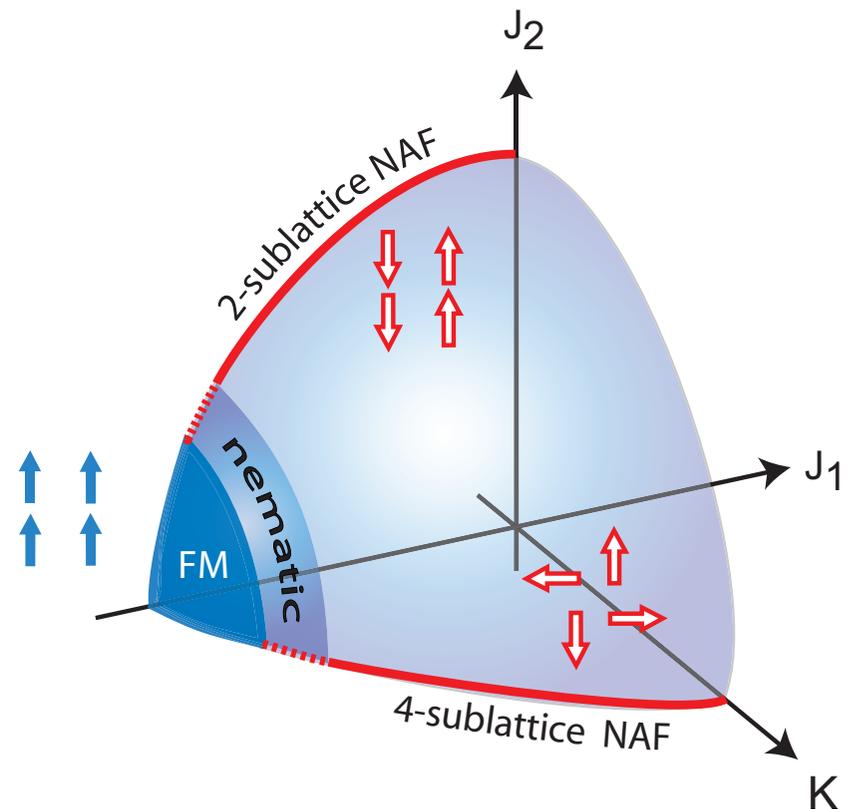


(How to have fun with) two-dimensional frustrated ferromagnets

nic shannon
(Bristol, UK)

Les Houches 20/6/6

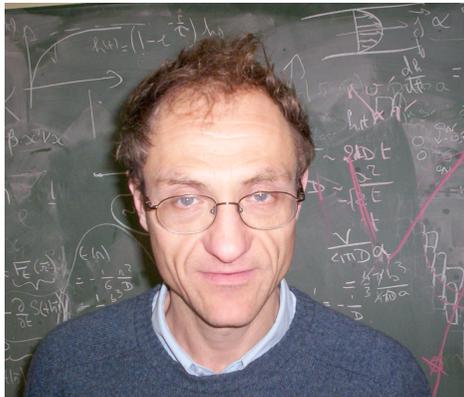


thanks to...



Tsutomu Momoi

RIKEN



Philippe Sindzingre

Paris VI, Jussieu

MAX-PLANCK-INSTITUT
FÜR CHEMISCHE PHYSIK FESTER STOFFE



Max-Planck-Institut für Physik komplexer Systeme
Nöthnitzer Str. 38 · D-01187 Dresden · Telefon +49(0)351 871-0 · eMail: info@mpipks-dresden.mpg.de



Institut du Développement et des Ressources
en Informatique Scientifique

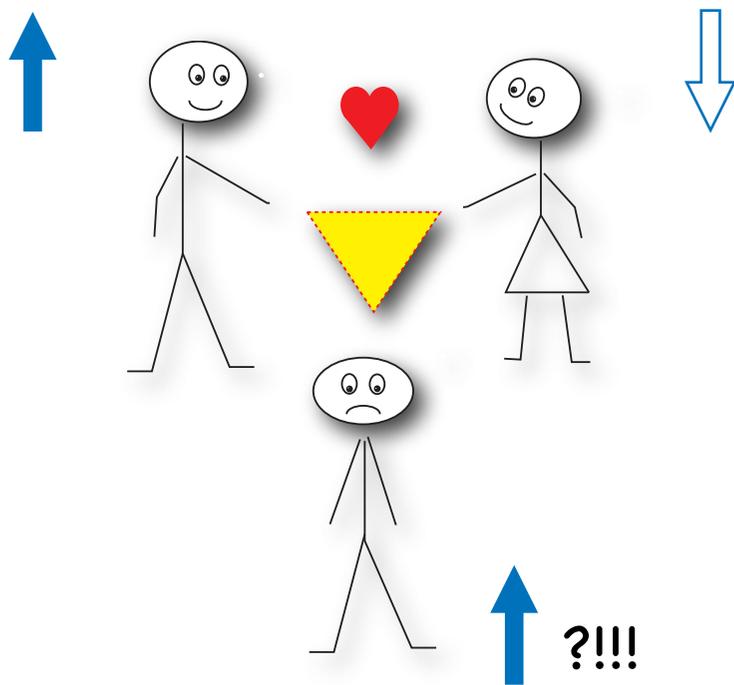
spin-1/2 on a triangular lattice

- spin liquids and the RVB idea -

spin-1/2 on a triangular lattice

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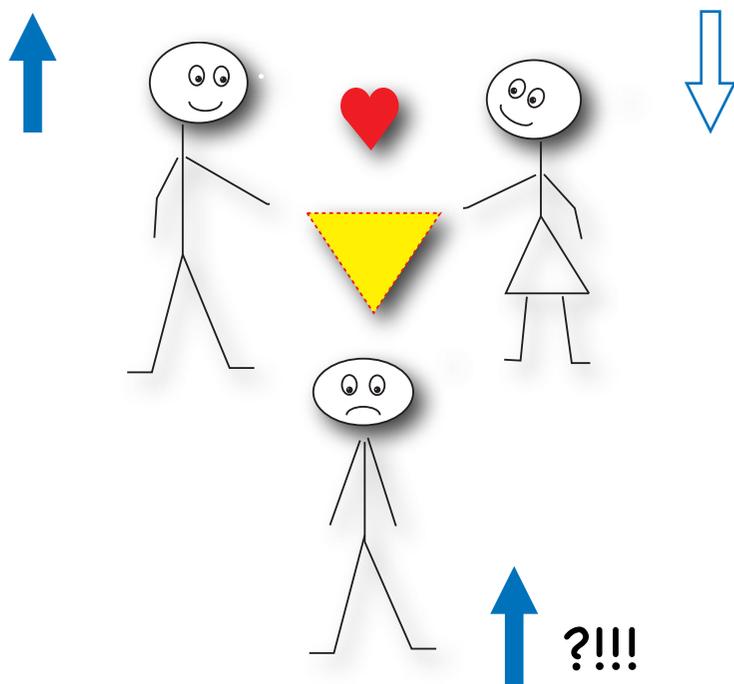
the eternal triangle



spin-1/2 on a triangular lattice

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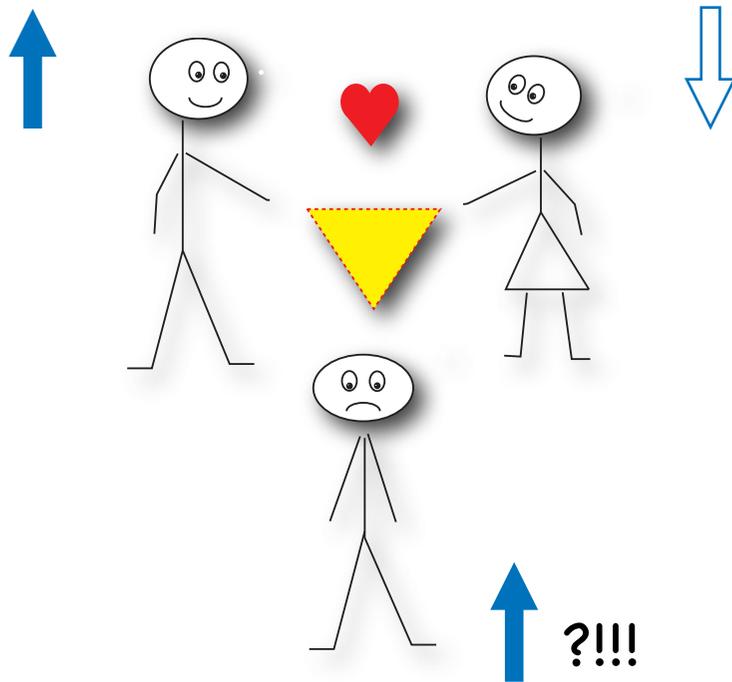


frustration, i.e. you can't please all
of the spins, all of the time

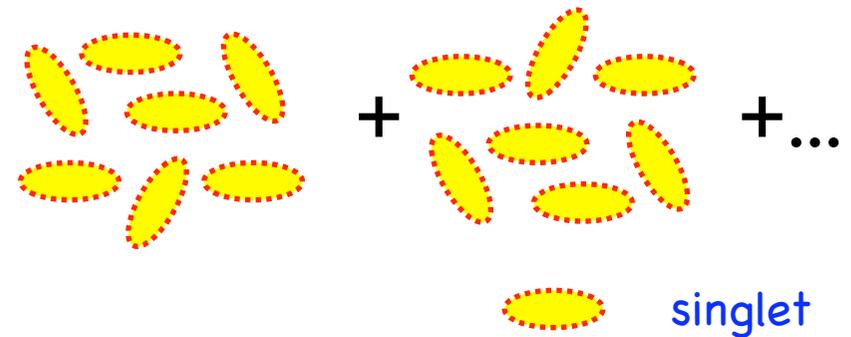
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Anderson's resonating valance bond
(RVB) state

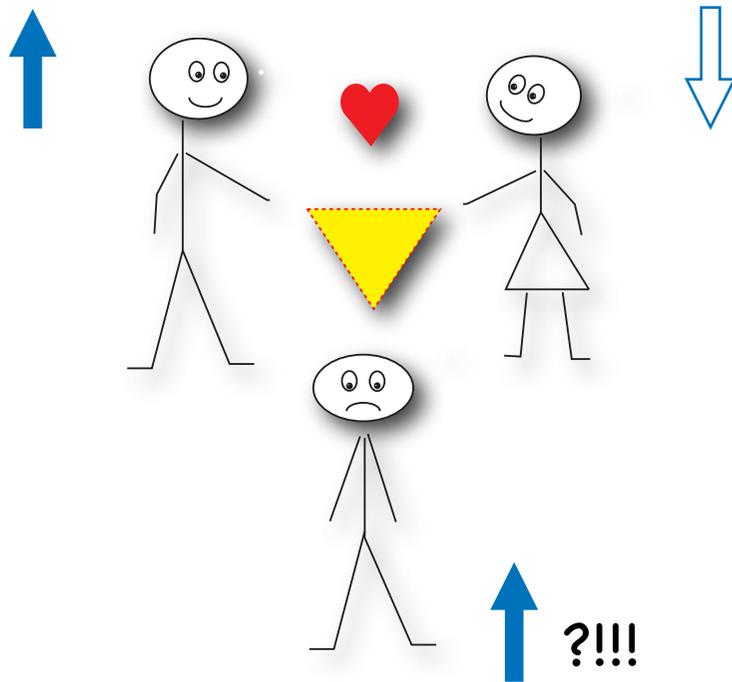


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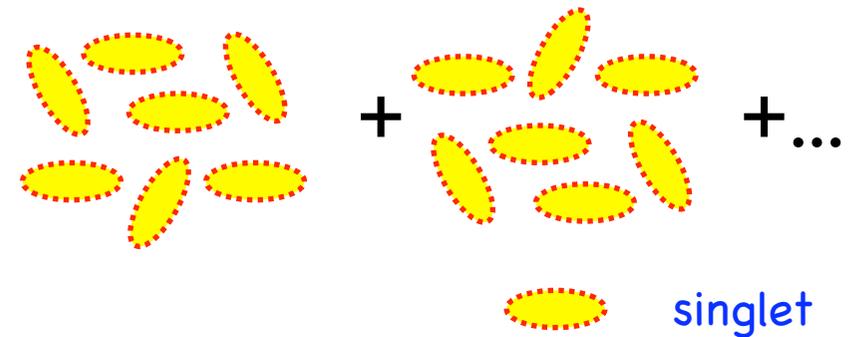
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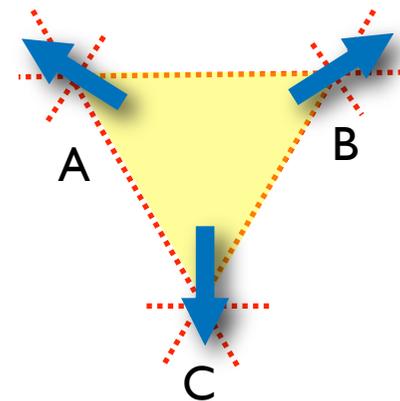


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Anderson's resonating valance bond (RVB) state



actual ground state of Heisenberg model on triangular lattice

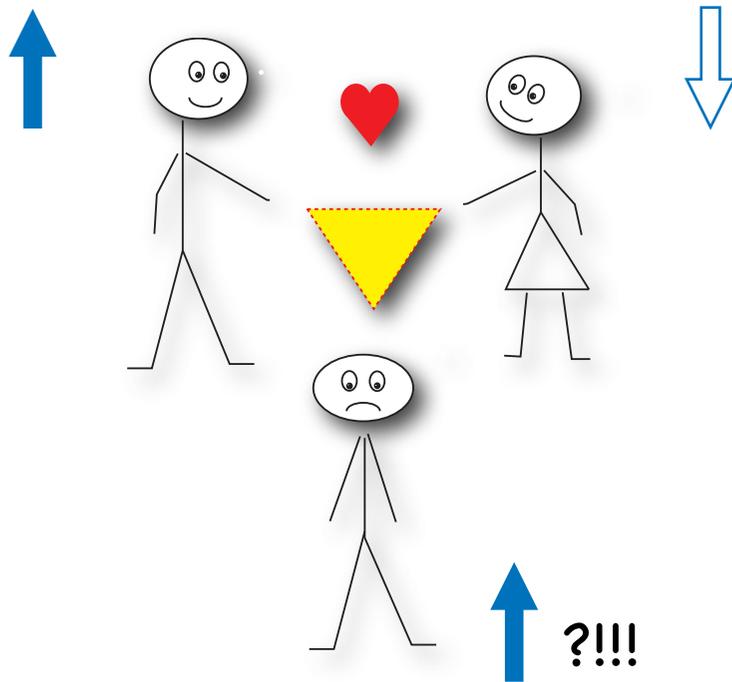


3-sublattice Néel state

spin-1/2 on a triangular lattice

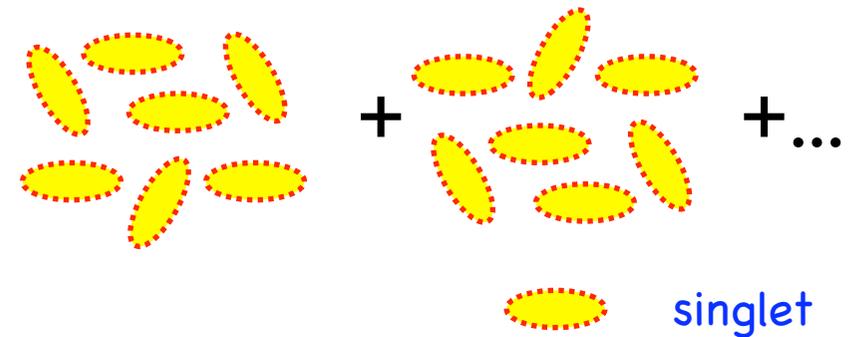
- spin liquids and the RVB idea -

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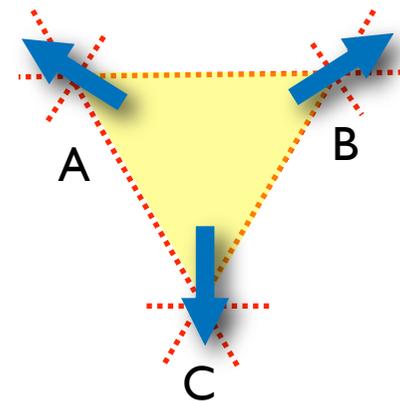


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3-sublattice Néel state

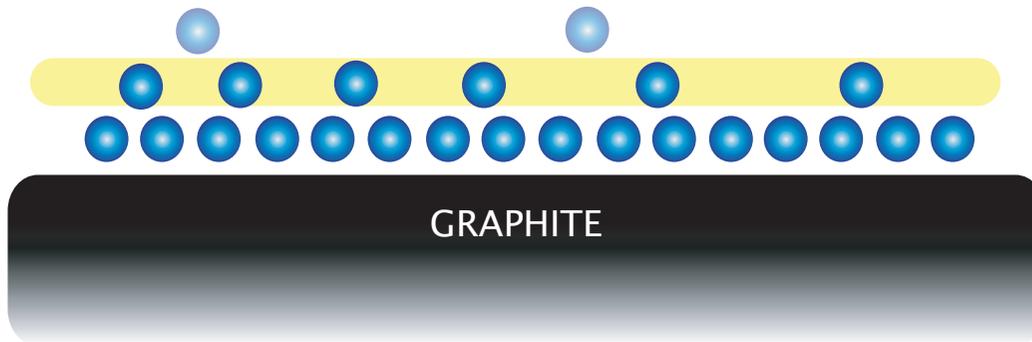
for a review, see e.g. Misguich and Lhullier in "Quantum Spin Systems" (2004 Diep)

in the beginning, God created He III...

the most perfect correlated Fermi system known to man

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the most perfect correlated Fermi system known to man

2D incarnation - He III on graphite



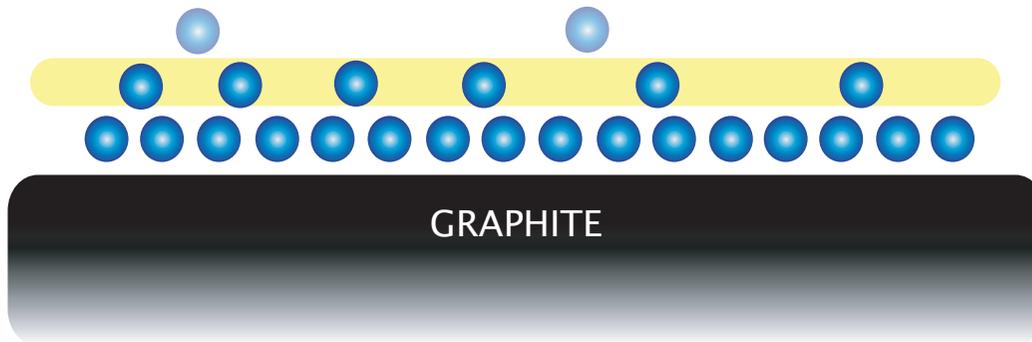
3rd layer - ignore

2nd layer - 2D FL/magnet

1st layer - paramagnetic solid

in the beginning, God created He III... the most perfect correlated Fermi system known to man

2D incarnation - He III on graphite

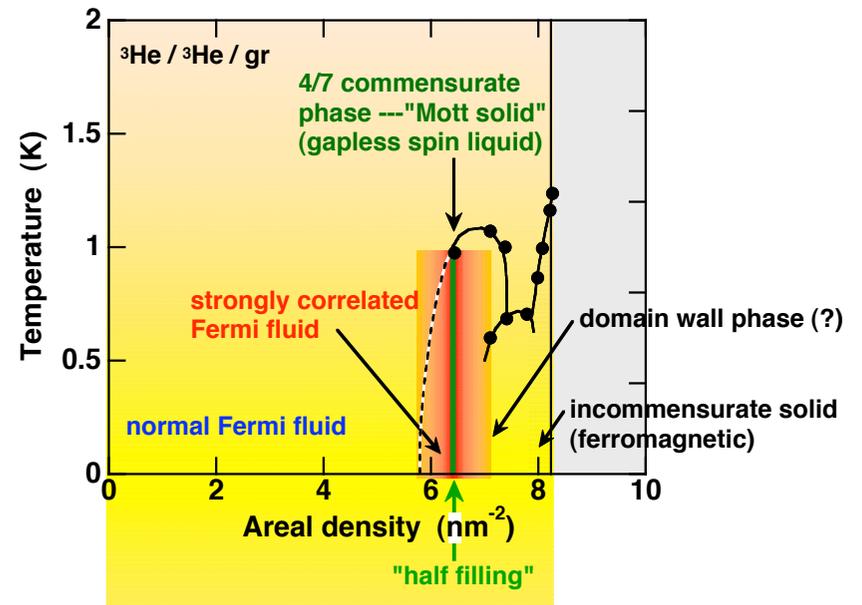


3rd layer - ignore
2nd layer - 2D FL/magnet
1st layer - paramagnetic solid

Fermi liquid in second layer becomes magnetic solid with increasing density :

high density solid is FM

something very special happens at low densities...



the first true spin liquid...

