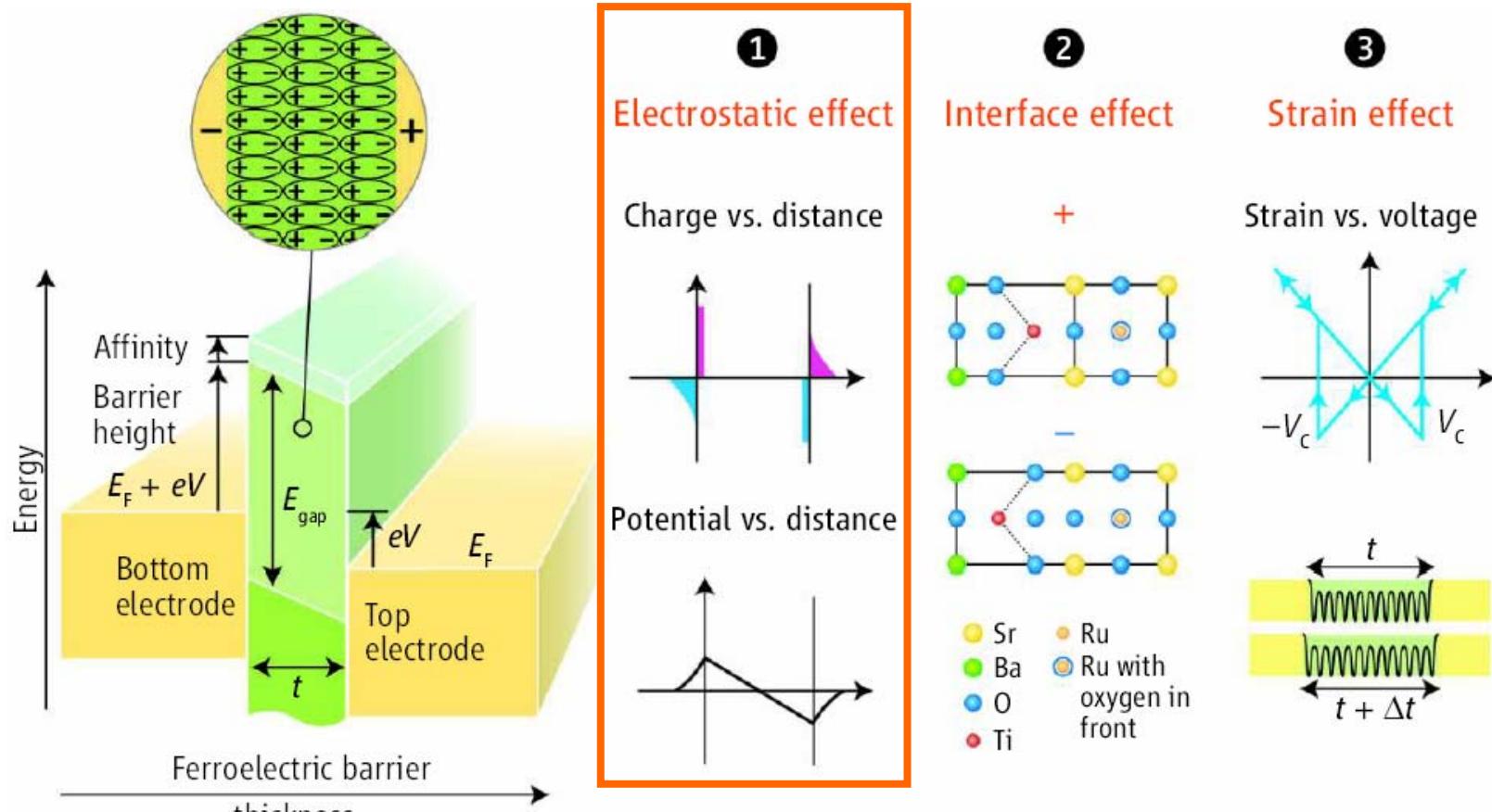
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Several mechanisms leading to current modulation upon polarization reversal

