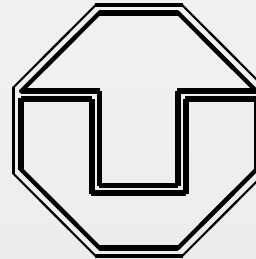


Sweep from Superfluid to Mott phase in the Bose-Hubbard model

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Motivation

- Equilibrium properties well understood for many systems, e.g., near a phase transition
- Dynamical (time-dependent) phase transitions
 - response time typically diverges
 - non-equilibrium properties
- Analogy to expanding universe
 - effective horizon (universal behaviour)
 - loss of causal contact
- Amplification of quantum fluctuations
 - seeds for pattern formation etc.
- Physical example: Bose-Hubbard model
- Relation to real cosmic inflation?

Bose-Hubbard Model

$$\hat{H} = J(t) \sum_{\alpha\beta} M_{\alpha\beta} \hat{a}_{\alpha}^{\dagger} \hat{a}_{\beta} + \frac{U}{2} \sum_{\alpha} (\hat{a}_{\alpha}^{\dagger})^2 \hat{a}_{\alpha}^2$$

Interaction U , tunnelling rate $J(t)$, lattice matrix $M_{\alpha\beta}$

Large integer filling $n = \langle \hat{n}_{\alpha} \rangle = \langle \hat{a}_{\alpha}^{\dagger} \hat{a}_{\alpha} \rangle \gg 1$

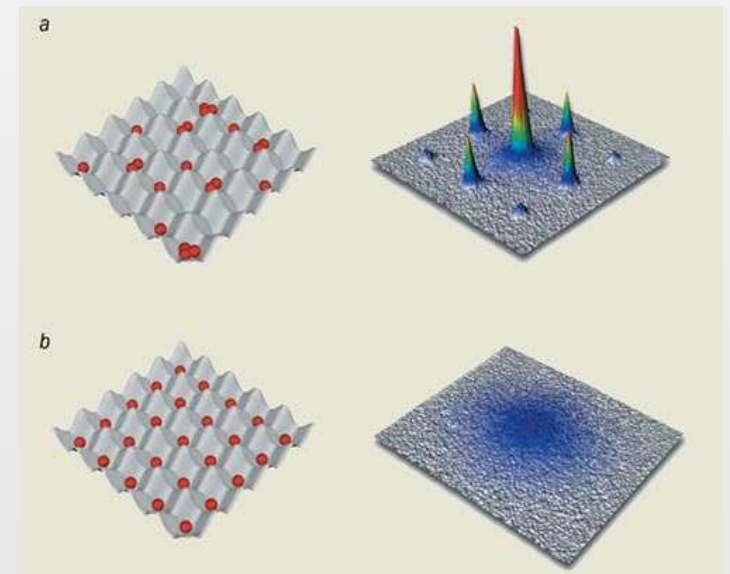
Superfluid \rightarrow **Mott phase transition** at $J_c = \mathcal{O}(U/n)$

$J \gg U/n \rightarrow$ superfluid phase

$$|\Psi\rangle_{\text{sf}} \propto \left(\sum_{\alpha} \hat{a}_{\alpha}^{\dagger} \right)^N |0\rangle$$

$J \ll U/n \rightarrow$ Mott insulator

$$|\Psi\rangle_{\text{Mott}} \propto \bigotimes_{\alpha} (\hat{a}_{\alpha}^{\dagger})^n |0\rangle$$