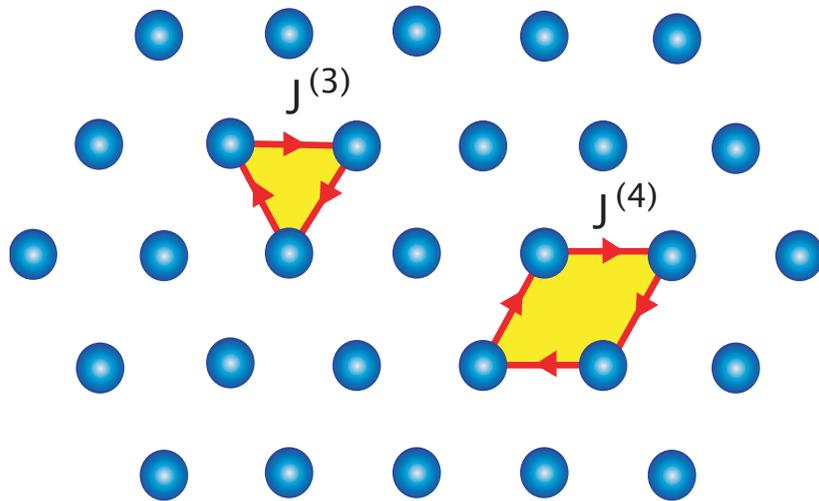
...in a 2D triangular lattice frustrated FM

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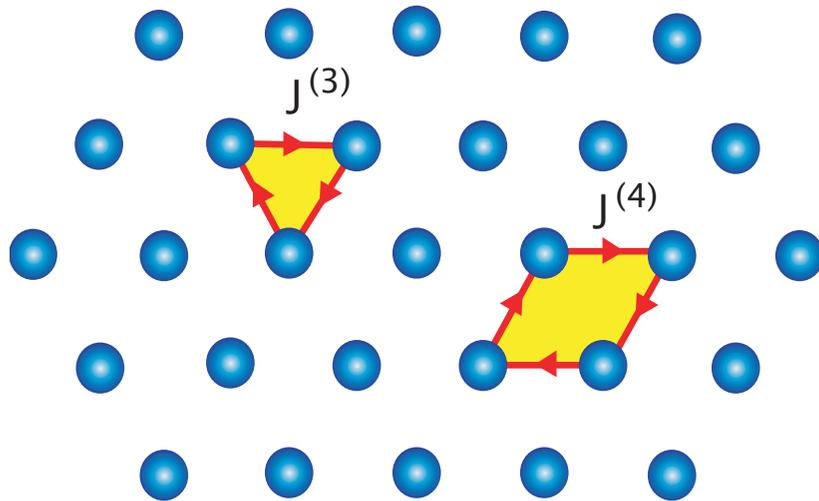


2nd layer magnetism controlled
by competition between
FM 3-spin exchange
and AF 4-spin exchange

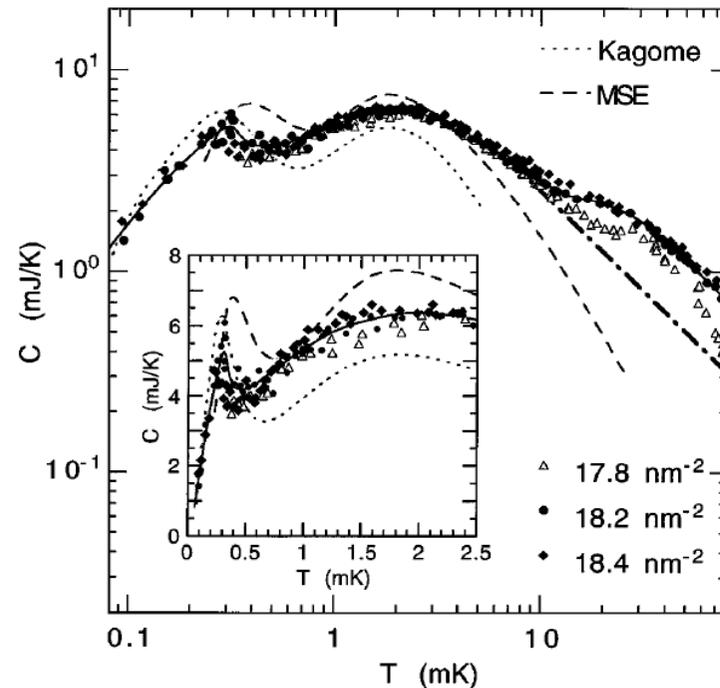
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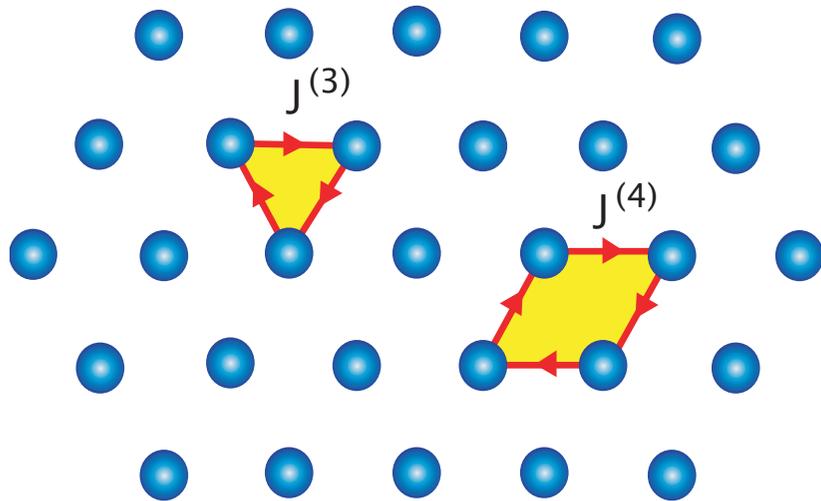


no magnetic order down to 0.1 mK !!!

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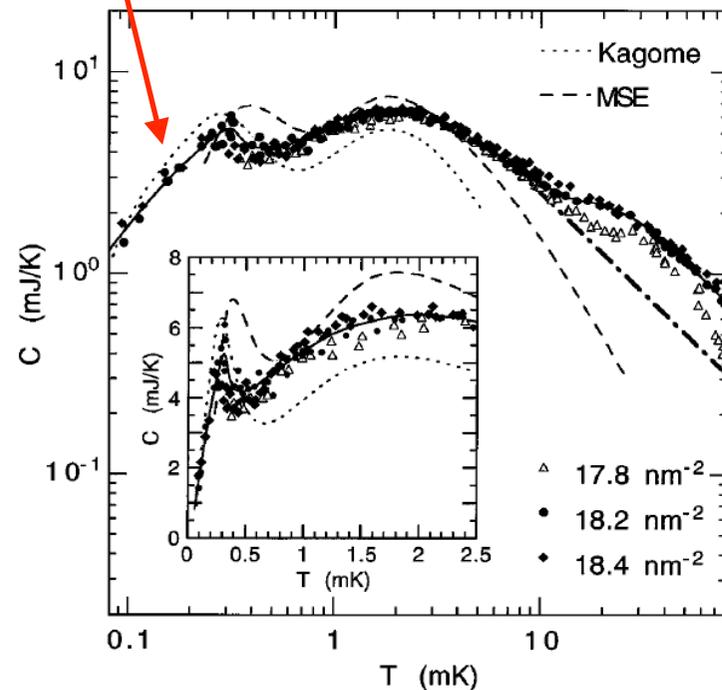
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linear specific heat
(c.f. 2D FM)

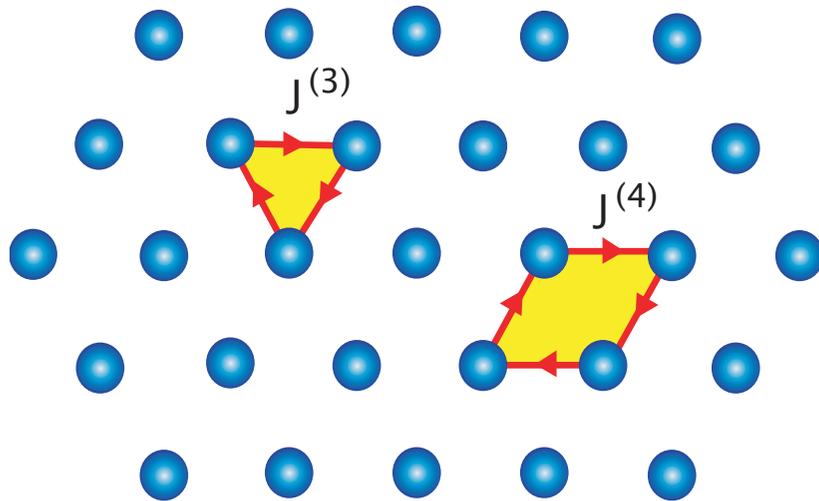


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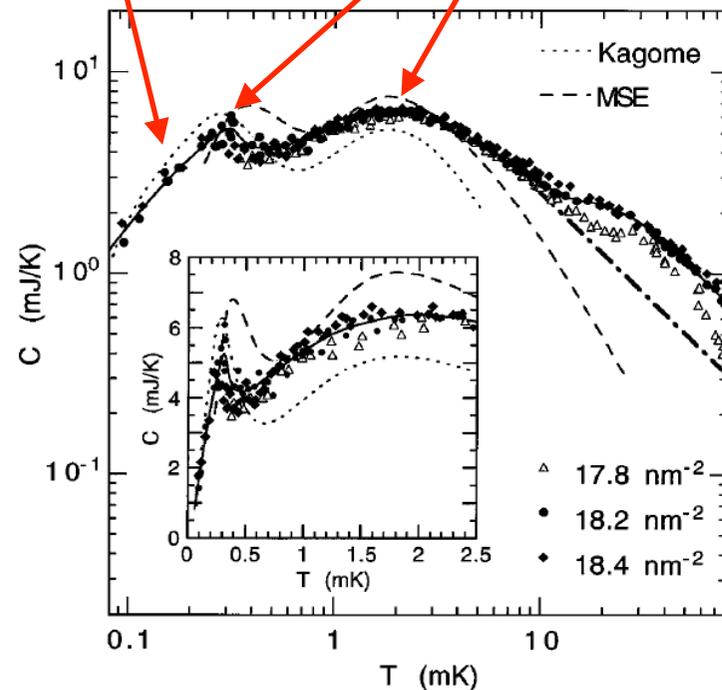
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2nd layer magnetism controlled
by competition between
FM 3-spin exchange
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linear specific heat
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double peak structure



no magnetic order down to 0.1 mK !!!



first example of a square lattice frustrated FM

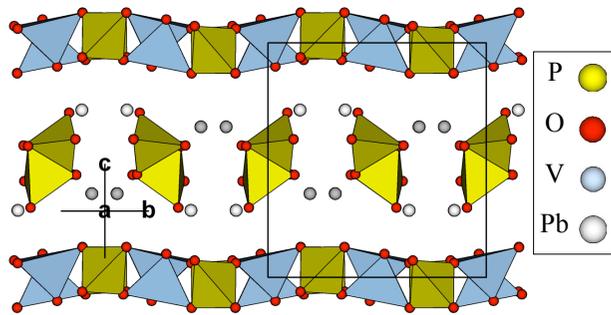
E. Kaul *et al.*, JMMM 272-276 (II), 922 (2004)

Pb₂VO(PO₄)₂

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Pb₂VO(PO₄)₂ : Structure



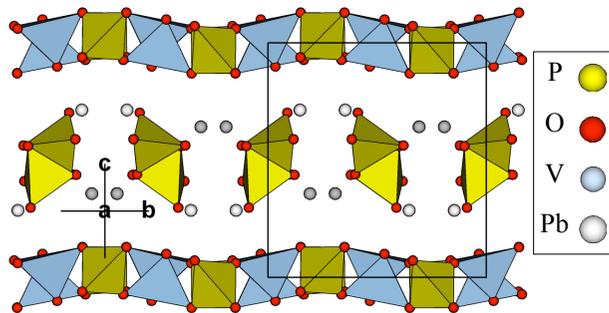
spin-1/2 V₄₊ in layered pyramids

Pb₂VO(PO₄)₂

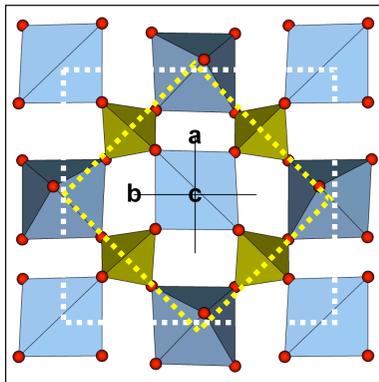
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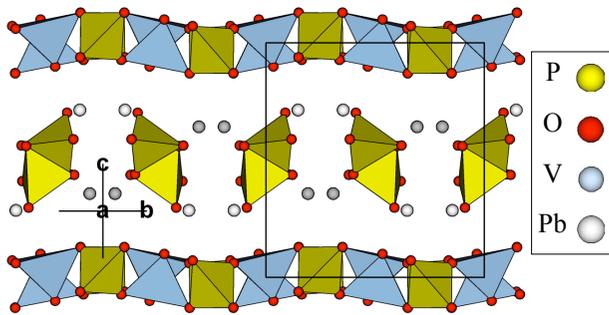
two different exchange paths
- both n.n. and n.n.n. bonds -

Pb₂VO(PO₄)₂

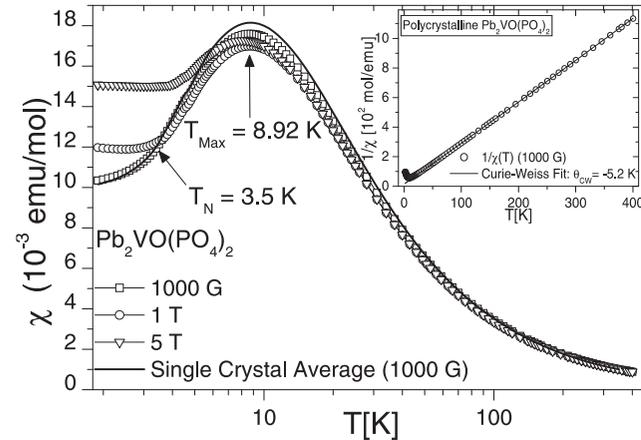
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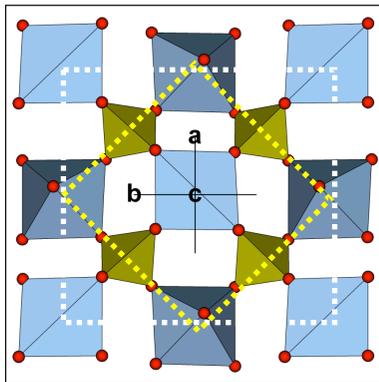
Pb₂VO(PO₄)₂ : Structure



spin-1/2 V₄₊ in layered pyramids



linear χ -inverse \Rightarrow frustrated magnet



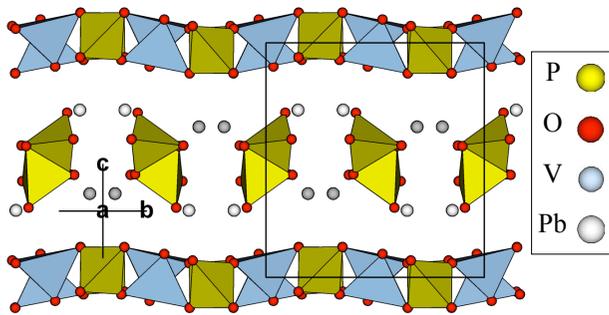
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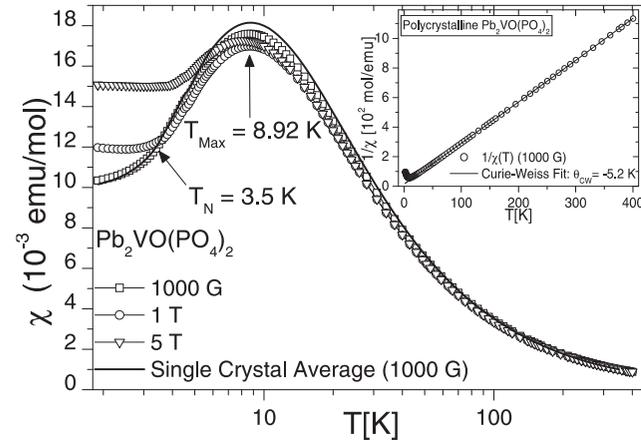
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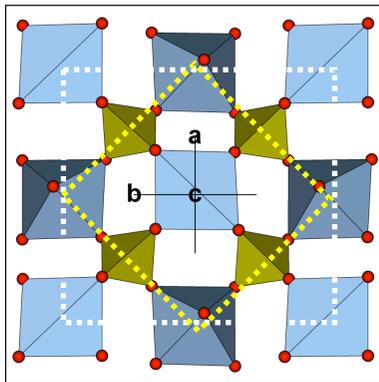
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$$J_2/J_1 \approx -2$$

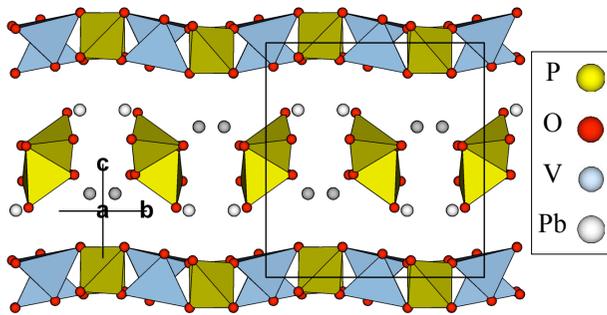
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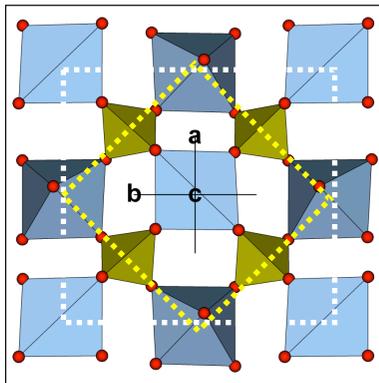
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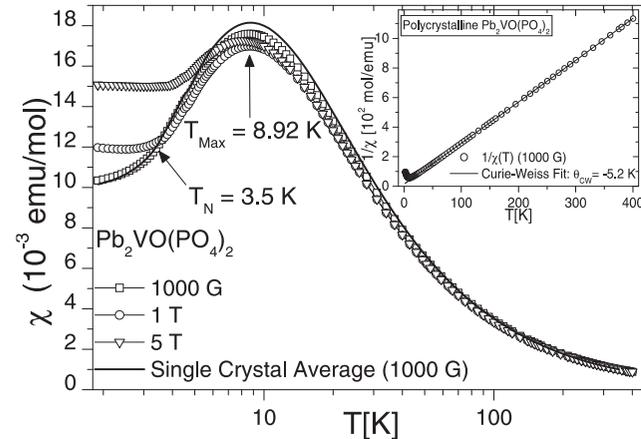
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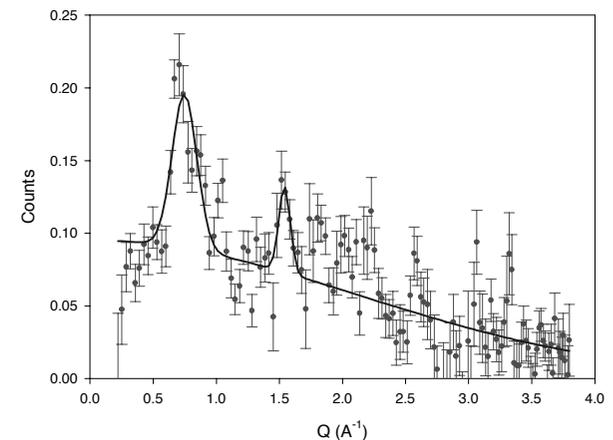


linear χ -inverse \Rightarrow frustrated magnet

$$J_2/J_1 \approx -2$$

ground state
is ($\pi, 0$)
collinear AF
with
reduced
moment

Average magnetic component for Pb₂VO(PO₄)₂ at 1.5 K



the "simplest" frustrated ferromagnet

extended FM Heisenberg model on square lattice

the "simplest" frustrated ferromagnet

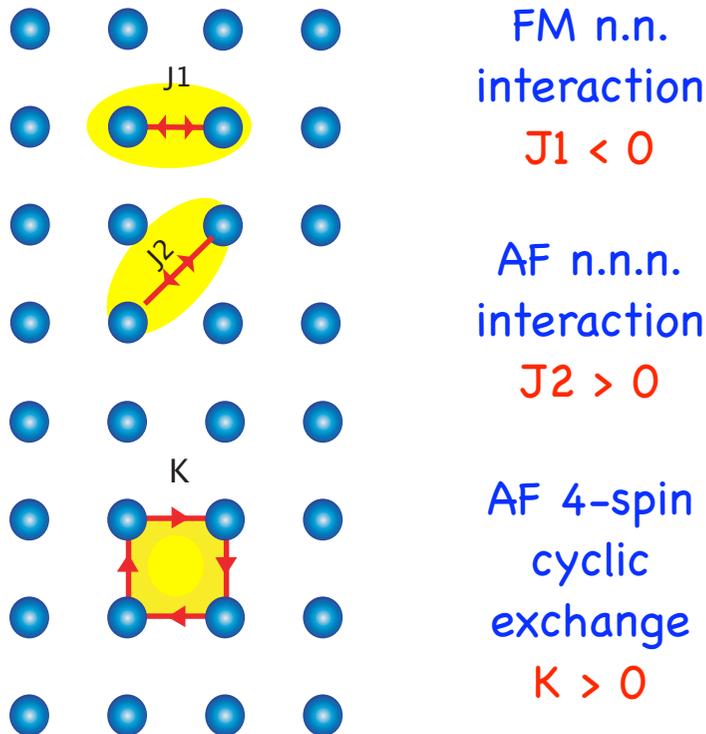
extended FM Heisenberg model on square lattice

$$\mathcal{H} = 2J_1 \sum_{\langle ij \rangle_1} \mathbf{S}_i \mathbf{S}_j + 2J_2 \sum_{\langle ij \rangle_2} \mathbf{S}_i \mathbf{S}_j + K \sum_{\langle 1234 \rangle} P_{1234} + P_{1234}^{-1}$$

the "simplest" frustrated ferromagnet

extended FM Heisenberg model on square lattice

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FM n.n.
interaction
 $J_1 < 0$