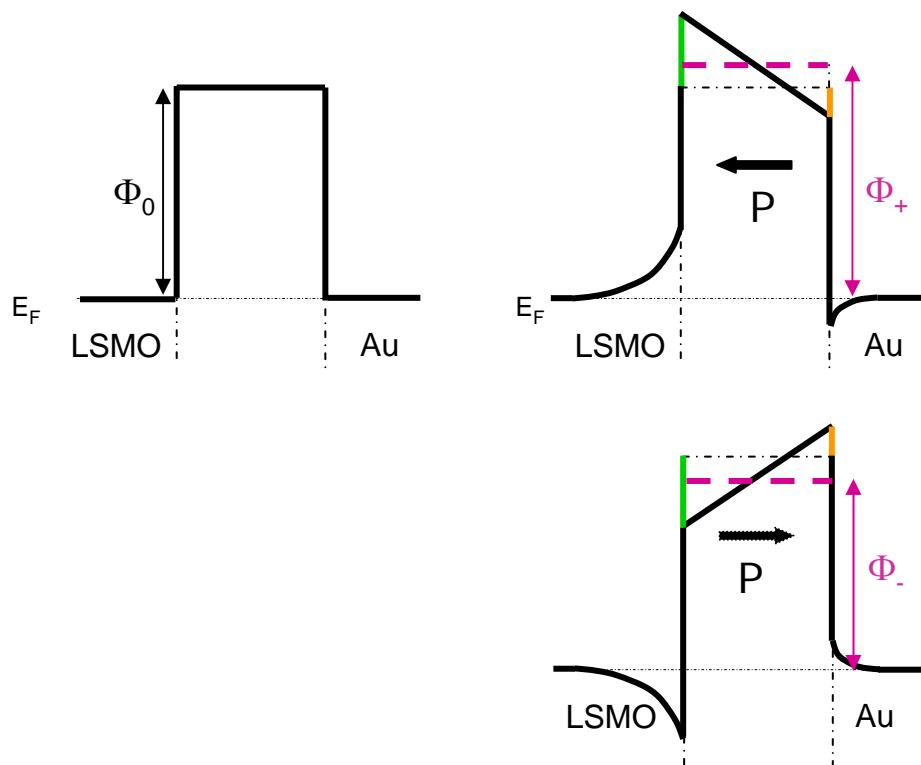
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Tunneling through a ferroelectric tunnel barrier



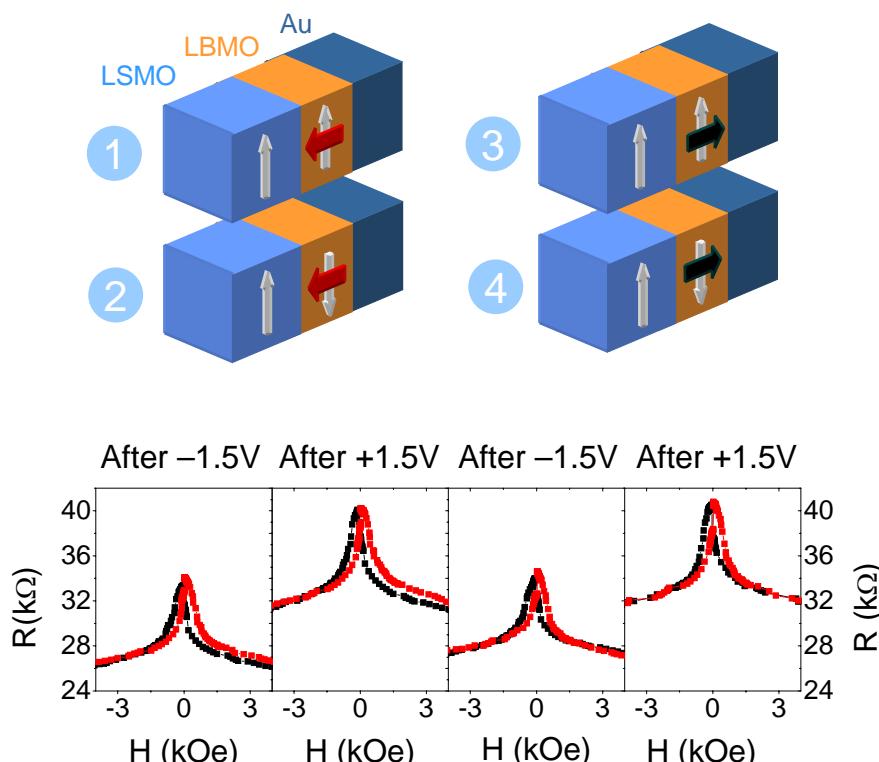
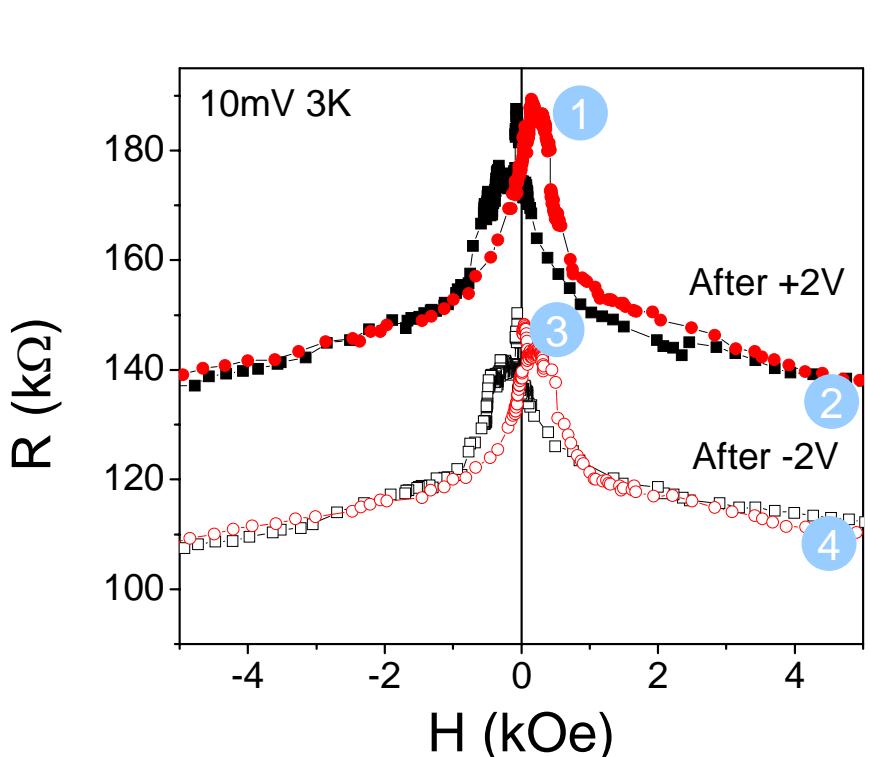
Large average barrier height
 → low tunnel current