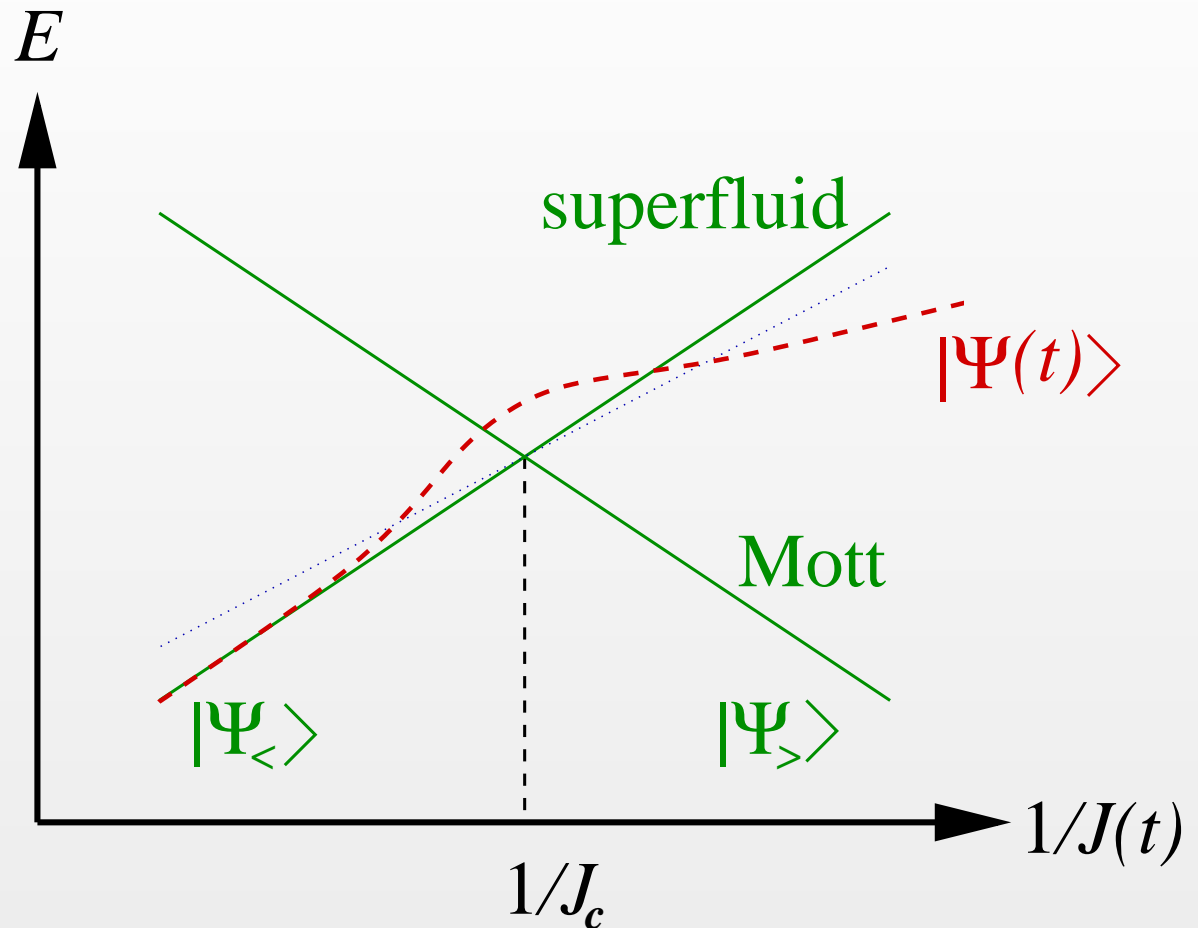
Dynamical Phase Transition

Zero temperature
 $T = 0$

External parameter
 $J = J(t)$

Two competing
ground states
 $|\Psi_{<}\rangle$ and $|\Psi_{>}\rangle$

Actual quantum
state $|\Psi(t)\rangle$



Sweeping through phase transition (critical point J_c)

→ non-equilibrium dynamics $J(t)$

Number Fluctuations

Normal-mode expansion with label κ , eigenvalues λ_κ

$$\left(\frac{\partial}{\partial t} \frac{1}{J(t)} \frac{\partial}{\partial t} + 8\lambda_\kappa [Un + 2J(t)\lambda_\kappa] \right) \delta\hat{n}_\kappa = 0$$