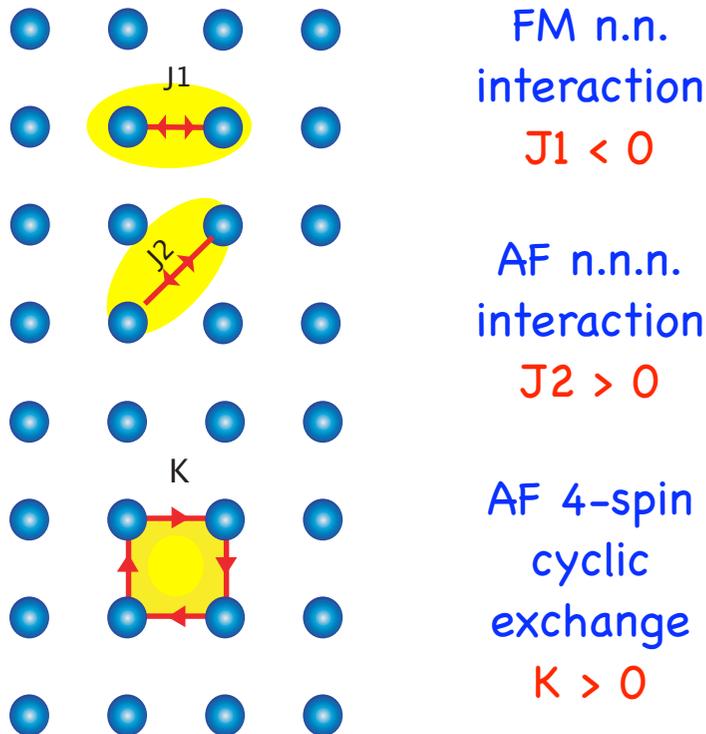
AF n.n.n.
interaction
 $J_2 > 0$

AF 4-spin
cyclic
exchange
 $K > 0$

the "simplest" frustrated ferromagnet

extended FM Heisenberg model on square lattice

$$\mathcal{H} = 2J_1 \sum_{\langle ij \rangle_1} \mathbf{S}_i \mathbf{S}_j + 2J_2 \sum_{\langle ij \rangle_2} \mathbf{S}_i \mathbf{S}_j + K \sum_{\langle 1234 \rangle} P_{1234} + P_{1234}^{-1}$$



FM n.n.
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 $J_1 < 0$

AF n.n.n.
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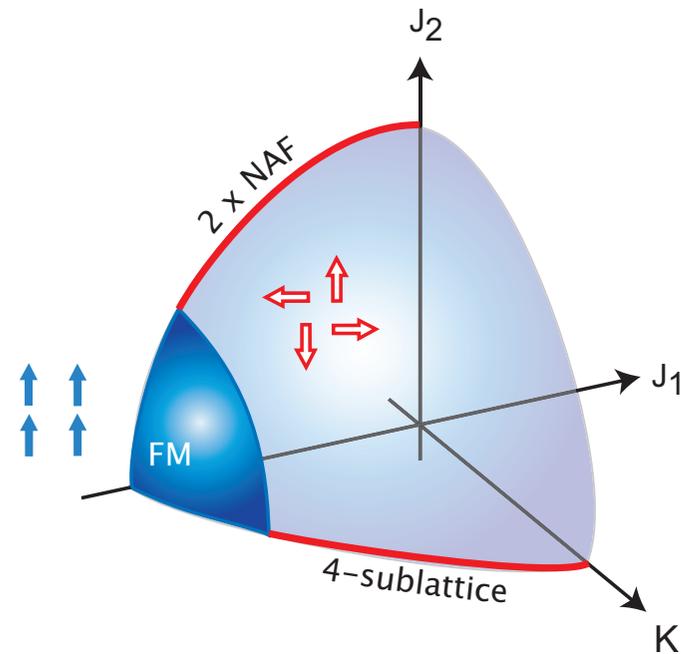
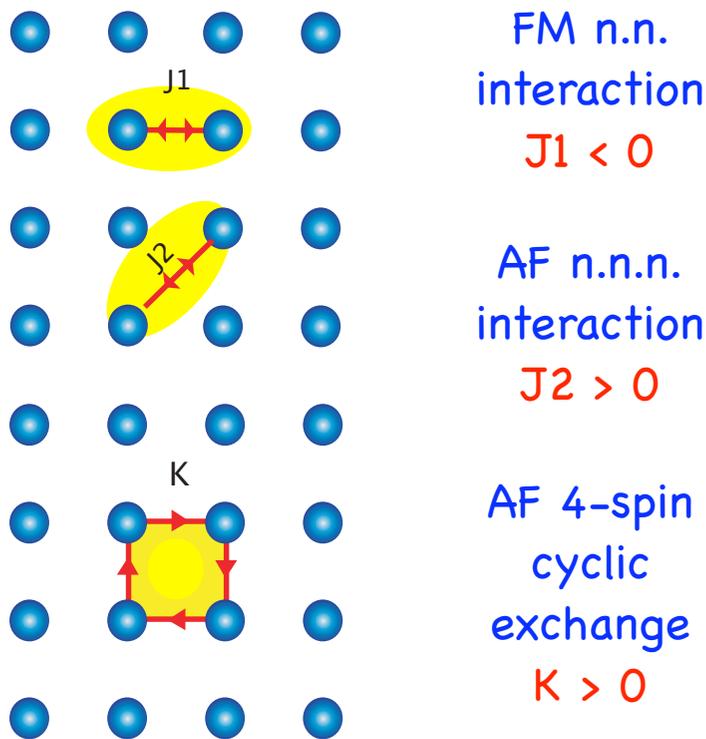
N.B.

$$P_{ijkl} + P_{ijkl}^{-1} = \vec{S}_i \cdot \vec{S}_j + \dots + 4 \left(\vec{S}_i \cdot \vec{S}_j \right) \left(\vec{S}_k \cdot \vec{S}_l \right) + \dots$$

the "simplest" frustrated ferromagnet

extended FM Heisenberg model on square lattice

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mean field phase diagram

N.B.

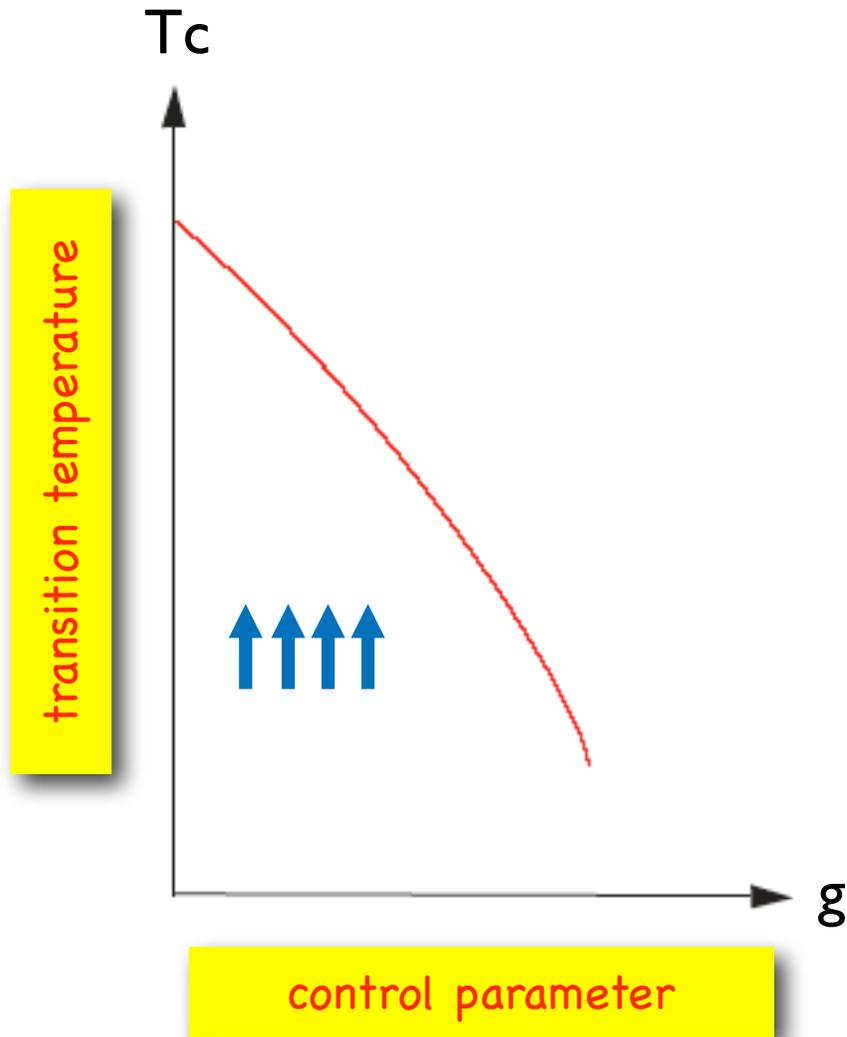
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quantum critical points

- and ferromagnetism -

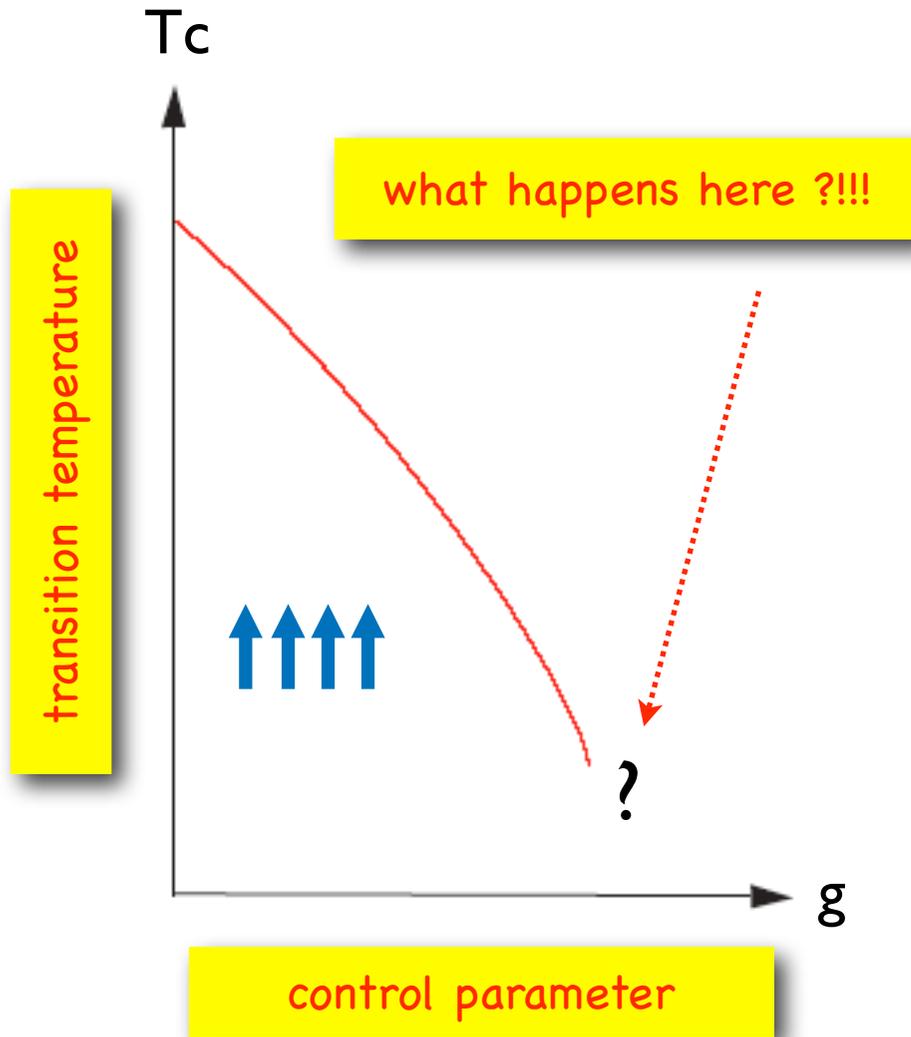
quantum critical points

- and ferromagnetism -



quantum critical points

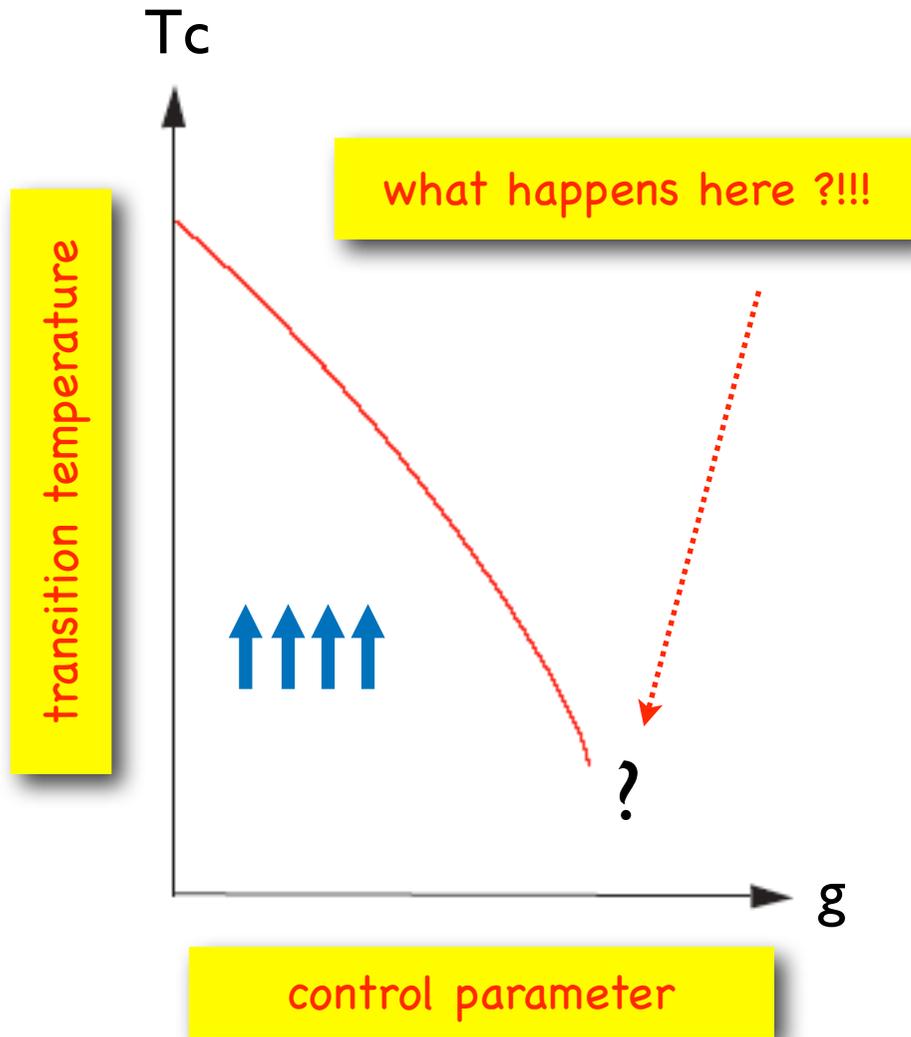
- and ferromagnetism -



quantum critical points

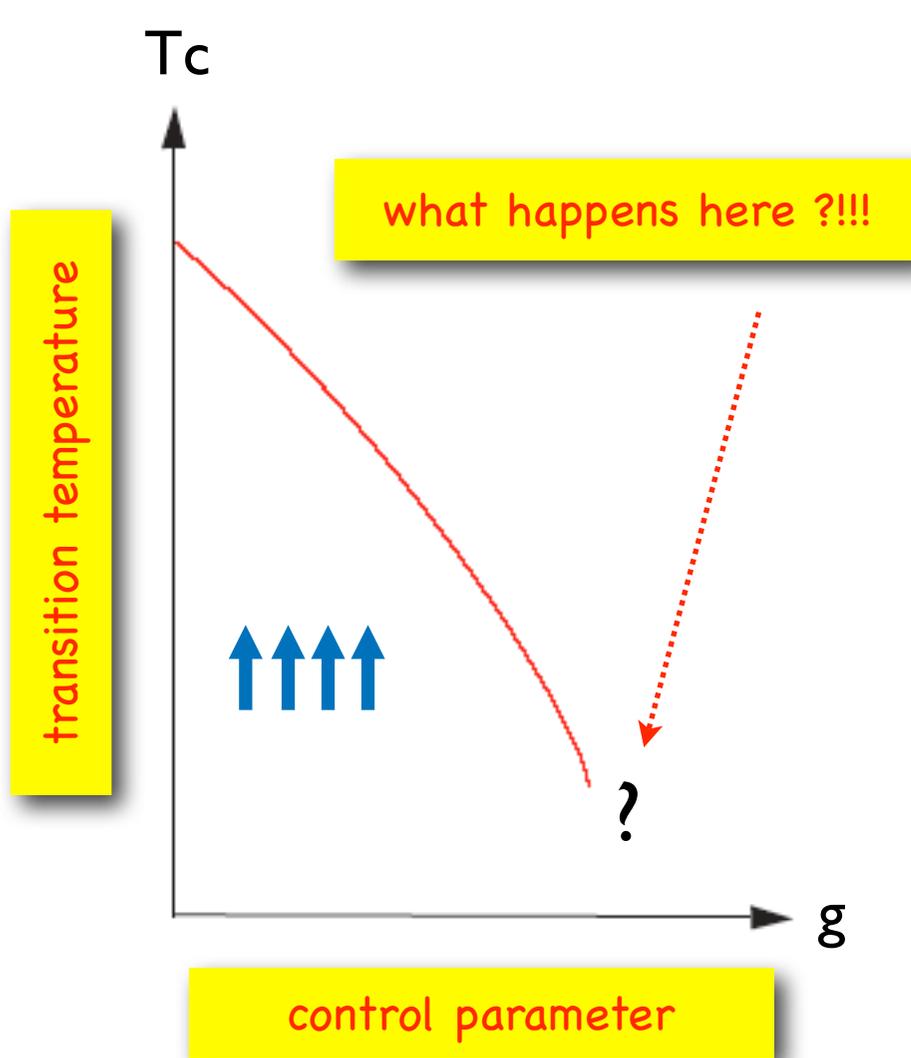
- and ferromagnetism -

manganites
(g=doping)
- phase separation



quantum critical points

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manganites
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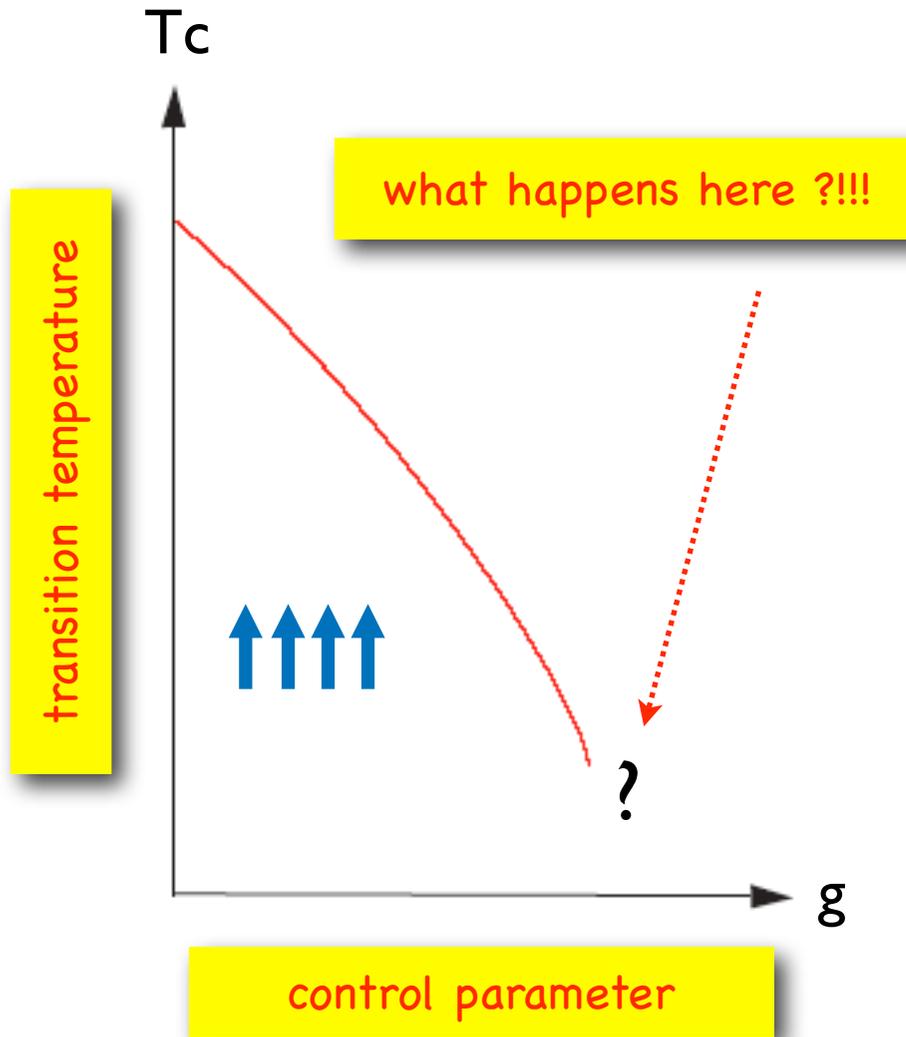
- phase separation

weak itinerant
ferromagnets
(g =pressure)

- superconductivity

quantum critical points

- and ferromagnetism -



manganites
(g=doping)

- phase separation

weak itinerant
ferromagnets
(g=pressure)

- superconductivity

frustrated
quantum spin
systems
(g=density,
chemical pressure)

- spin liquid ?!!!

how does the FM die ?

- nature of spin excitations at boundary with AF -

how does the FM die ?

- nature of spin excitations at boundary with AF -

“one magnon” dispersion :

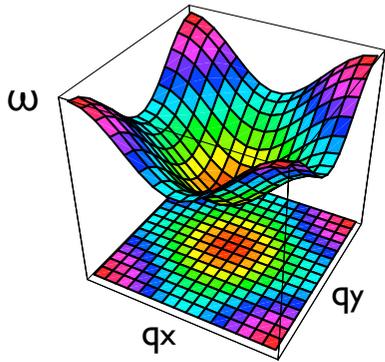
$$\omega(\mathbf{q}) = 8(|J_1| - 2K - J_2) - 4(|J_1| - 2K)[\cos q_x + \cos q_y] + 8J_2 \cos q_x \cos q_y$$

how does the FM die ?