Low average barrier height
 → large tunnel current

**Hysteretic I(V) curves are expected
 An electroresistance effect should occur upon polarization reversal**

See details in [Zhuravlev et al, PRL 94, 246802 \(2005\)](#)

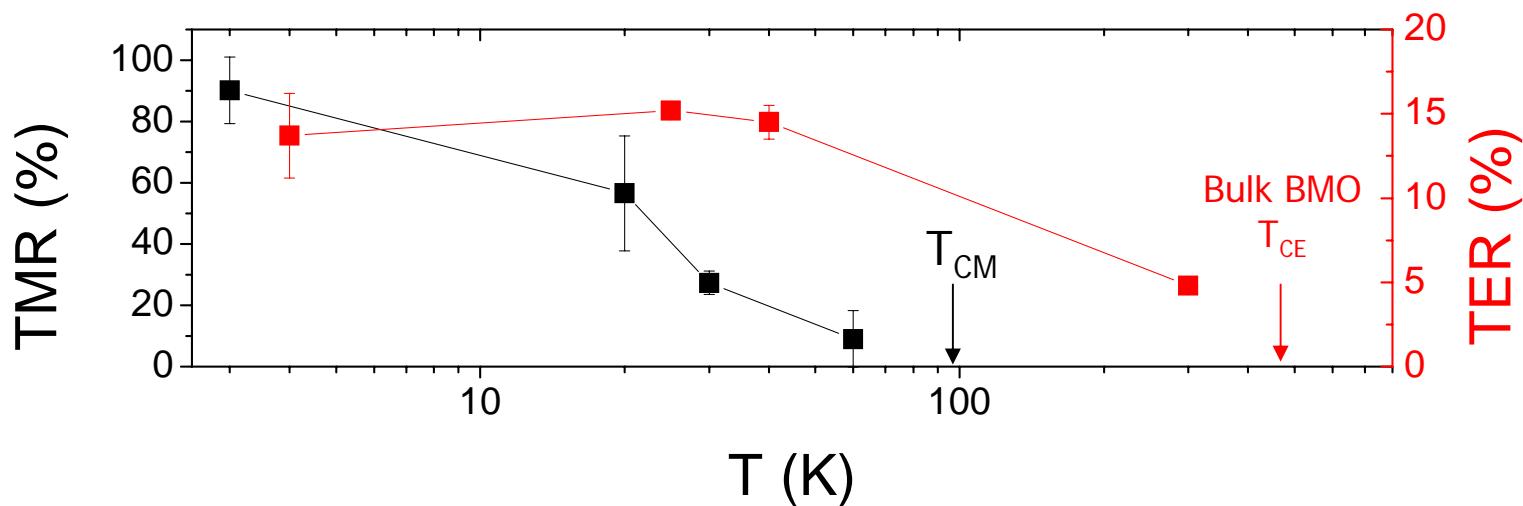
Combination of tunnel electroresistance and magnetoresistance



- ➁ Demonstration of a 4-resistance state system with a multiferroic barrier
- ➂ Effect was observed on several junctions and is reproducible

Nature Materials 6, 196 (2007)

Temperature dependence of TER and TMR



Temperature dependence of TER and TMR consistent with an origin related to ferroelectricity and ferromagnetism, respectively

BiFeO₃

Single phase films can be grown in a narrow range of P and T
They are G-type antiferromagnetic and ferroelectric at RT
LSMO/BFO/Co junctions show a large positive TMR
BFO induced exchange bias on CoFeB

(La,Bi)MnO₃ based junctions

La_{0.1}Bi_{0.9}MnO₃ films are ferromagnetic and ferroelectric
LBMO junctions show TMR and electroresistance (4 resistance states)
We suggest electroresistance effect due to ferroelectric switching in the barrier
Control of the magnetic state by electric field ?

Perspectives :

Magnetization manipulation with an electric field at RT
New materials, ideally ferromagnetic multiferroics with high T_Cs, M_S
and P are required

See review on Oxide Spintronics by M. Bibes and A. Barthélémy, IEEE Trans. Electron Devices 54, 1003 (2007)

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ESF Thiox program, French ANR program, Spanish Ministry for Research and EGIDE