Exponential sweep of tunnelling rate (experiments)

$$J(t) = J_0 \exp\{-\gamma t\}$$

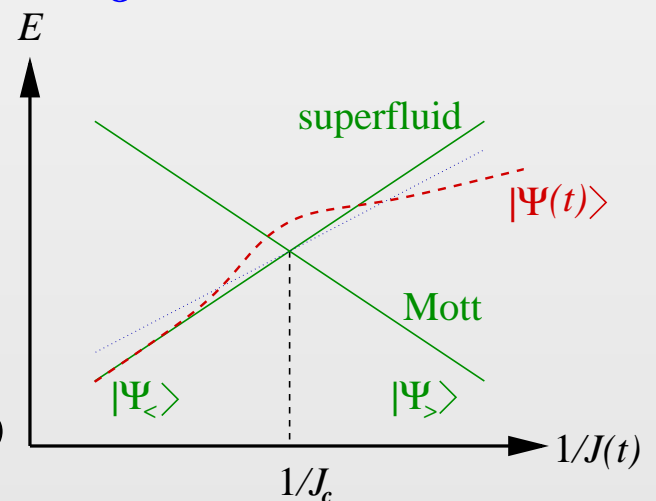
Scaling solution with $\tau_\kappa = -4\lambda_\kappa J(t)/\gamma$