- nature of spin excitations at boundary with AF -

“one magnon” dispersion :

$$\omega(\mathbf{q}) = 8(|J_1| - 2K - J_2) - 4(|J_1| - 2K)[\cos q_x + \cos q_y] + 8J_2 \cos q_x \cos q_y$$



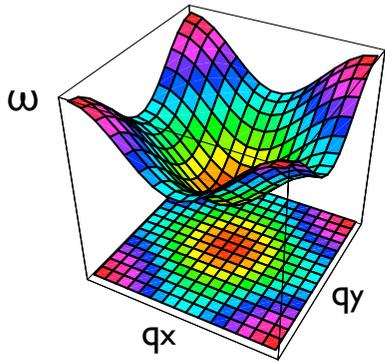
$$J=-1, J_2=K=0$$

how does the FM die ?

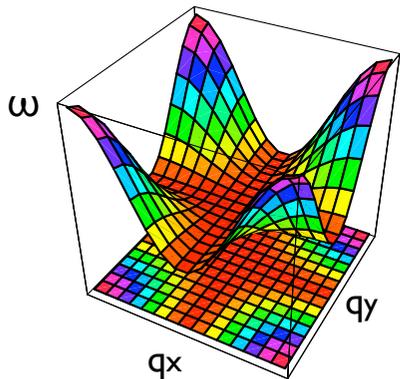
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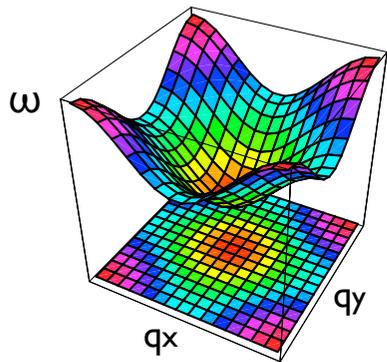
$J=-1, J_2=1/2, K=0$

how does the FM die ?

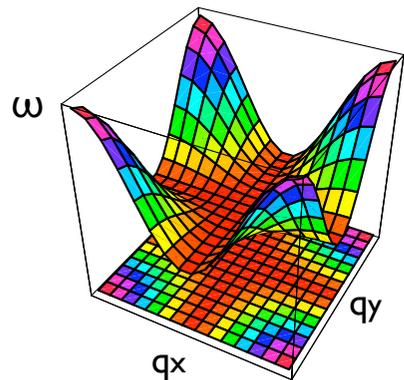
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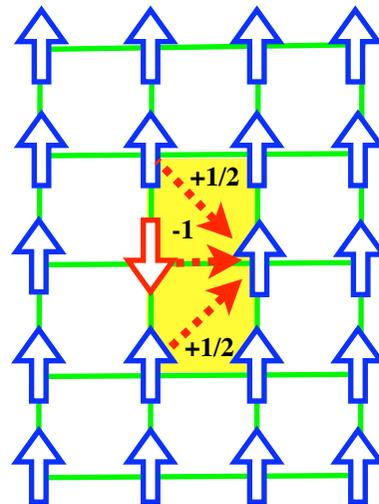
$J=-1, J_2=K=0$



$J=-1, J_2=1/2, K=0$

limiting case #1 :
 $J_1=-1, J_2 = 1/2, K=0$

J1-J2 model



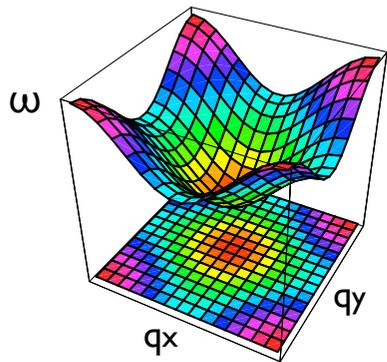
line zeros for $q_x = 0, q_y = 0$

how does the FM die ?

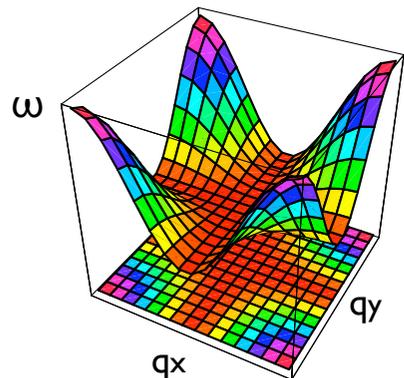
- nature of spin excitations at boundary with AF -

"one magnon" dispersion :

$$\omega(\mathbf{q}) = 8(|J_1| - 2K - J_2) - 4(|J_1| - 2K)[\cos q_x + \cos q_y] + 8J_2 \cos q_x \cos q_y$$



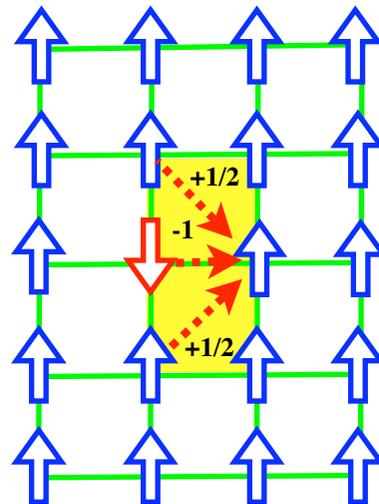
$J=-1, J_2=K=0$



$J=-1, J_2=1/2, K=0$

limiting case #1 :
 $J_1=-1, J_2 = 1/2, K=0$

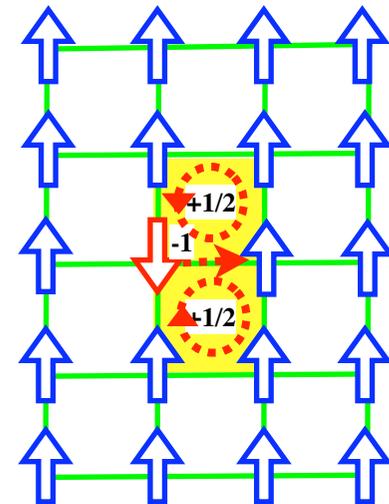
J1-J2 model



line zeros for $q_x = 0, q_y = 0$

limiting case #2 :
 $J_1=-1, J_2 = 0, K=1/2$

square lattice
MSE model



entire dispersion vanishes !!!

what kind of excitation works ?

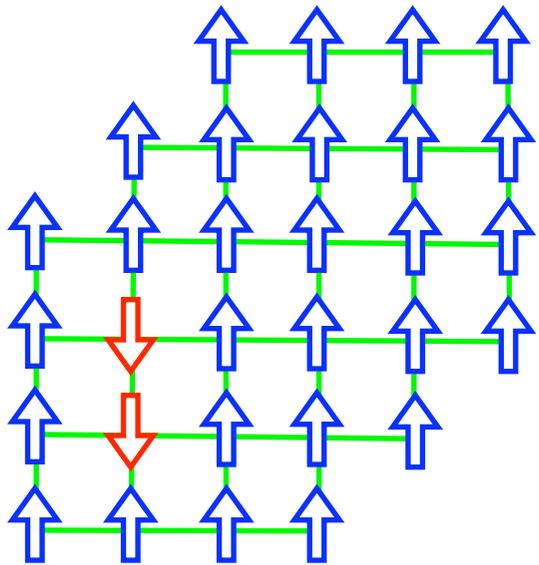
- two magnons are better than one -

what kind of excitation works ?

- two magnons are better than one -

square lattice MSE model

$J_1 = -1, J_2 = 0, K = 1/2$

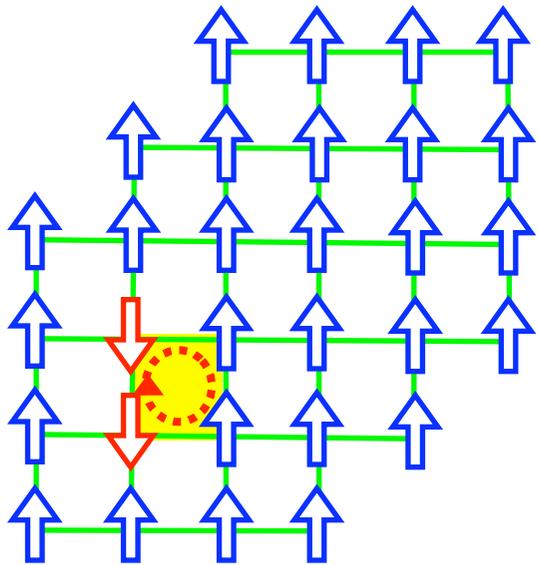


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square lattice MSE model

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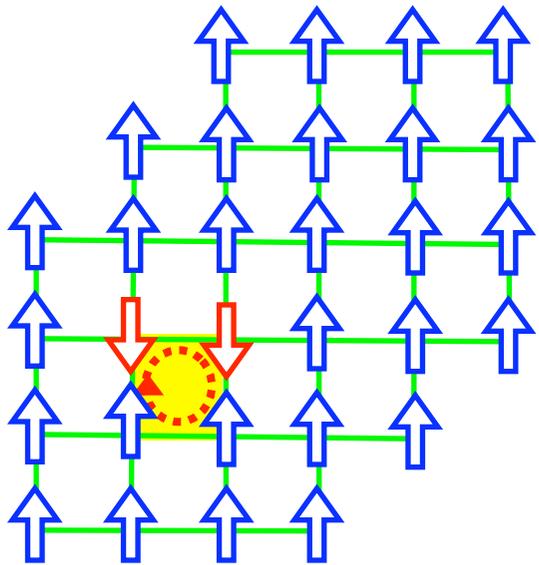


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square lattice MSE model

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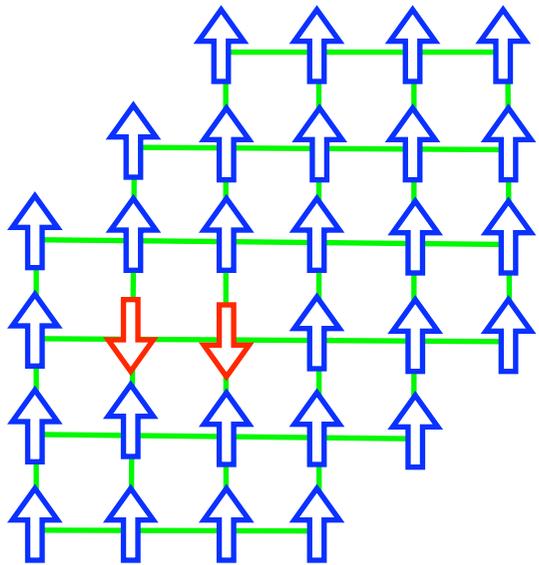


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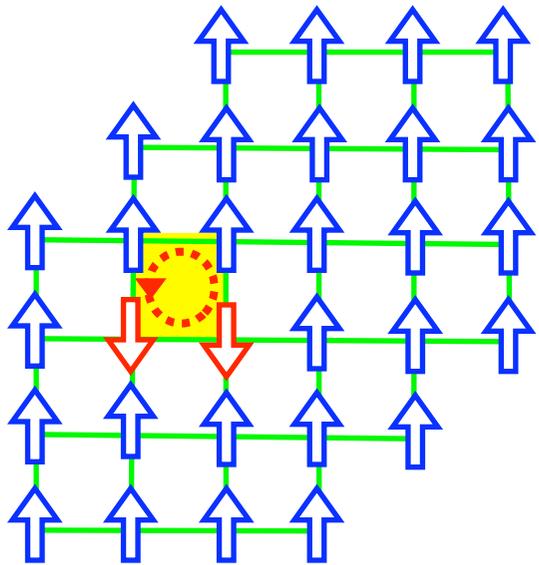


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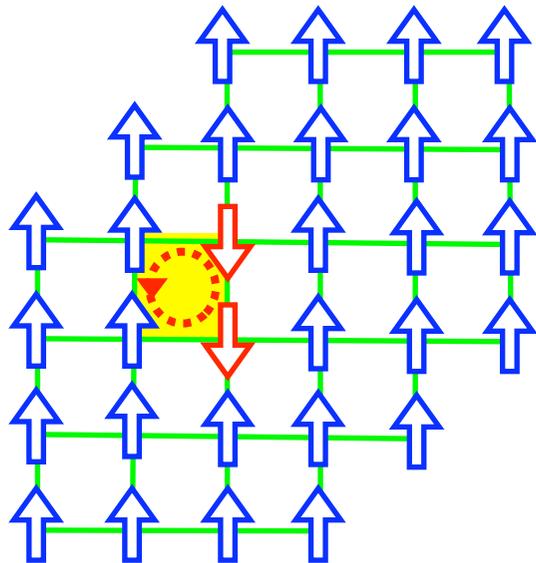


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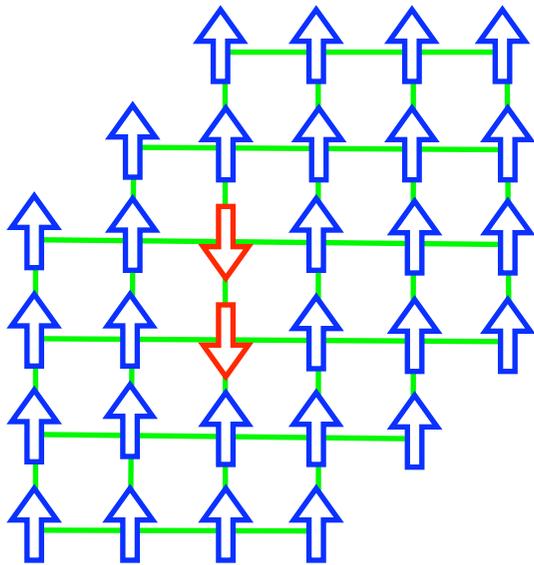


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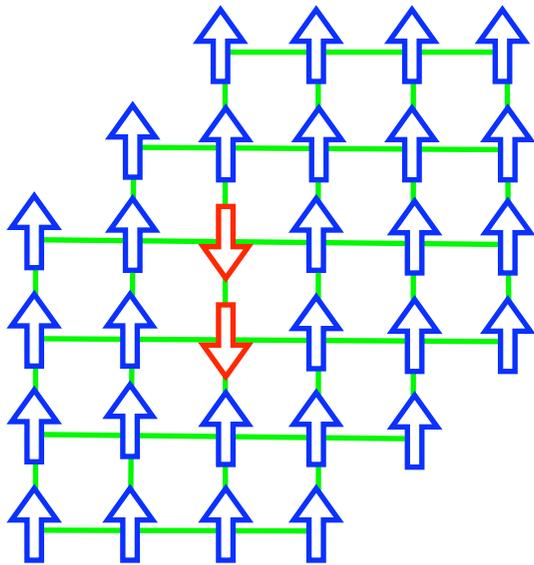


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but **pairs** of magnons can
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simple trial wave function
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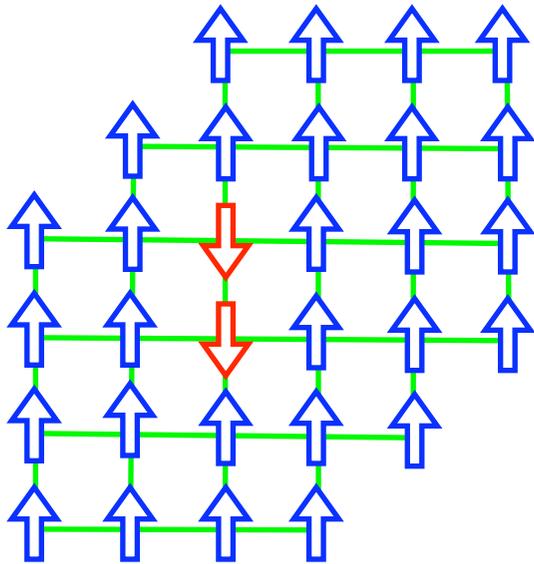
$$\frac{1}{\sqrt{2}} \left\{ \left| \begin{array}{cc} \uparrow & \uparrow \\ \downarrow & \downarrow \end{array} \right\rangle - \left| \begin{array}{cc} \downarrow & \uparrow \\ \downarrow & \uparrow \end{array} \right\rangle \right\} \exp(i\mathbf{q}\cdot\mathbf{r}/2)$$

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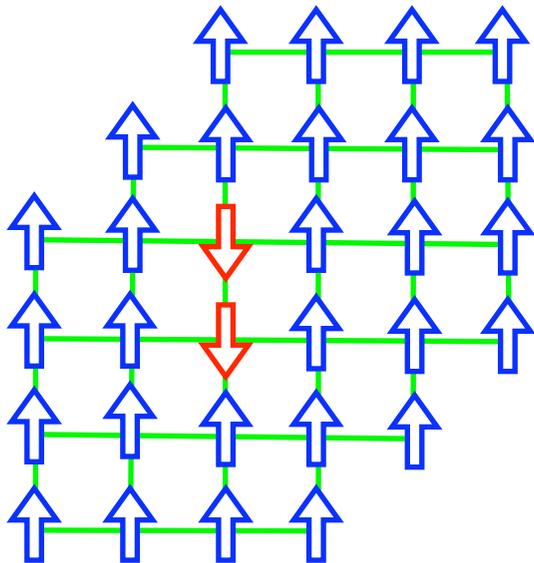
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for special point
 $J_1 = -1, J_2 = 0, K = 1/2$
this wave function is an
exact eigenstate

so what is the first instability of the FM ?