$$\left(\frac{\partial^2}{\partial \tau_\kappa^2} + \left[1 - \frac{2}{\tau_\kappa} \frac{Un}{\gamma} \right] \right) \delta\hat{n}_\kappa = 0$$

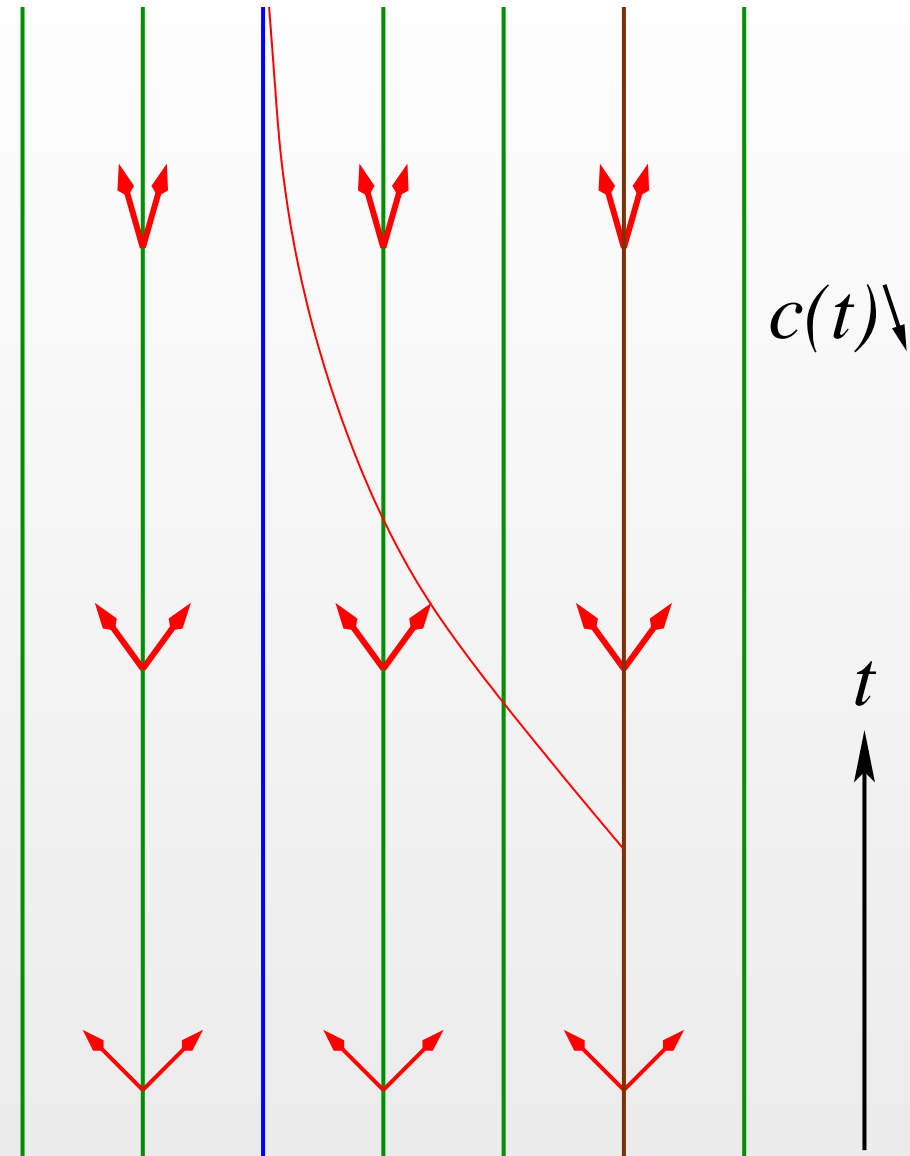
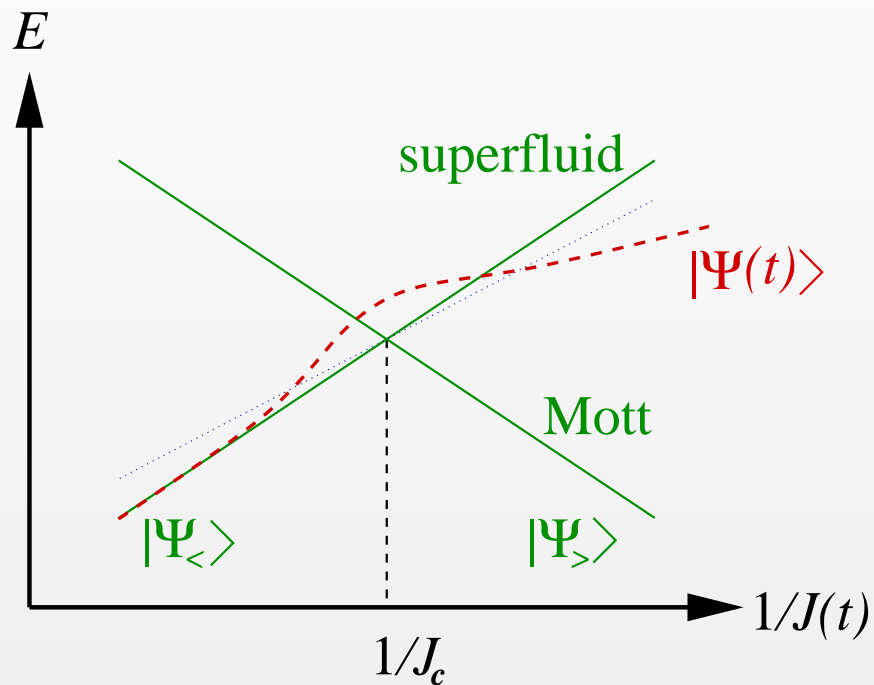
Adiabaticity parameter

$$\nu = \frac{Un}{\gamma} = \frac{\mu}{\gamma}$$

Fast $\nu \ll 1$ vs slow $\nu \gg 1$ sweep



“Cosmic” Horizon



Analogue of cosmic horizon

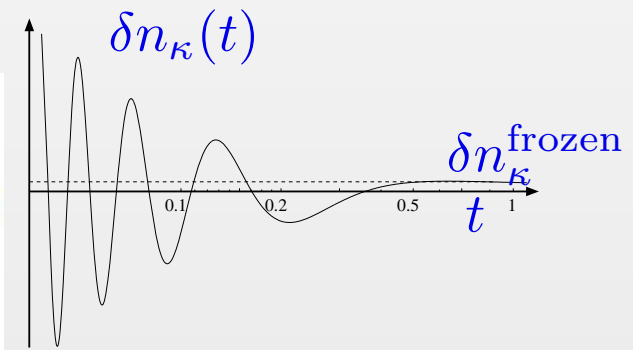
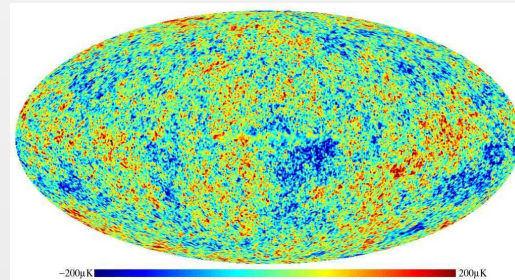
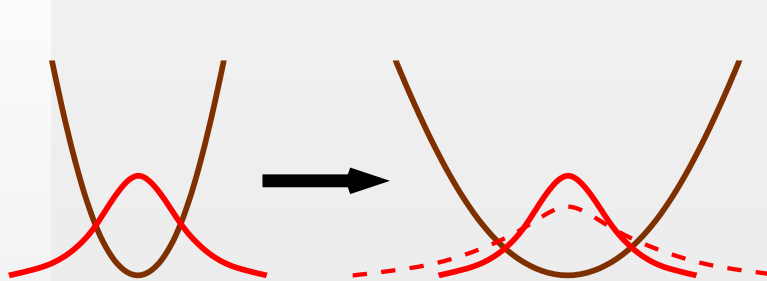
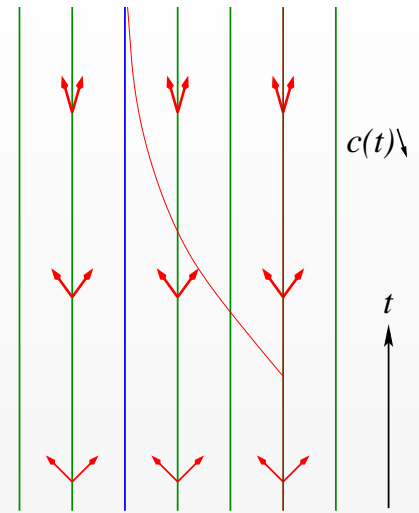
$$\Delta_{\text{horizon}}(t) = \int_t^{\infty} dt' \sqrt{J(t')Un}$$

Quantum Fluctuations

Size of “cosmic” horizon always decreases

$$\frac{d}{dt} \Delta_{\text{horizon}}(t) < 0$$

Oscillation $\lambda \ll \Delta r(t) \rightarrow$ horizon crossing \rightarrow
 \rightarrow freezing $\lambda \gg \Delta r(t)$ and squeezing



Amplification of quantum fluctuations $\hbar\omega/2$

Analogous to early universe (WMAP)

Frozen Number Fluctuations

$$\langle \delta \hat{n}_\kappa^2 \rangle = n \frac{1 - e^{-2\pi\nu}}{2\pi\nu} + \mathcal{O}(t\lambda_\kappa e^{-\gamma t})$$

Off-site number correlations decay exponentially

$$\langle \hat{n}_\alpha \hat{n}_\beta \rangle - \langle \hat{n}_\alpha \rangle \langle \hat{n}_\beta \rangle = \langle \delta \hat{n}_\alpha \delta \hat{n}_\beta \rangle = \mathcal{O}(\gamma t e^{-\gamma t})$$

On-site number variations

$$\Delta^2(n_\alpha) = \langle \hat{n}_\alpha^2 \rangle - \langle \hat{n}_\alpha \rangle^2 = \langle \delta \hat{n}_\alpha^2 \rangle = n \frac{1 - e^{-2\pi\nu}}{2\pi\nu}$$

Fast $\nu \ll 1$ sweep:

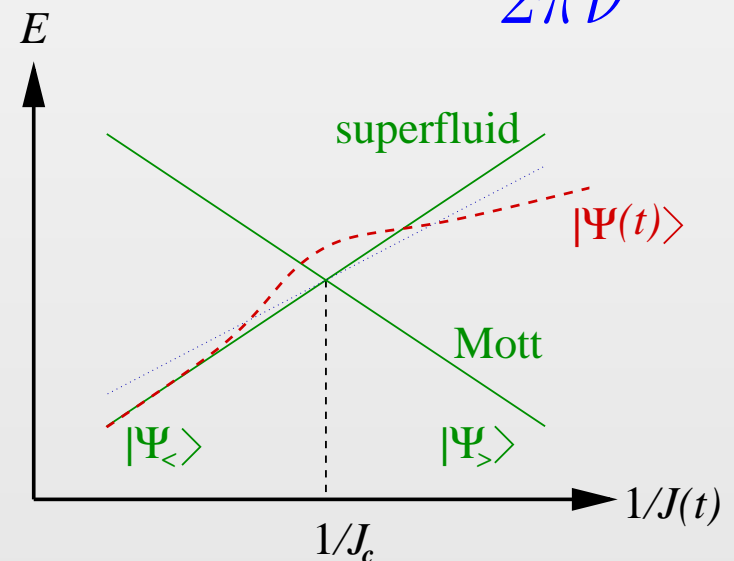
$$\Delta^2(n_\alpha) \rightarrow n$$

(Poissonian \rightarrow superfluid)