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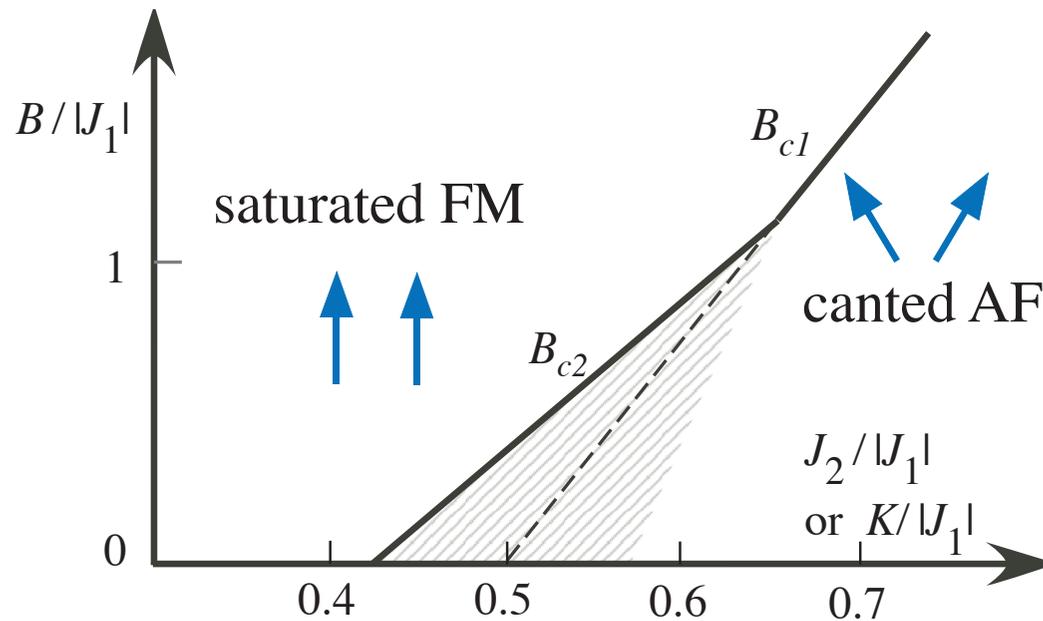
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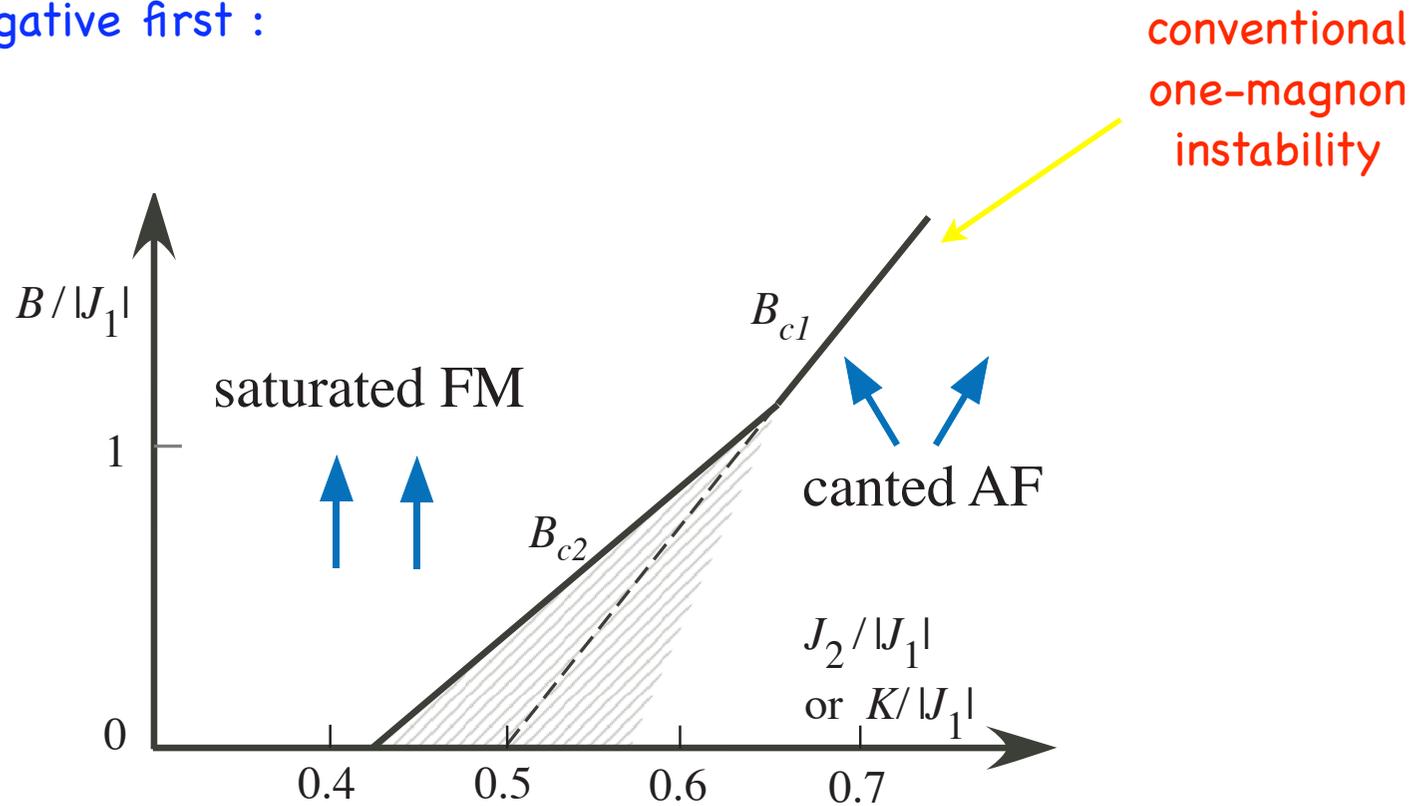
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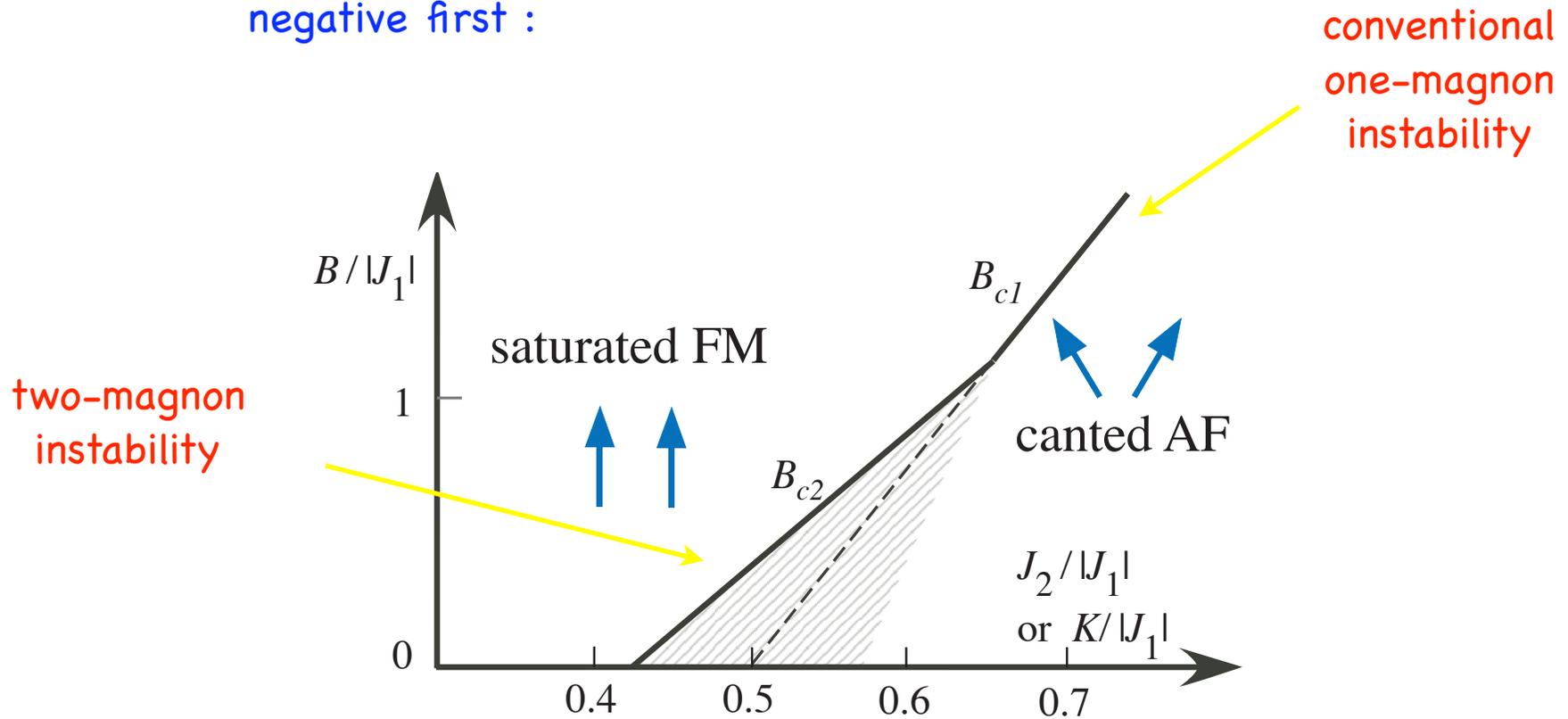
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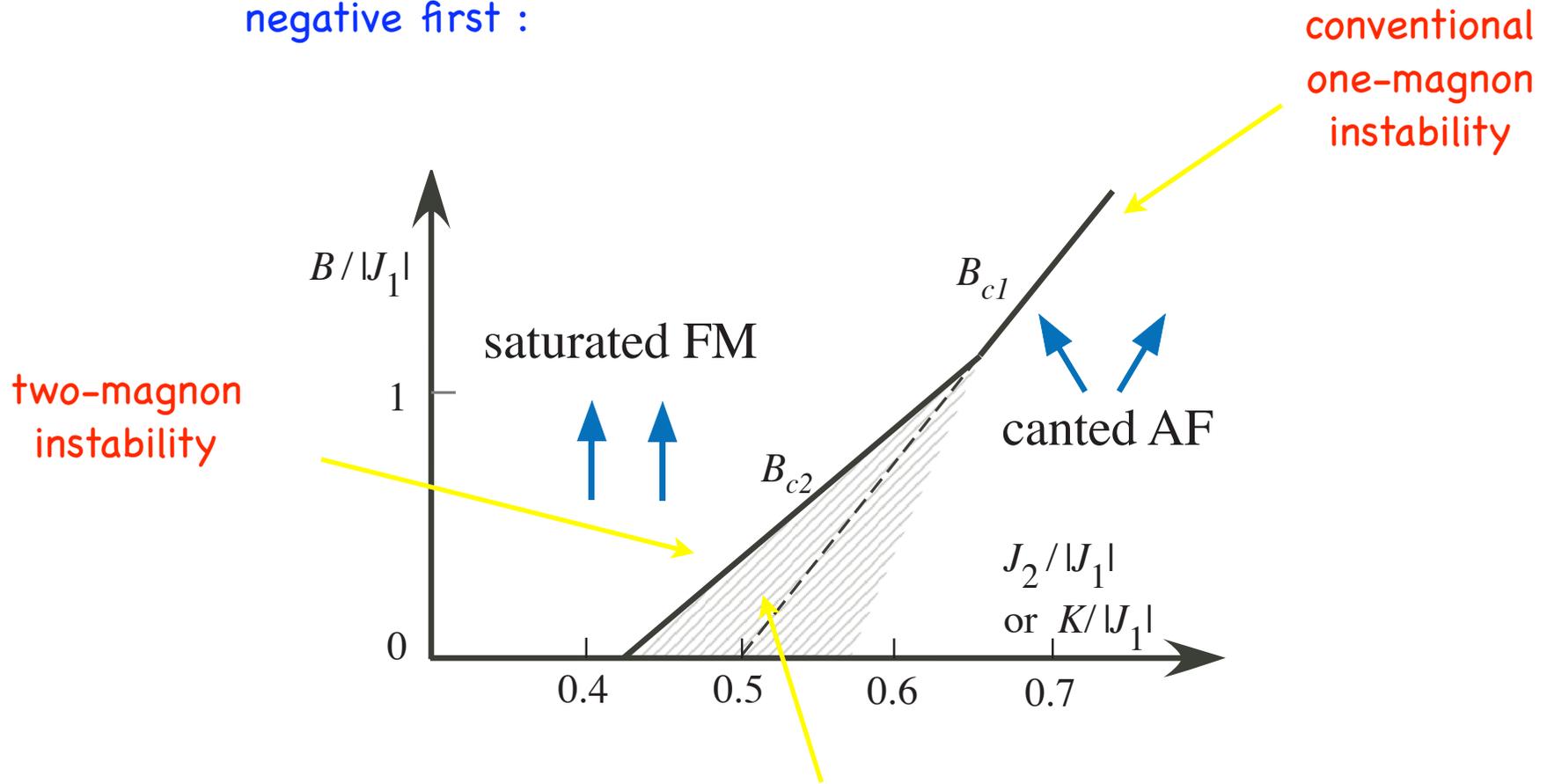
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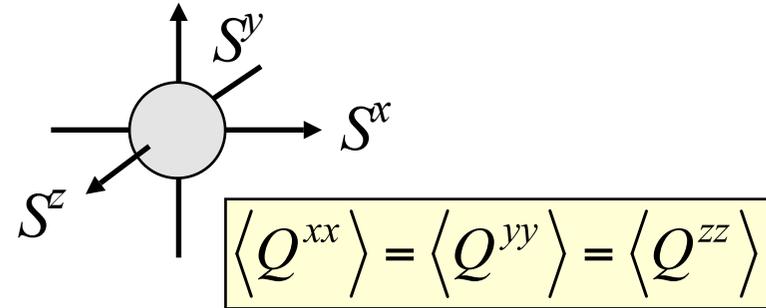
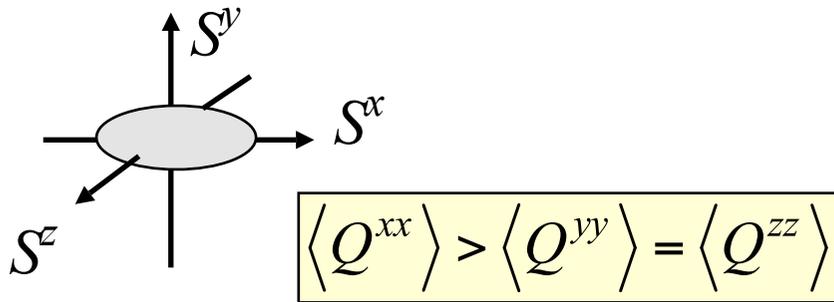
what is the nature of this phase ?!!!

a new idea - nematic order
- systems that don't know up from down -

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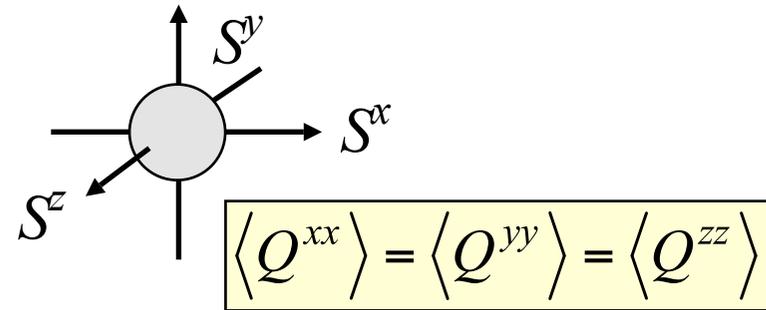
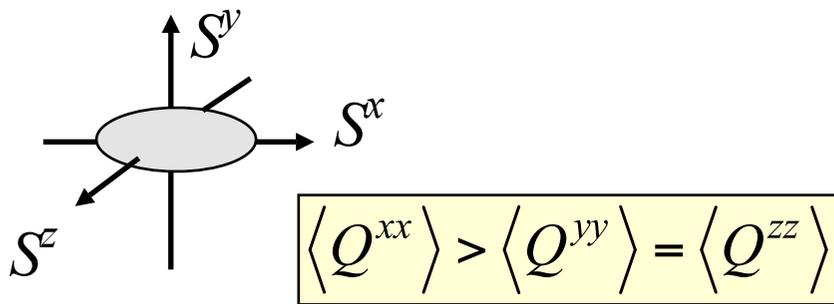
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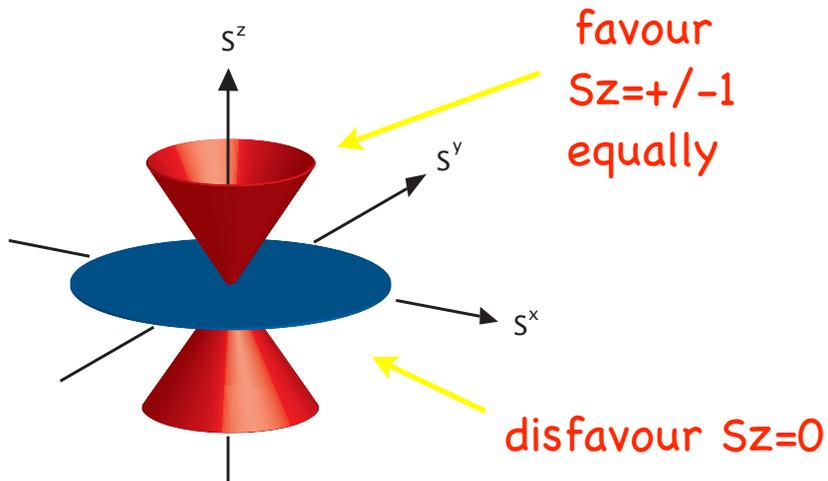
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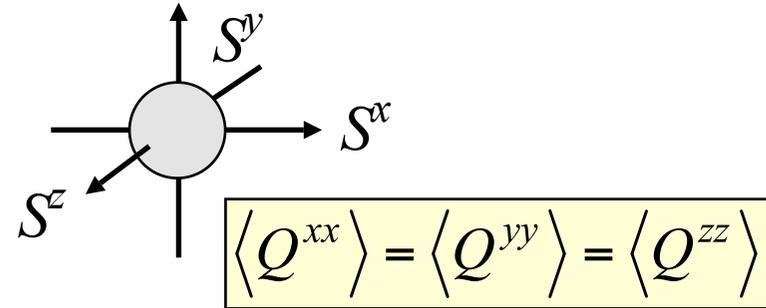
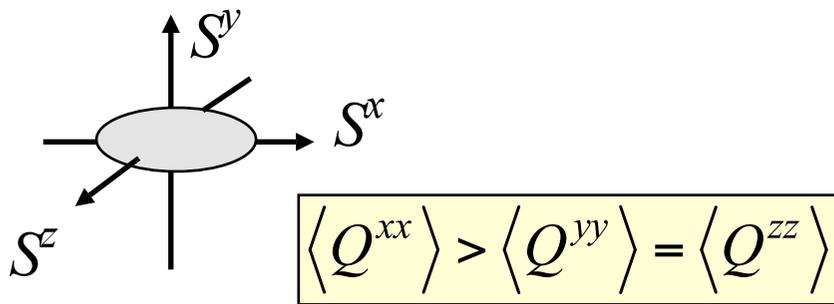
site-wise nematic works for spin-1 :



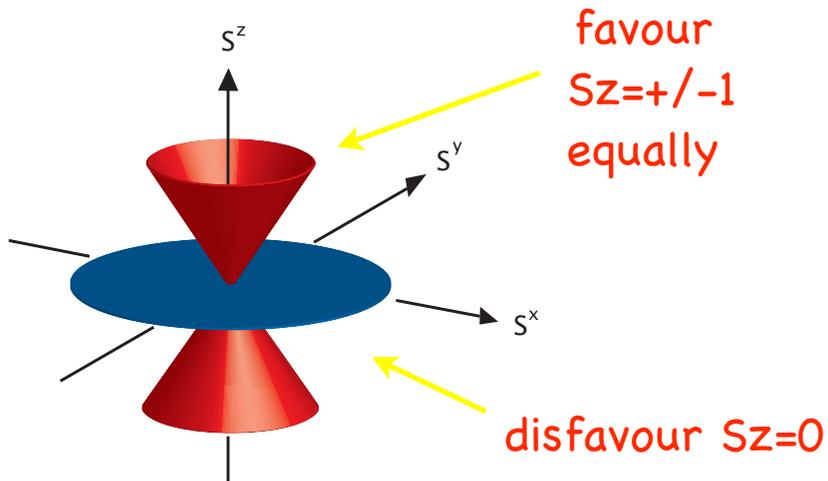
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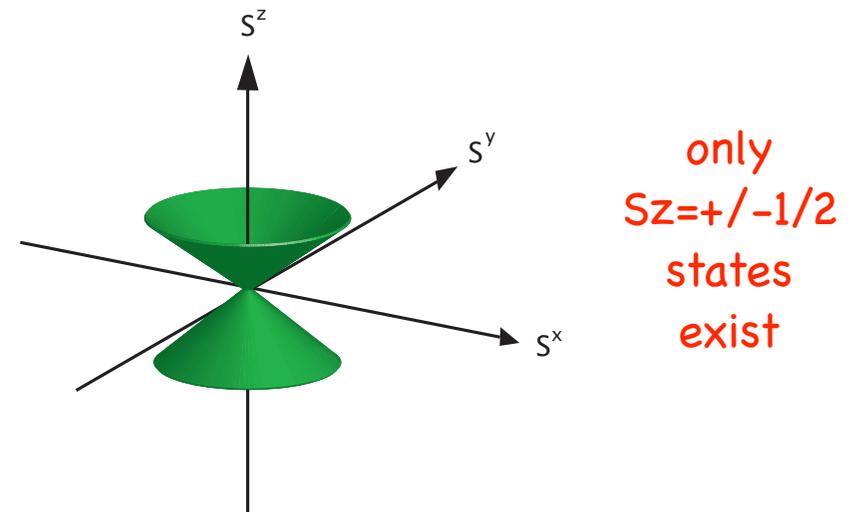
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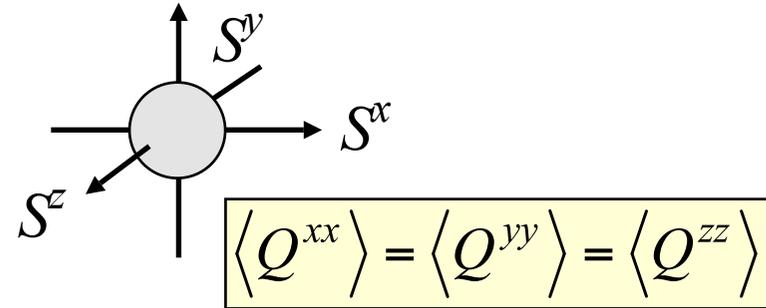
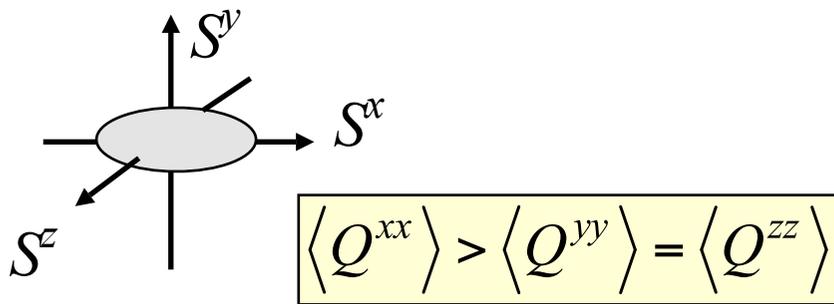
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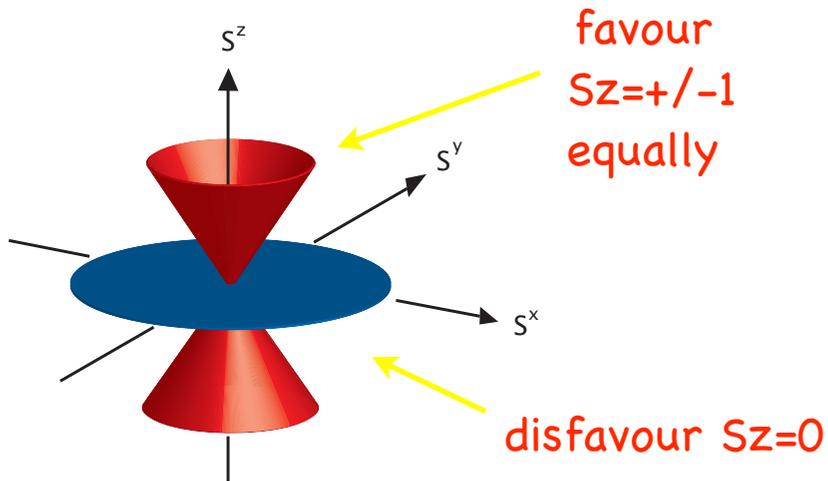
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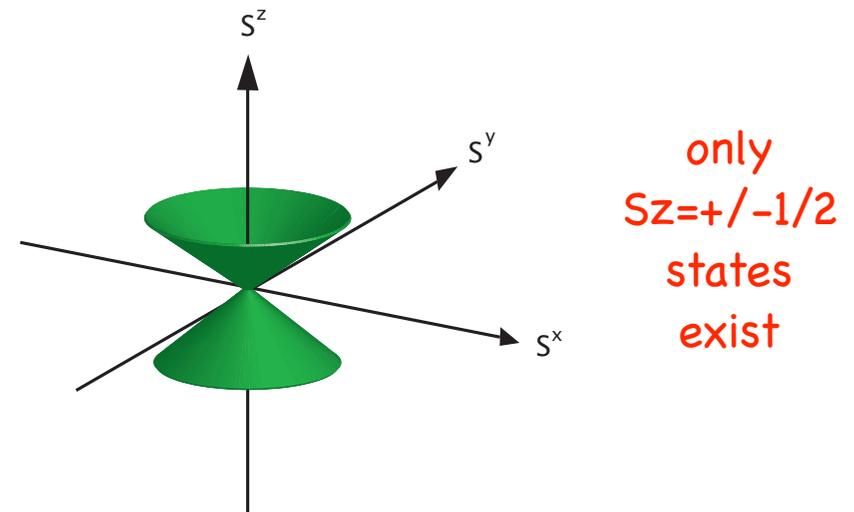
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for a spin-1 example see, e.g. : K. Harada and N. Kawashima, PRB **65**, 052403(2002)

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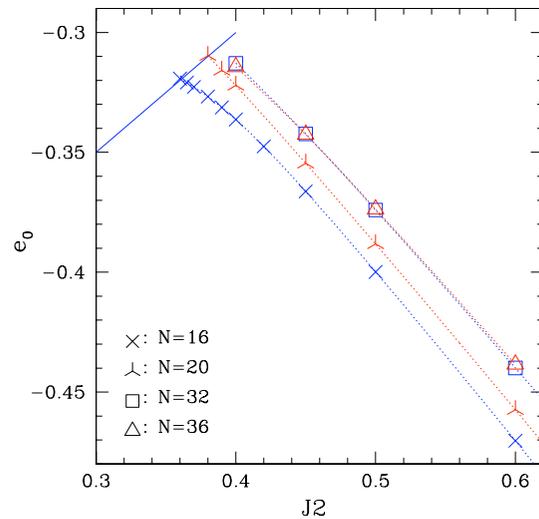
can we see the nematic in numerics ?