Slow $\nu \gg 1$ sweep:

$$\Delta^2(n_\alpha) \rightarrow 0$$

(\rightarrow Mott)



Off-diagonal long-range order

Phase fluctuations grow (instability)

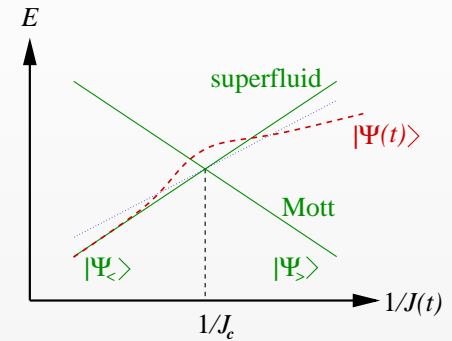
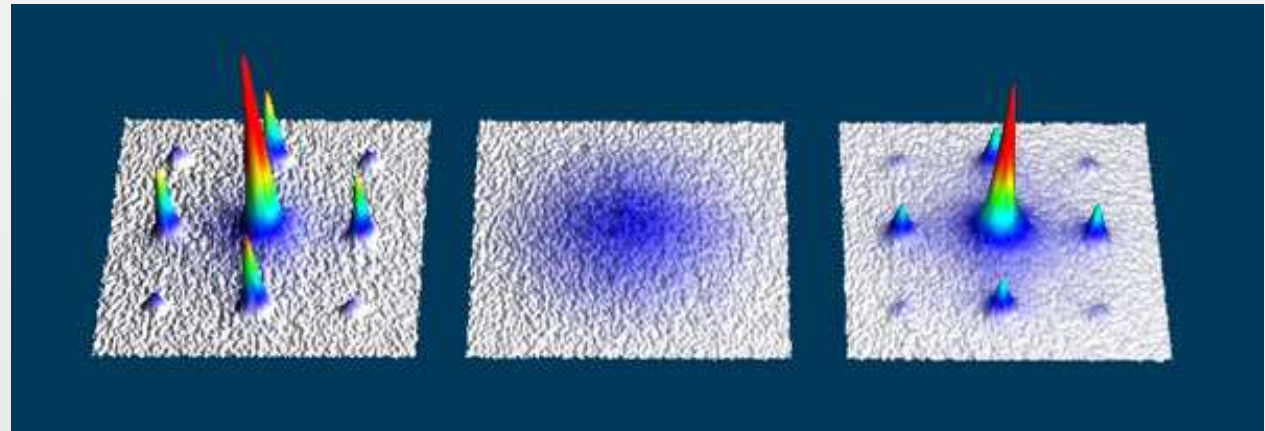
$$\langle \delta \hat{\phi}_{\kappa}^2 \rangle = \nu \frac{1 - e^{-2\pi\nu}}{2\pi n} \gamma^2 t^2 + \mathcal{O}(\gamma t \ln \lambda_{\kappa})$$

Decay of off-diagonal long-range order

$$\langle \hat{a}_{\alpha}^{\dagger}(t) \hat{a}_{\beta}(t) \rangle \approx n \exp\{-U^2 t^2 \Delta^2(n_{\alpha})\}$$

Loss of phase coherence due to horizon

→ peak at
 $k = 0$
 decreases
 (Greiner
et al)



Note: revival for $Ut \in 2\pi\mathbb{N}$ (similar to spin echo)

Decay of Superfluid

Superfluid fraction defined via $\mathbf{j} = \rho_{\text{sf}} \nabla \Phi$ is determined by $\langle \hat{a}_\ell^\dagger \hat{a}_{\ell+1} \rangle$ and thus also decreases

$$\frac{n_{\text{sf}}}{n} \approx \exp \left\{ -U^2 t^2 n \frac{1 - e^{-2\pi\nu}}{2\pi\nu} \right\}$$

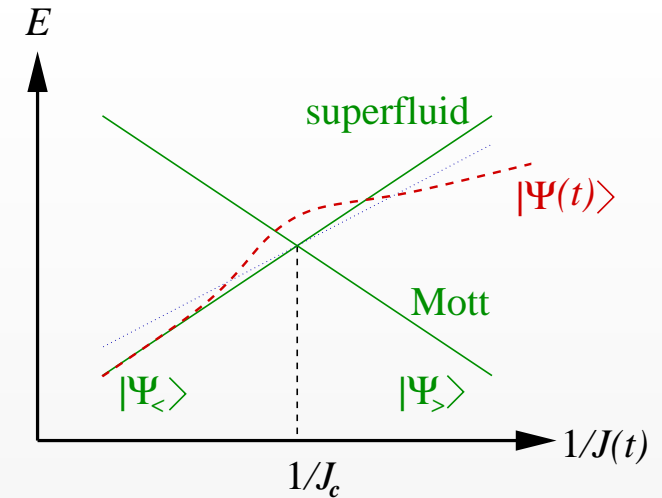
Rapid sweep $\nu \ll 1 \rightarrow$ decay with $\exp\{-nU^2 t^2\}$
(independent of γ)

Adiabatic sweep $\nu \gg 1 \rightarrow$ decay of superfluid fraction much slower $\exp\{-nU^2 t^2 / (2\pi\nu)\}$

R. S., M. Uhlmann, Y. Xu and U. R. Fischer,

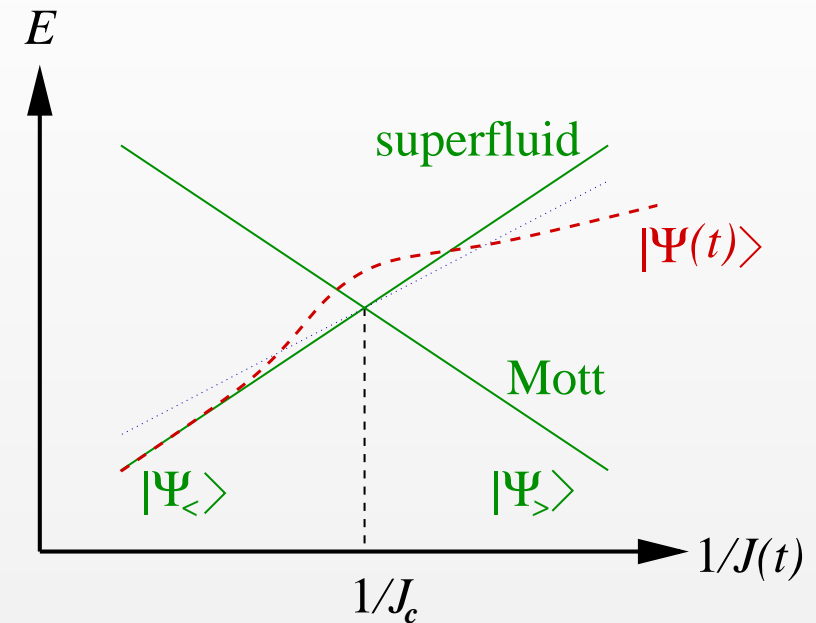
Sweeping from the superfluid to Mott phase in the Bose-Hubbard model,

cond-mat/0605121



Similarities to Cosmic Inflation

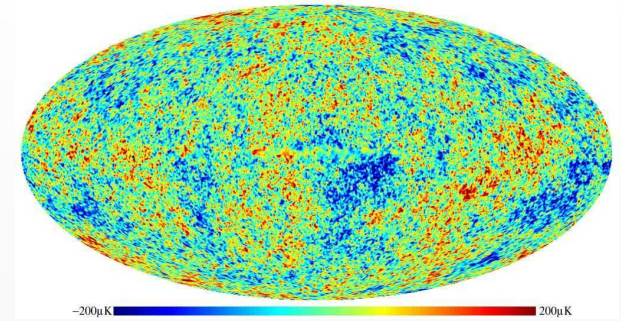
- Release of energy
→ (p)re-heating
- Robust against initial
(small-scale) perturbations
- Universality
(no fine-tuning)
- Amplification of quantum fluctuations



But: different spectrum in general

- Preferred frame (rest frame of medium)
- No unique/constant propagation speed
- Neglect of (quantum) back-reaction

Speculations...