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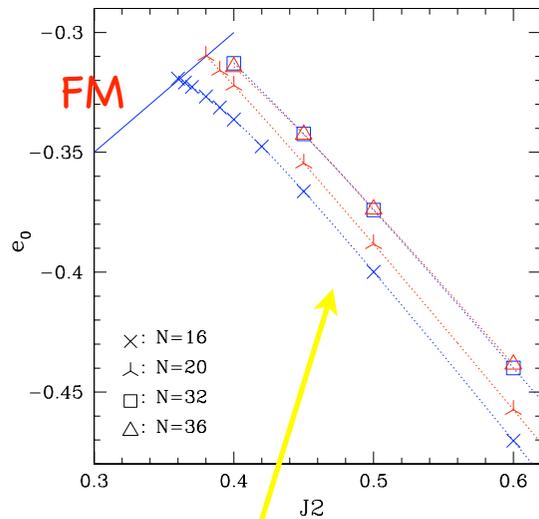
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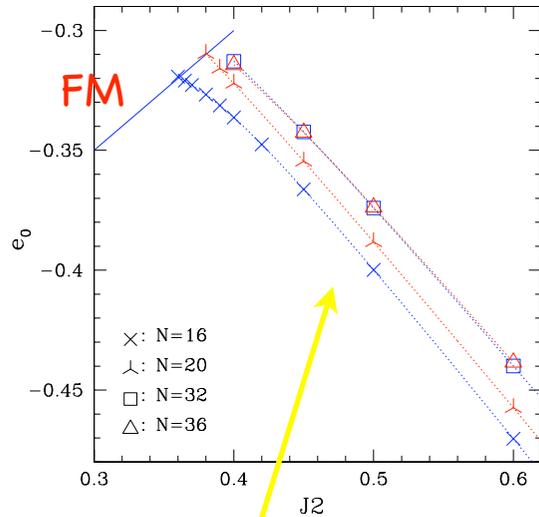


ground state energy
($S=0$) of finite size cluster

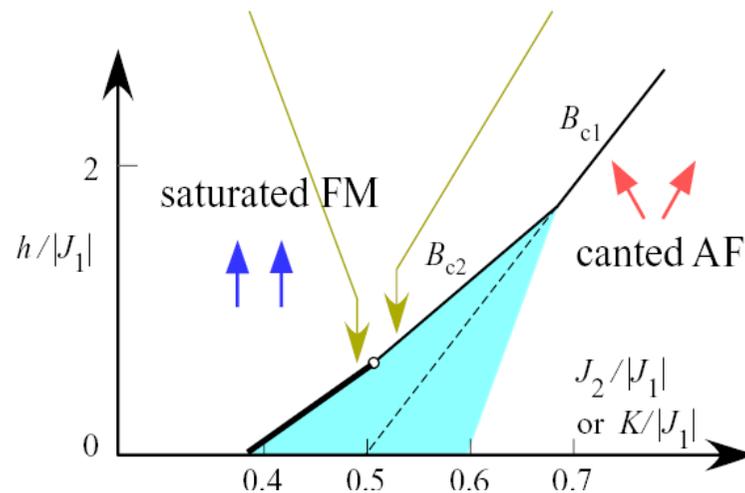
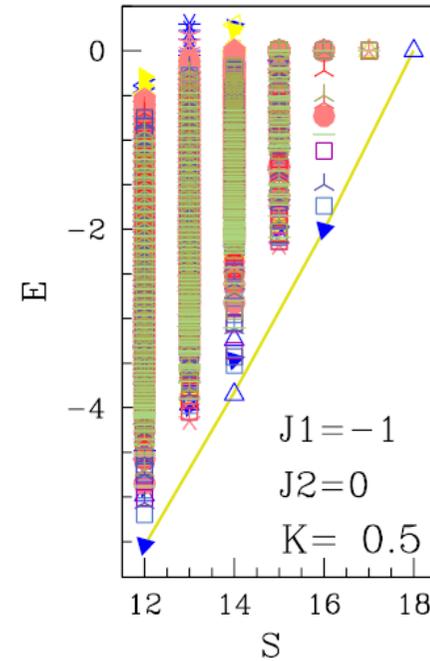
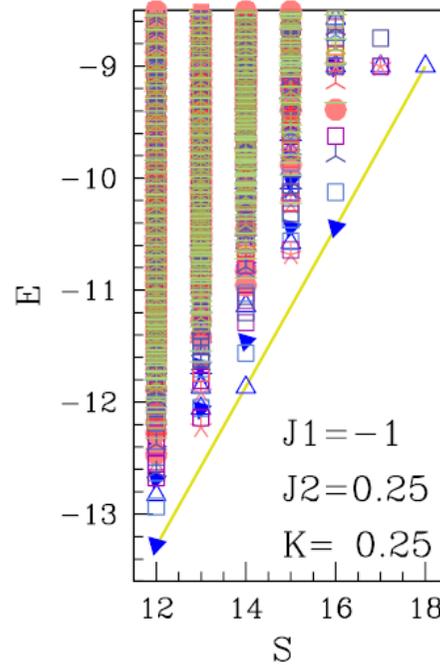
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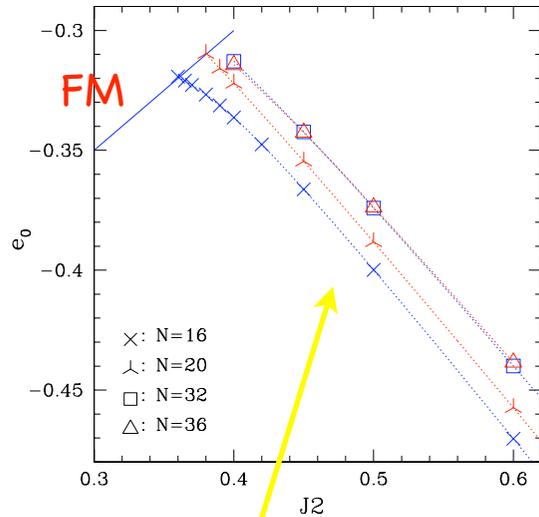
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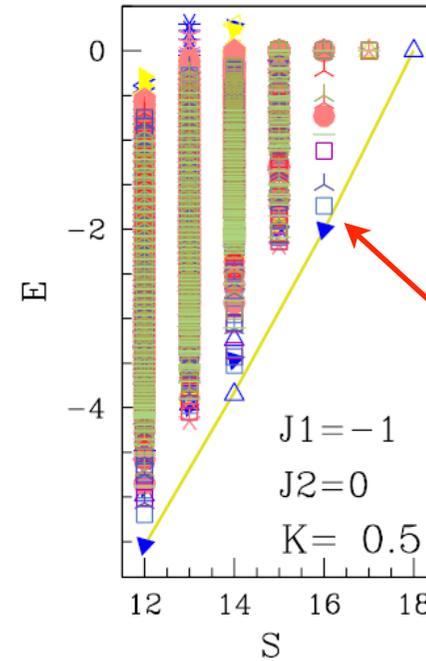
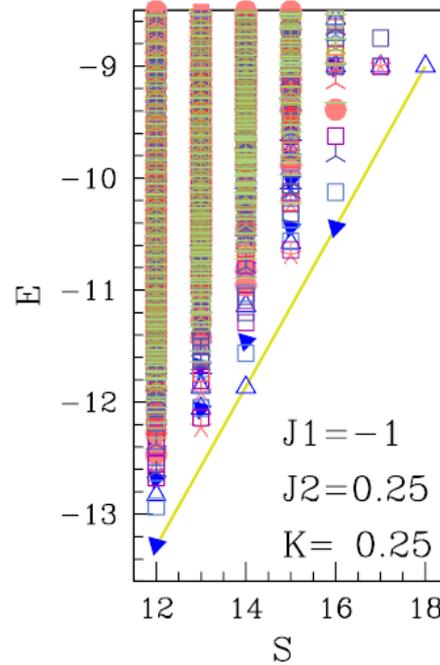
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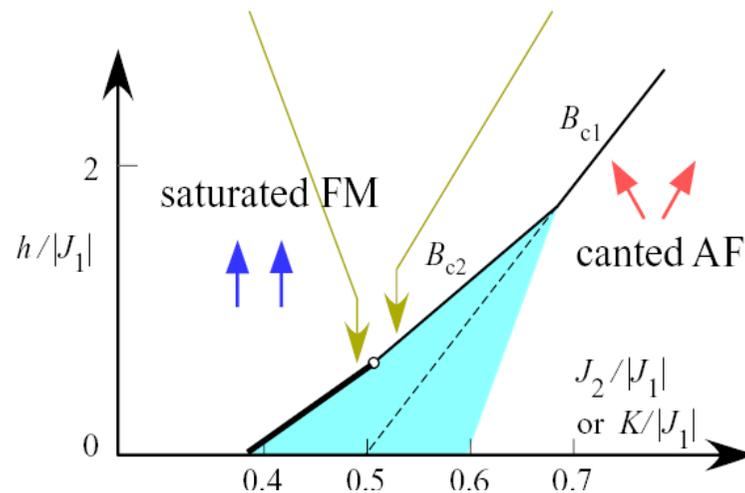
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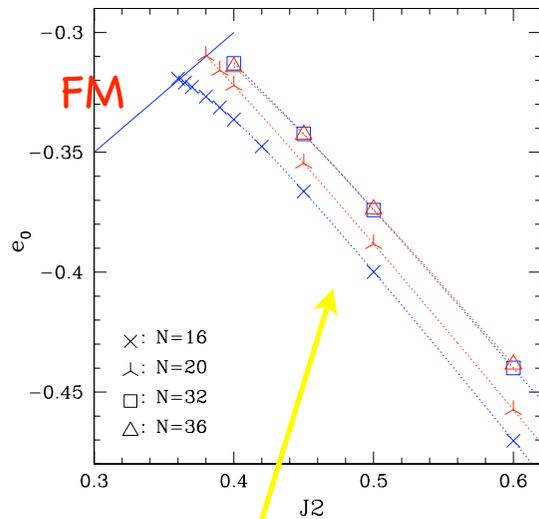
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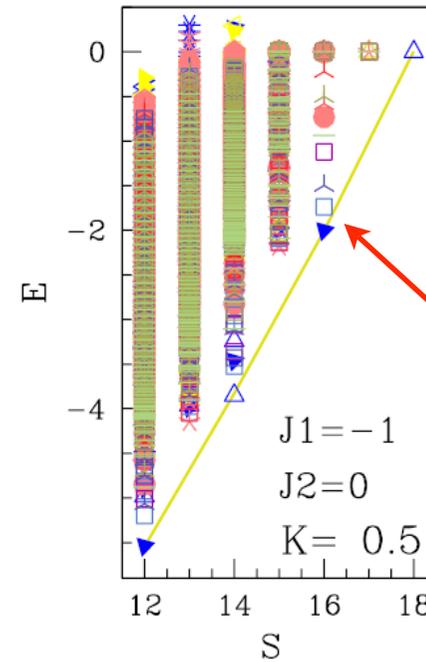
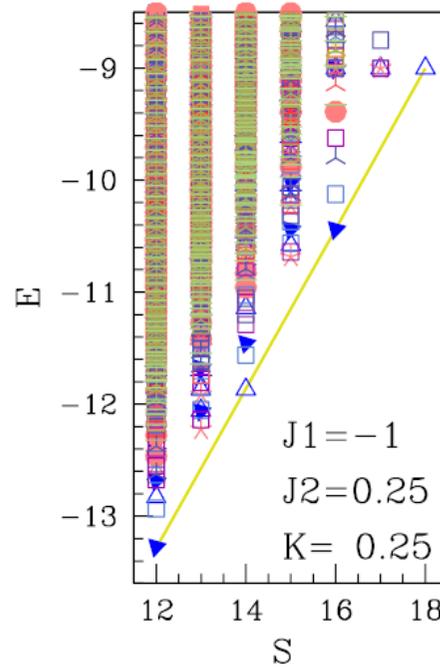
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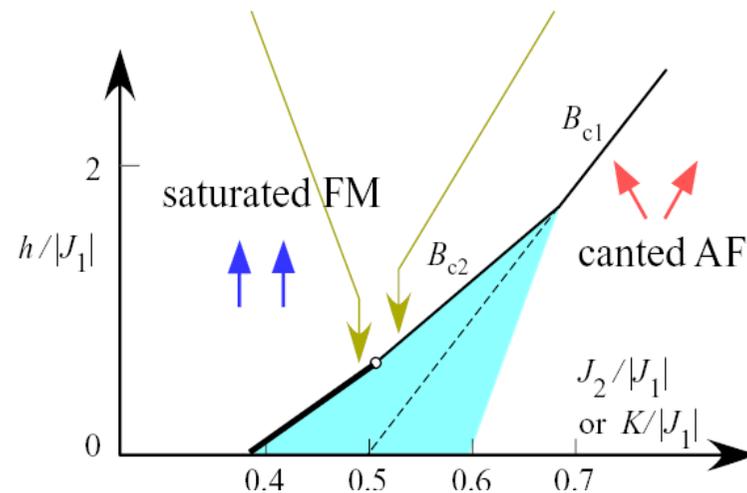


ground state energy ($S=0$) of finite size cluster



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Bose-Einstein condensation of bi-magnons at high field can be seen in numerics ✓



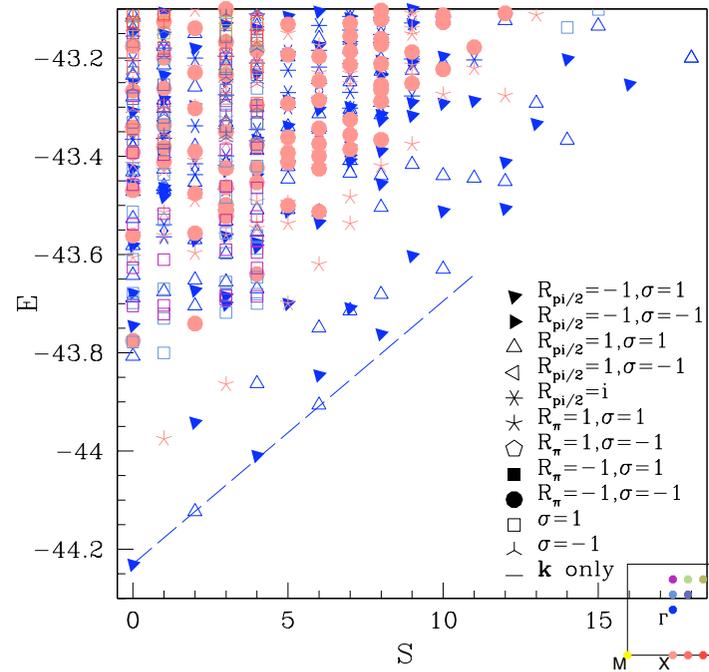
what about the ground state ?

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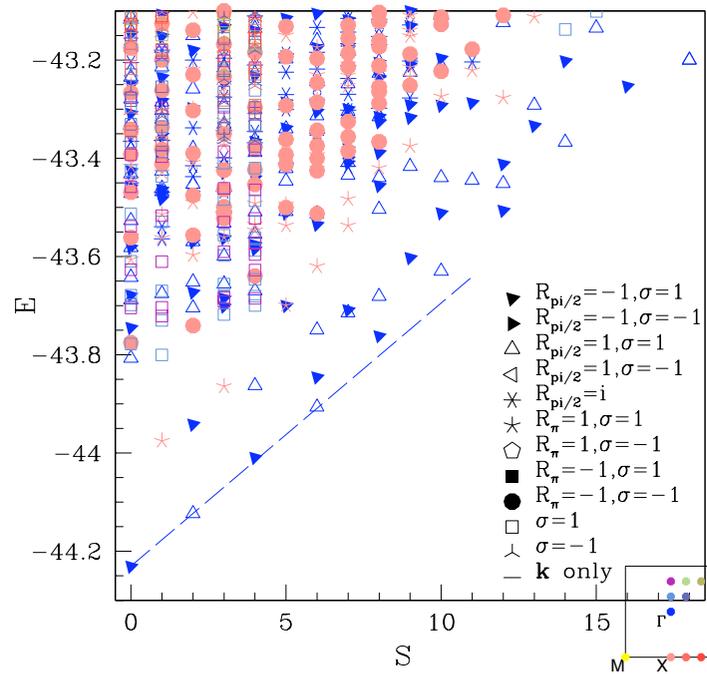
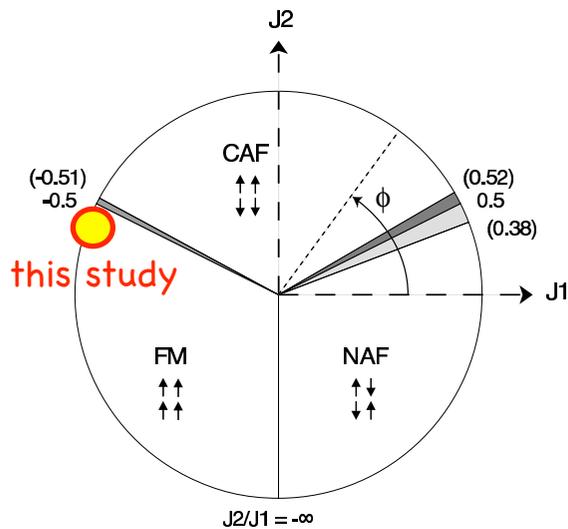
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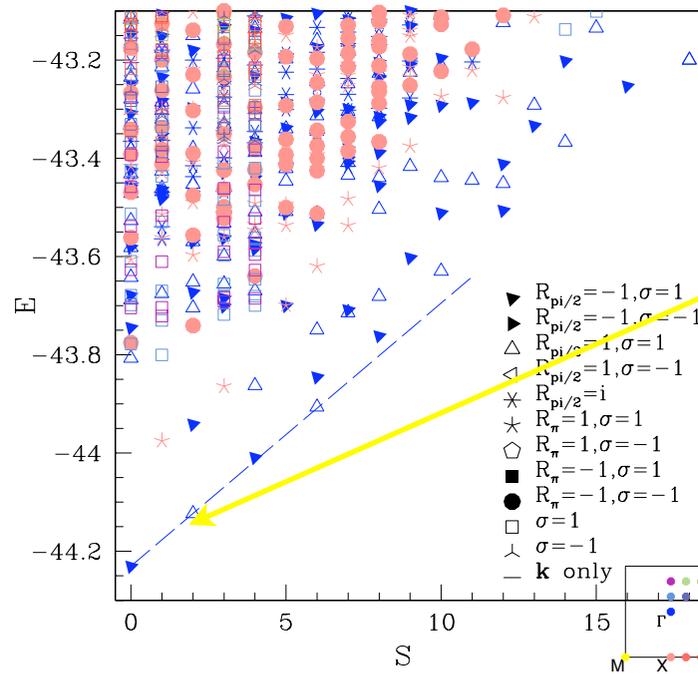
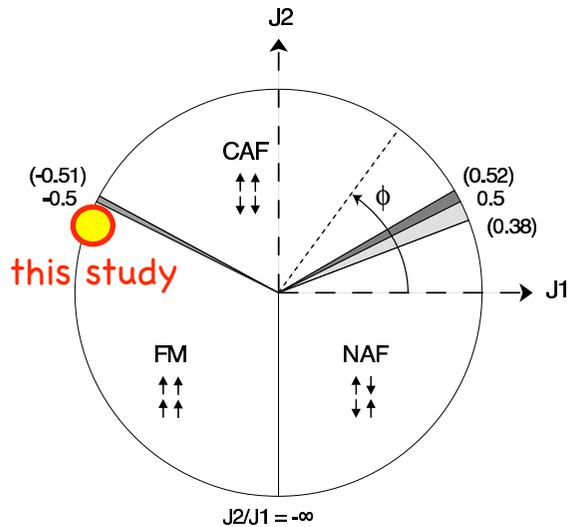
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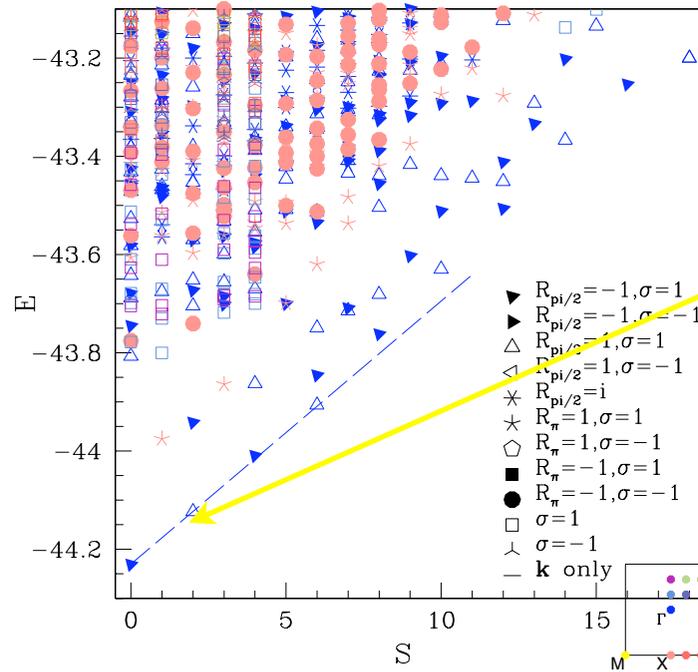
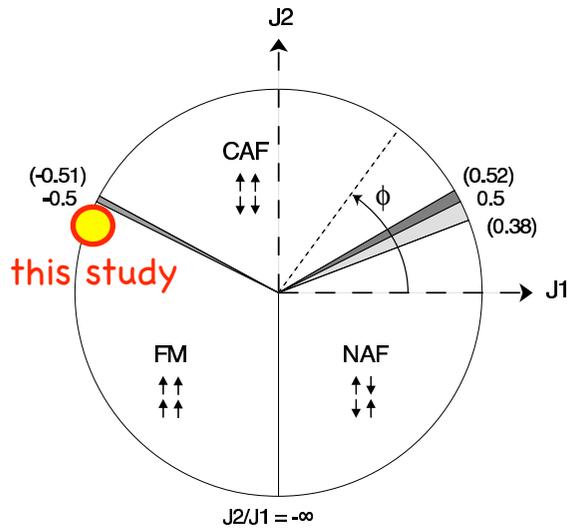


spectrum contains
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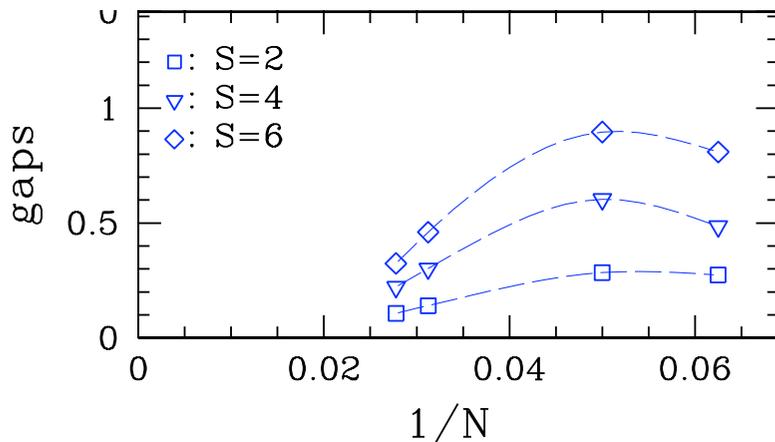
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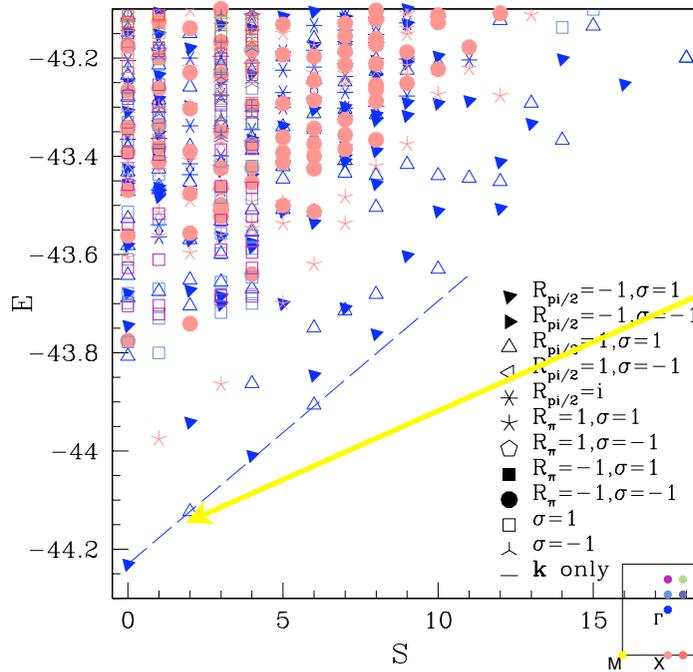
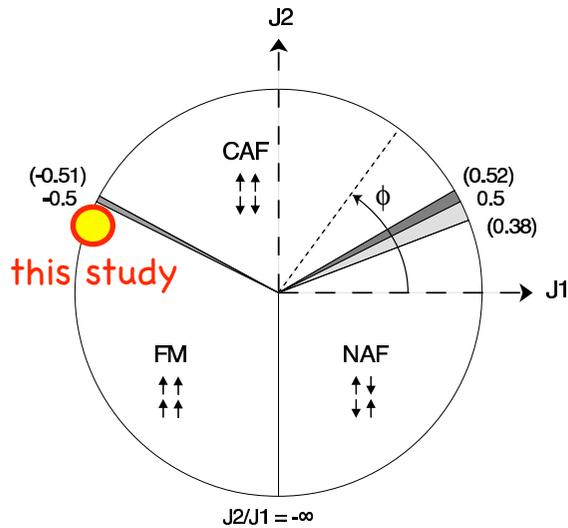
Gaps in even spin sectors scale to zero :



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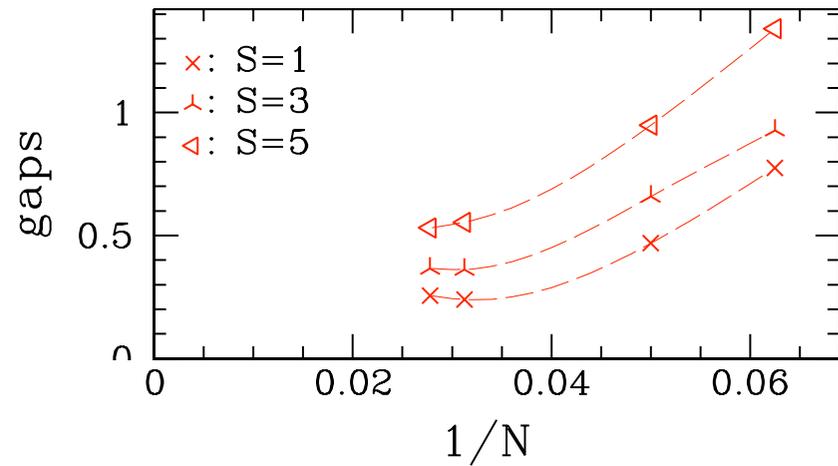
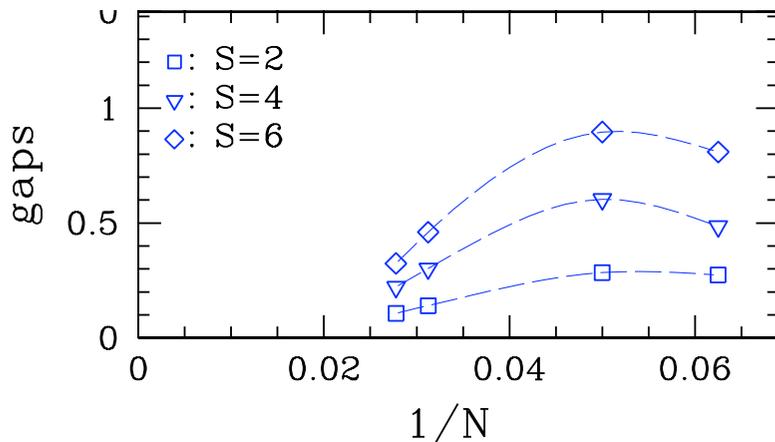
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Gaps in even spin sectors scale to zero :

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but is it really the same state ?

- nematic correlation in ground state -

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nematic correlation
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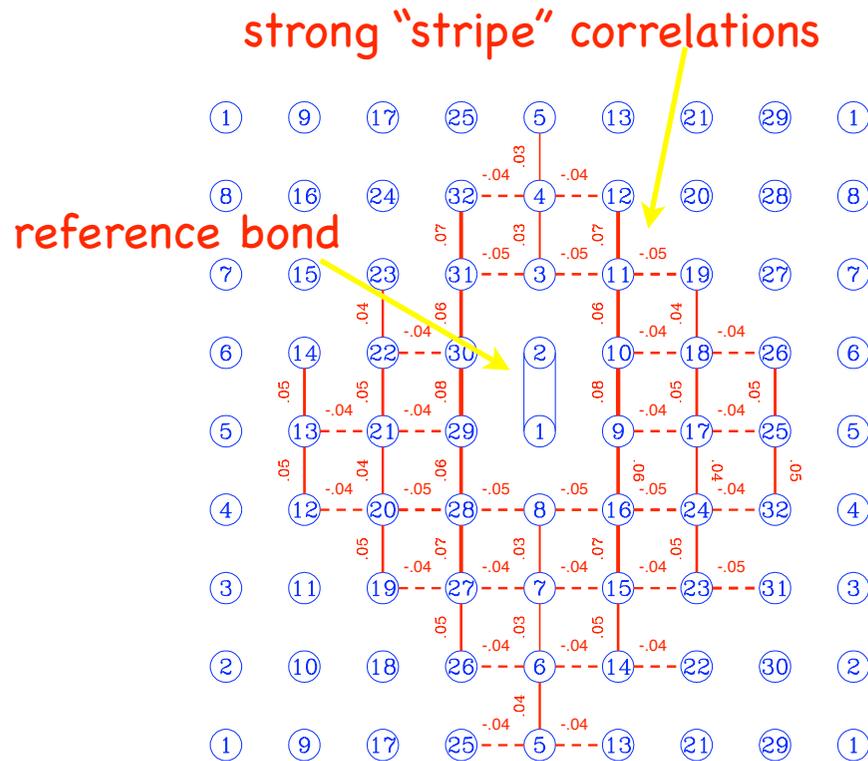
$$C(i, j, k, l) = \sum_{\alpha\beta} \langle \mathcal{O}^{\alpha\beta}(\mathbf{r}_i, \mathbf{r}_j) \mathcal{O}^{\alpha\beta}(\mathbf{r}_k, \mathbf{r}_l) \rangle$$

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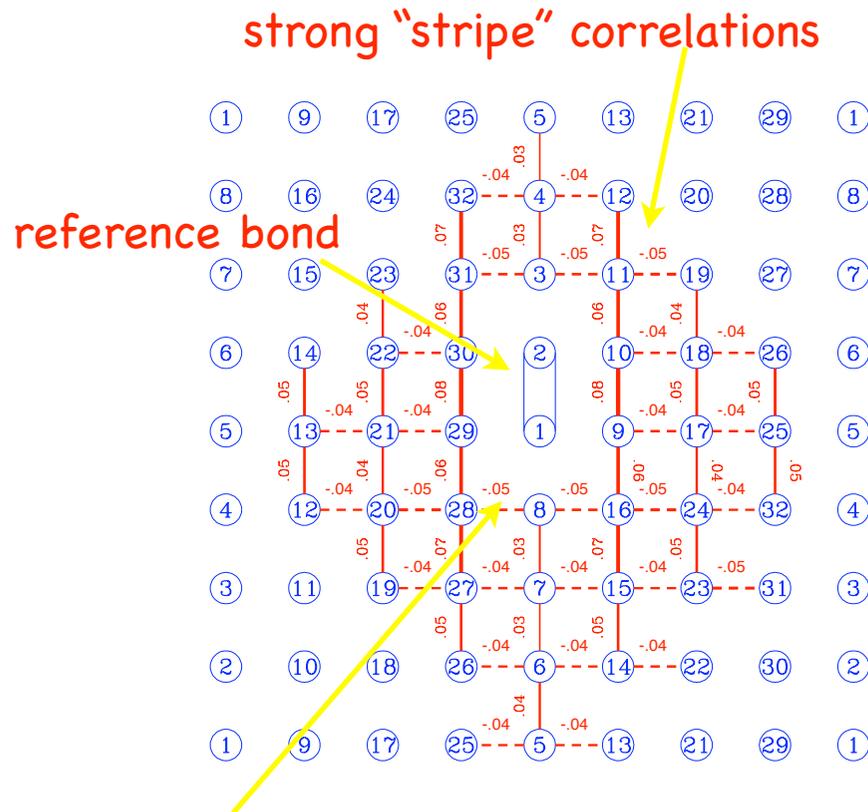


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"-" \Rightarrow d-wave sym

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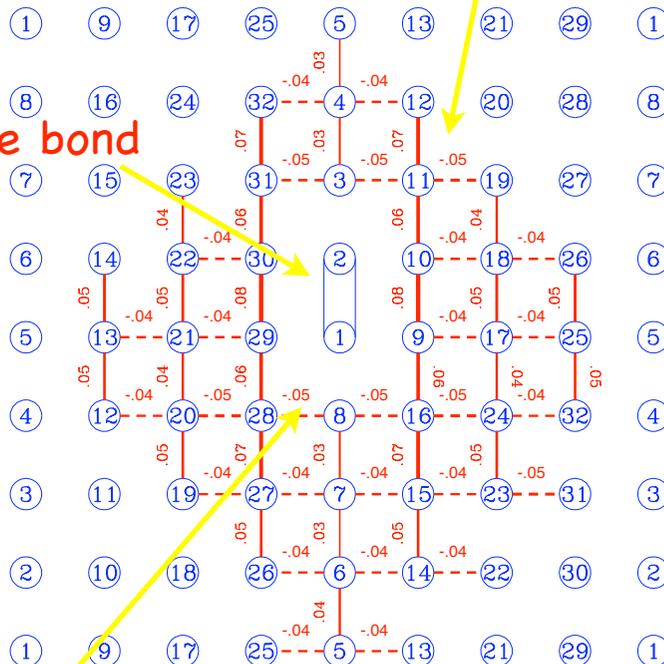
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strong "stripe" correlations

reference bond



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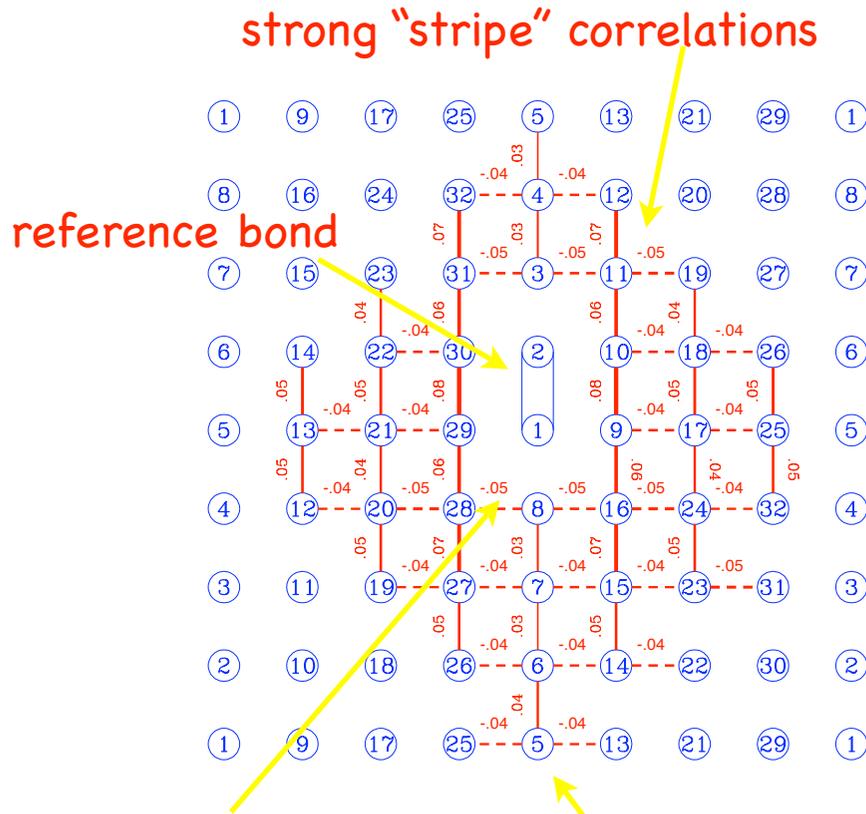
slow decay
of correlations

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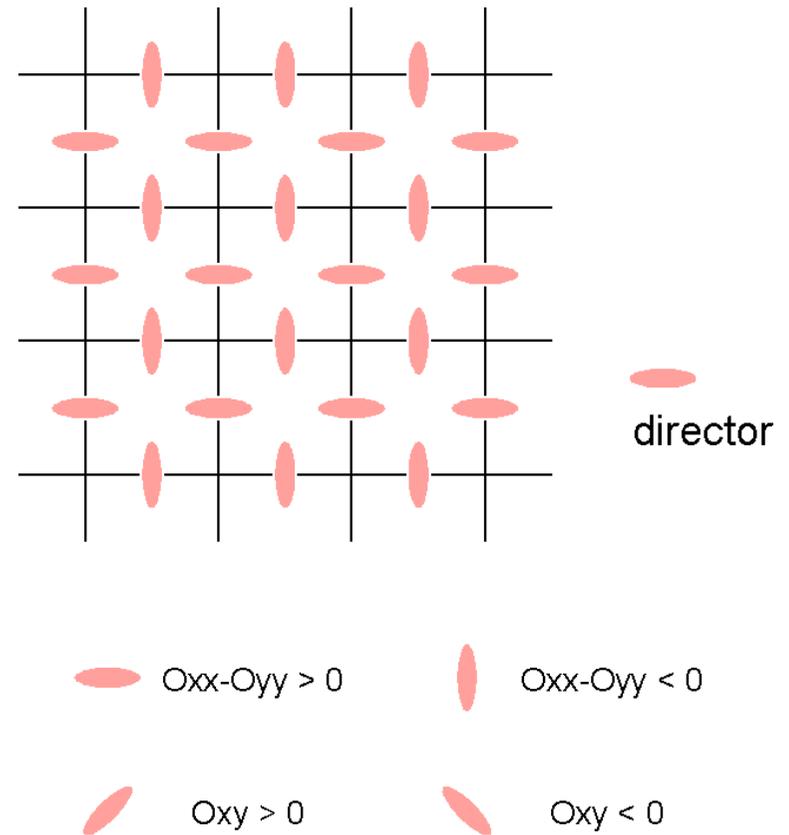
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slow decay
of correlations

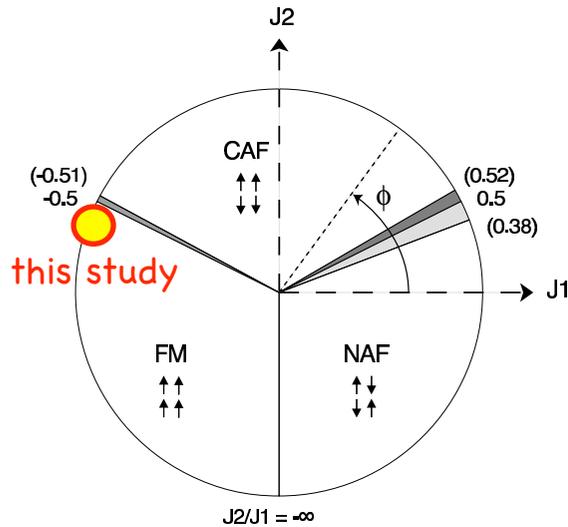


thermodynamics...

- formation of triplets and "emergent" nematic order -

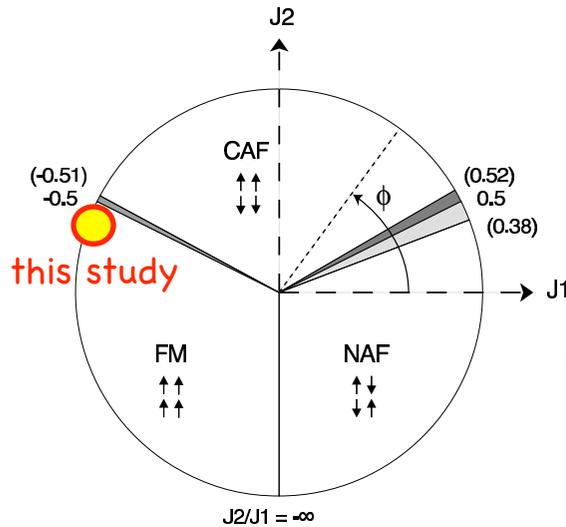
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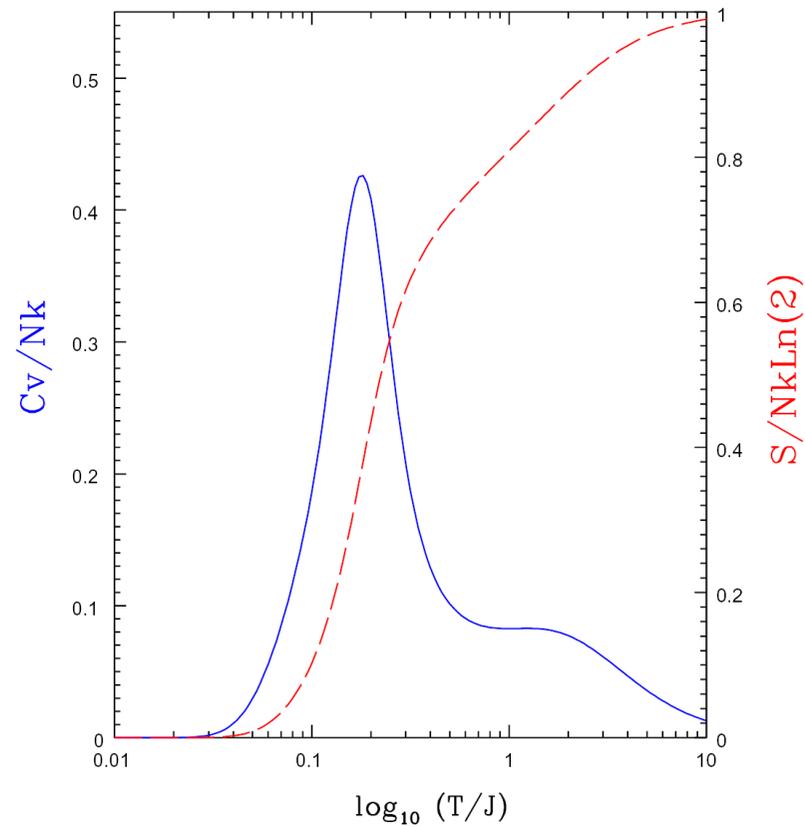