Postulate:

- No (locally) preferred frame
- Unique/constant propagation speed

$$\mathcal{A} = \frac{1}{2} \int dt d^3r \frac{\dot{\Phi}^2 - (\nabla\Phi)^2}{t^2}$$

- \leftrightarrow scale-invariance $\mathcal{A}[\lambda t, \lambda \mathbf{r}] = \mathcal{A}[t, \mathbf{r}]$
- Dominated by (quantum) back-reaction?

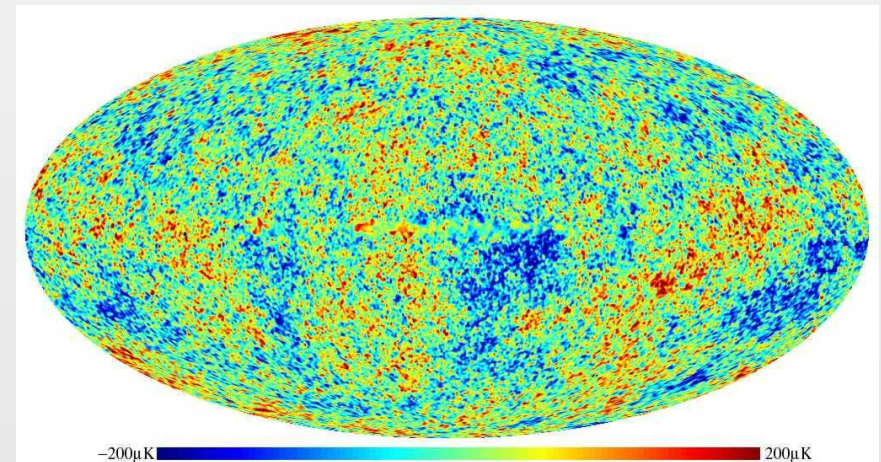
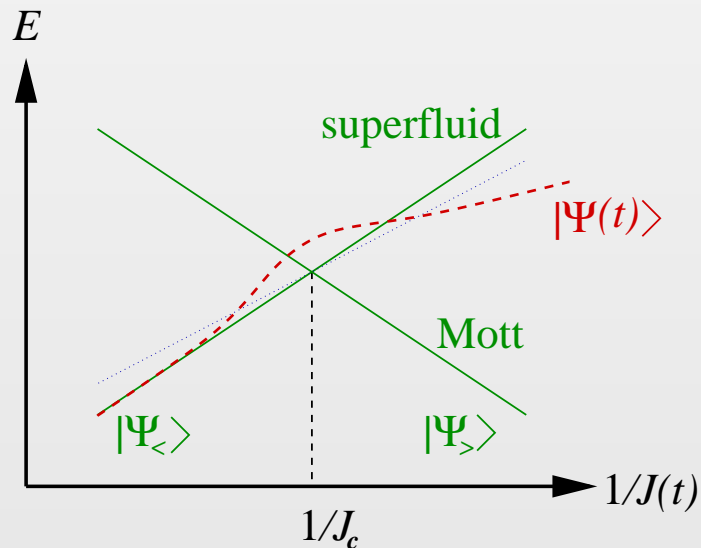
\rightarrow correct $1/k^3$ -spectrum (conformal de Sitter metric)

Was cosmic inflation just a phase transition?

R. S., Phys. Rev. Lett. **95**, 135703 (2005)

Summary

- Analogy between cosmic inflation and dynamical quantum phase transitions
- Effective “cosmic” horizons
 - loss of causal contact
 - non-adiabatic behaviour
 - amplification of quantum fluctuations
- Relation to real cosmic inflation?



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