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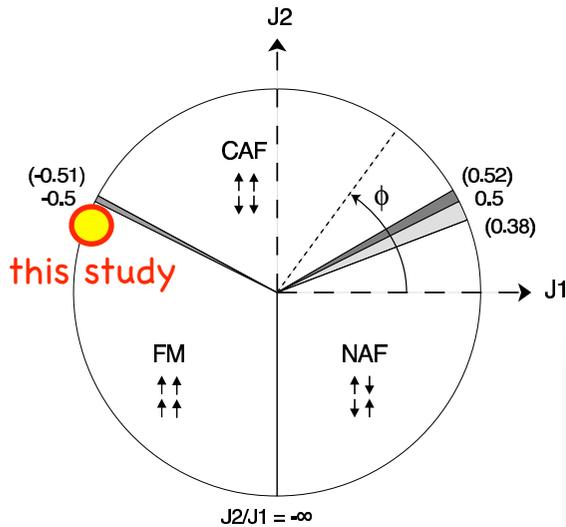


heat
capacity

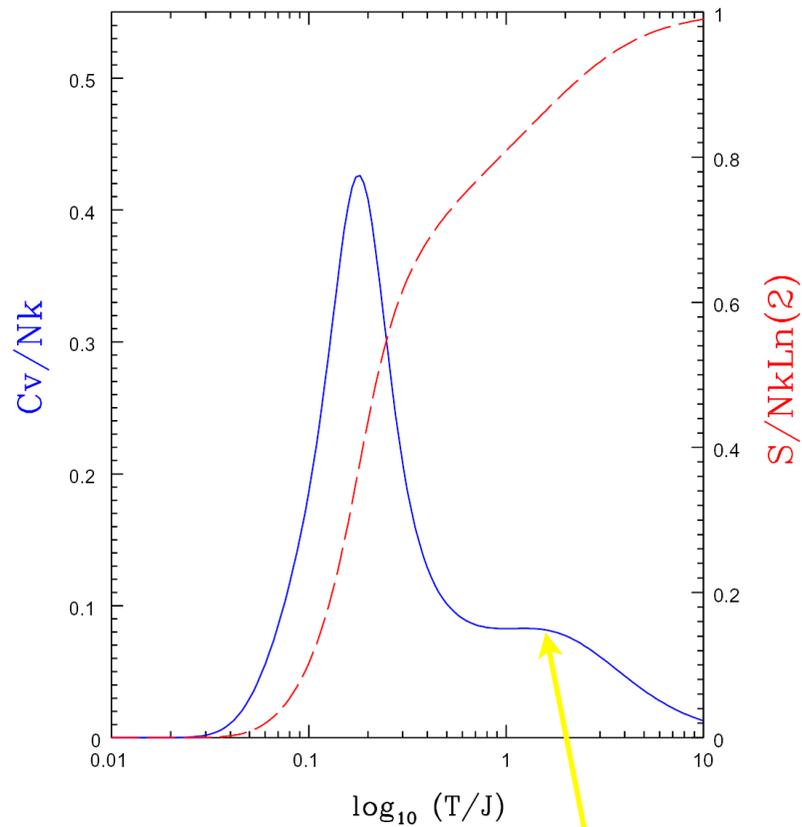


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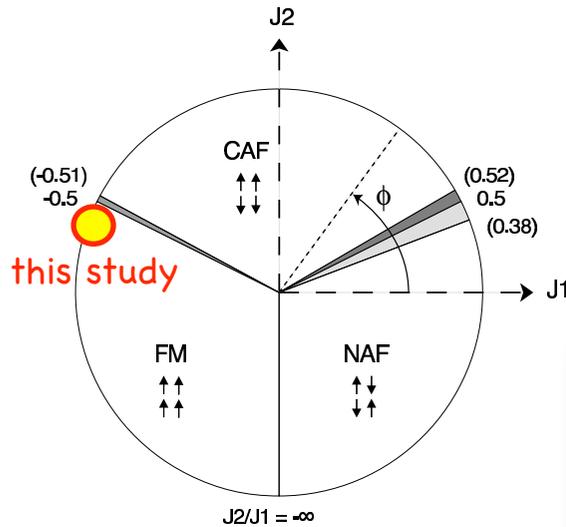
heat
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natural energy
scale $\sim J$

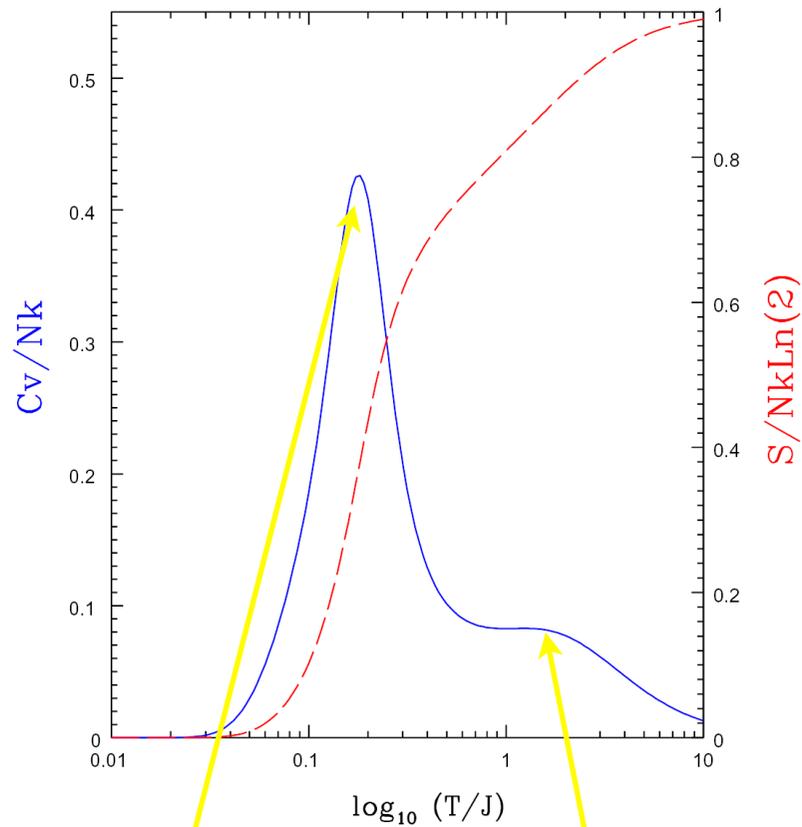
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this study

heat capacity

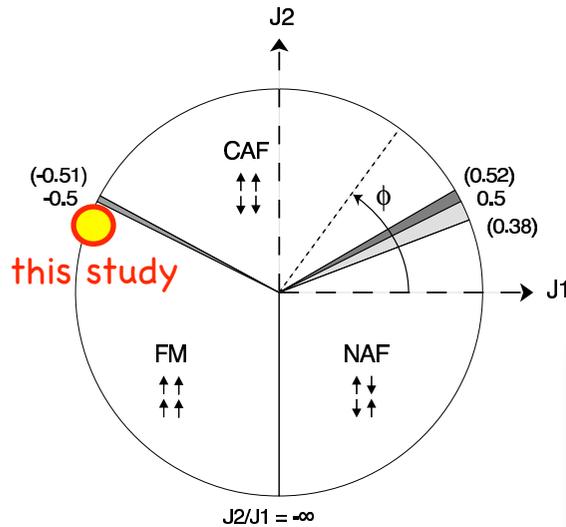


new low
"emergent"
energy scale

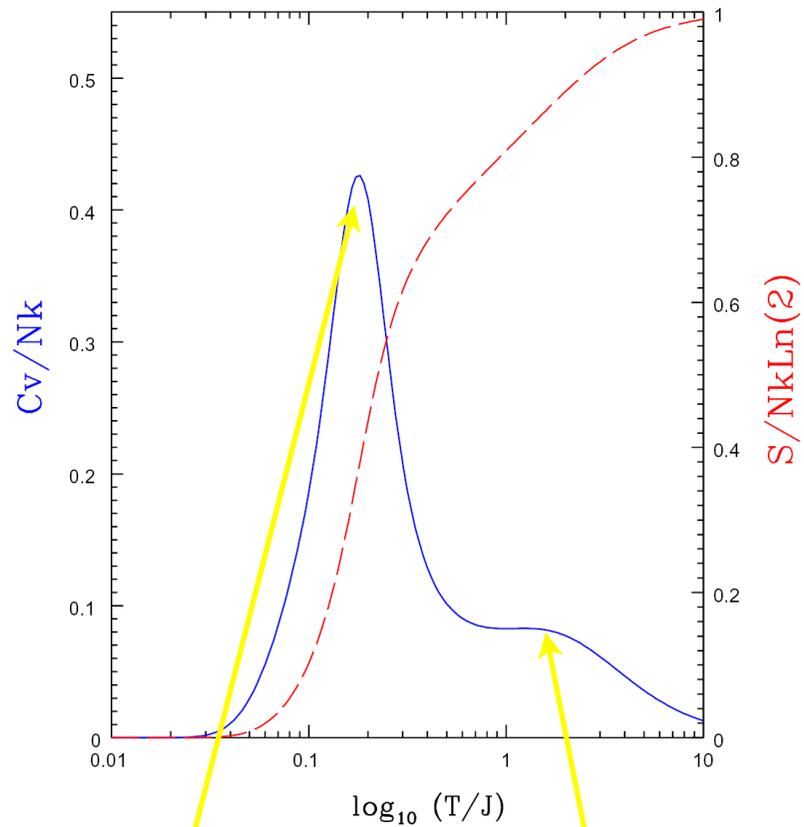
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heat capacity



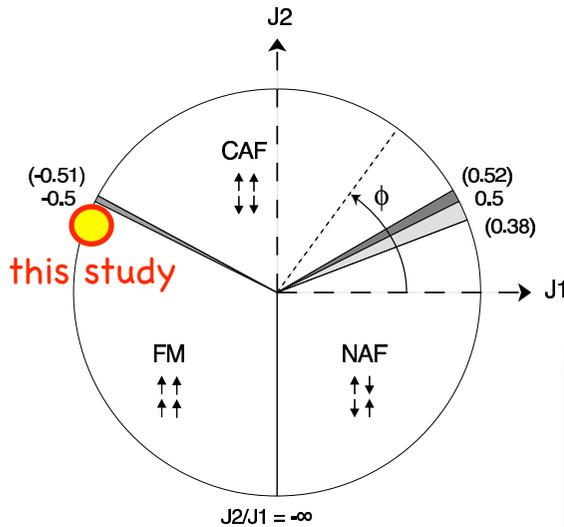
entropy

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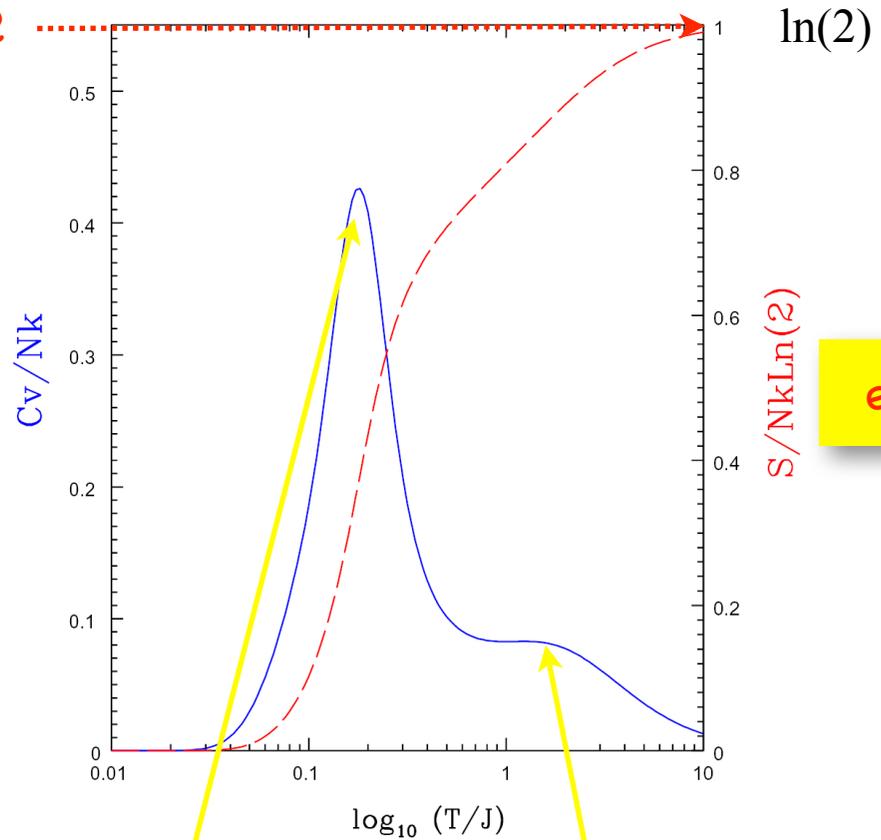
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heat capacity

spin-1/2

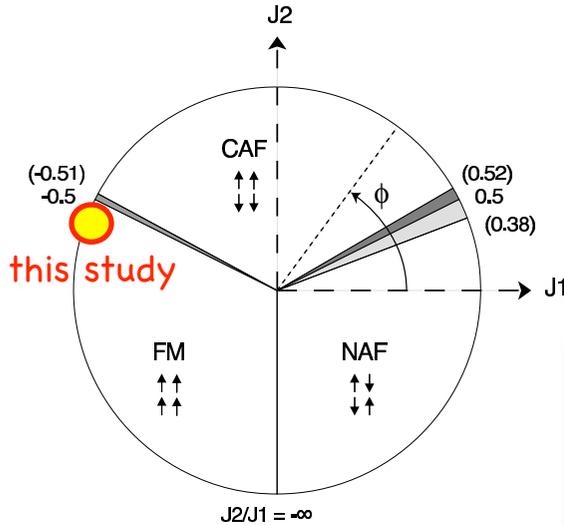


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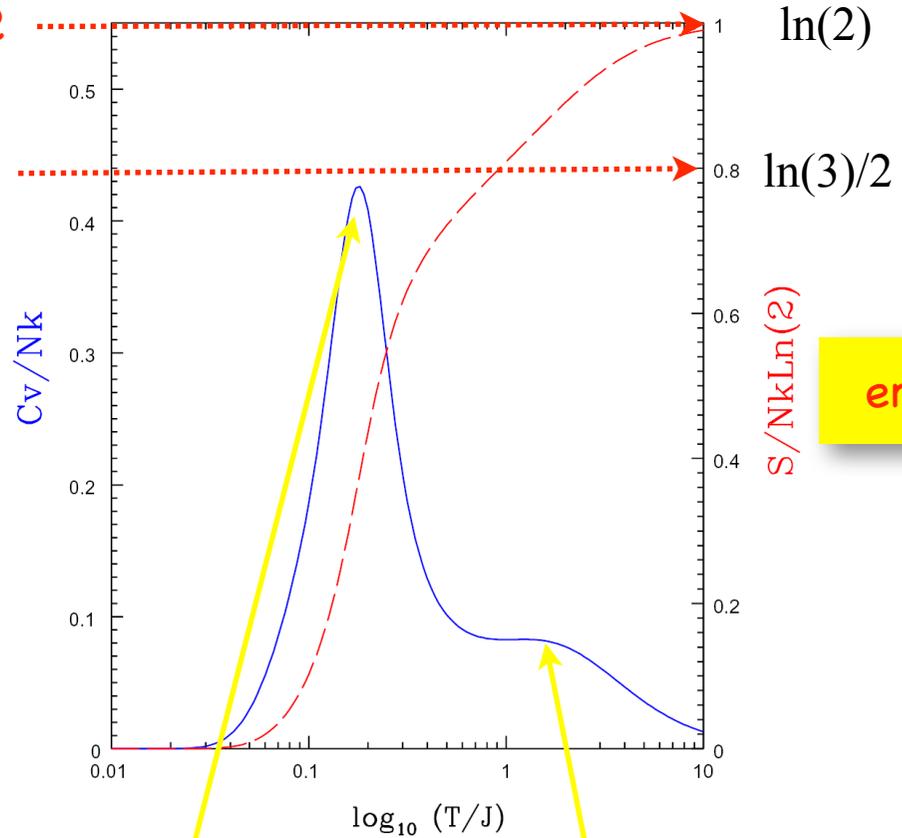
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spin-1/2
spin-1
on
bonds:

heat
capacity



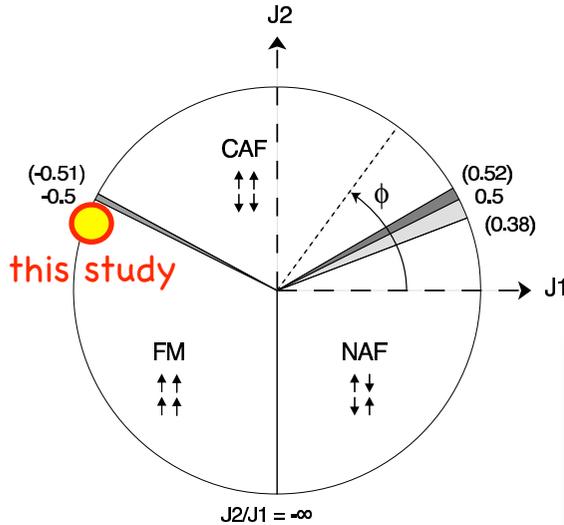
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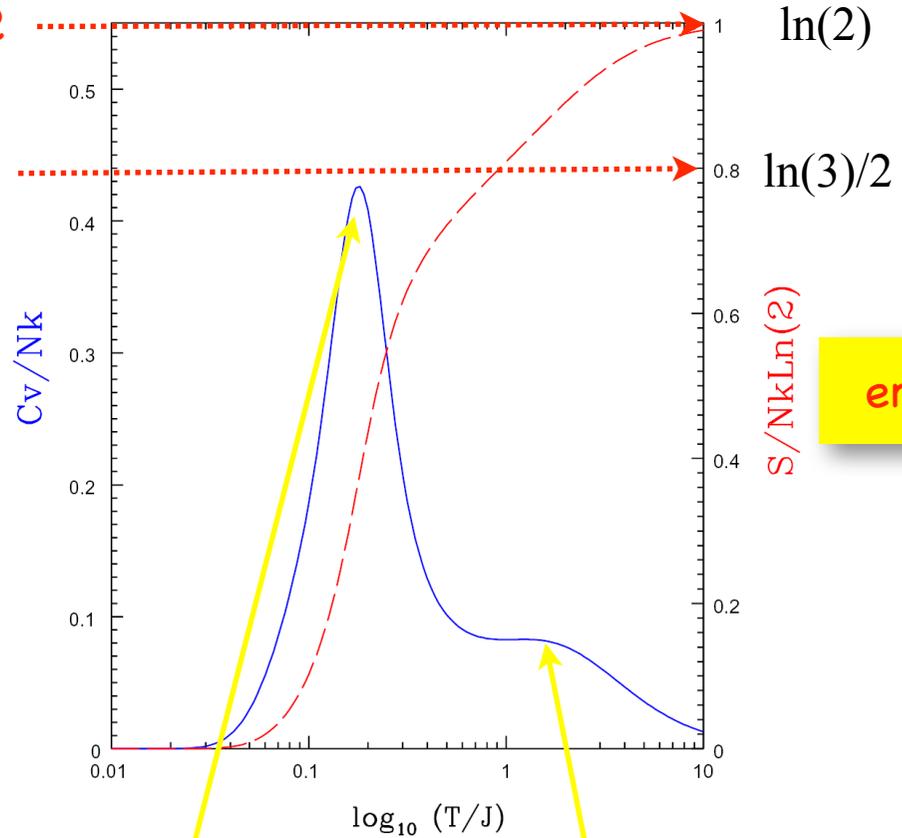
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spin-1/2
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on
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heat
capacity



entropy

entropy in peak is
associated with nematic
ordering of spin-1 bond
degrees of freedom

new low
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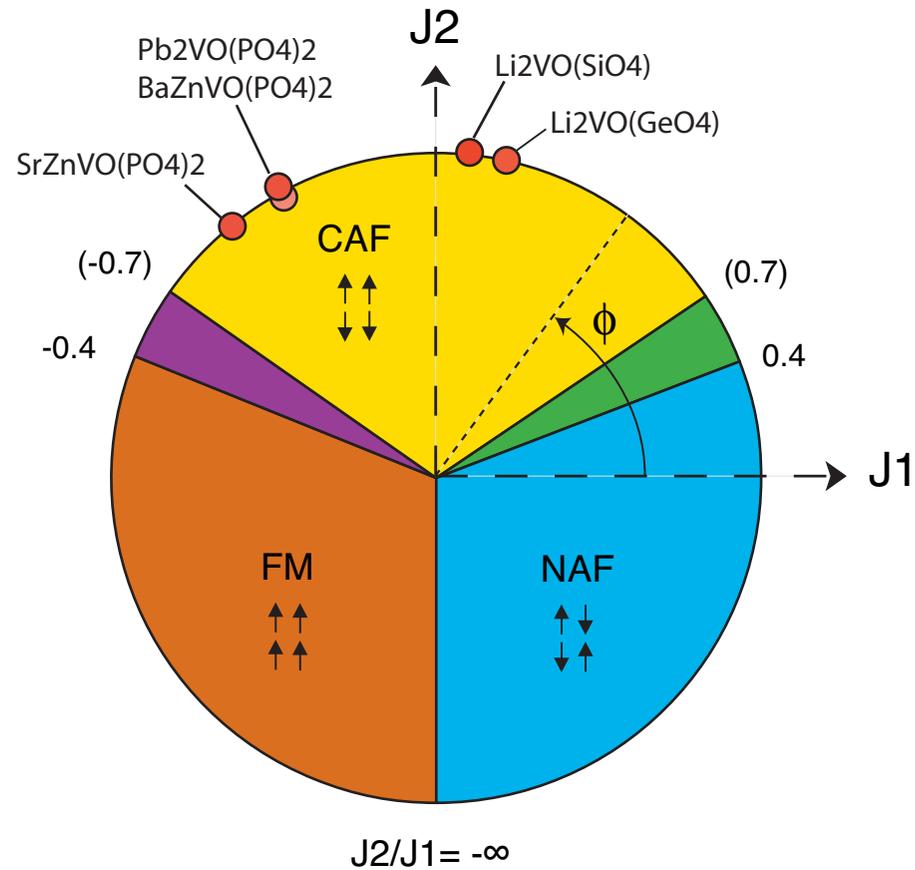
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could this ever be observed in nature ?

- more quasi-2D vanadates ! -

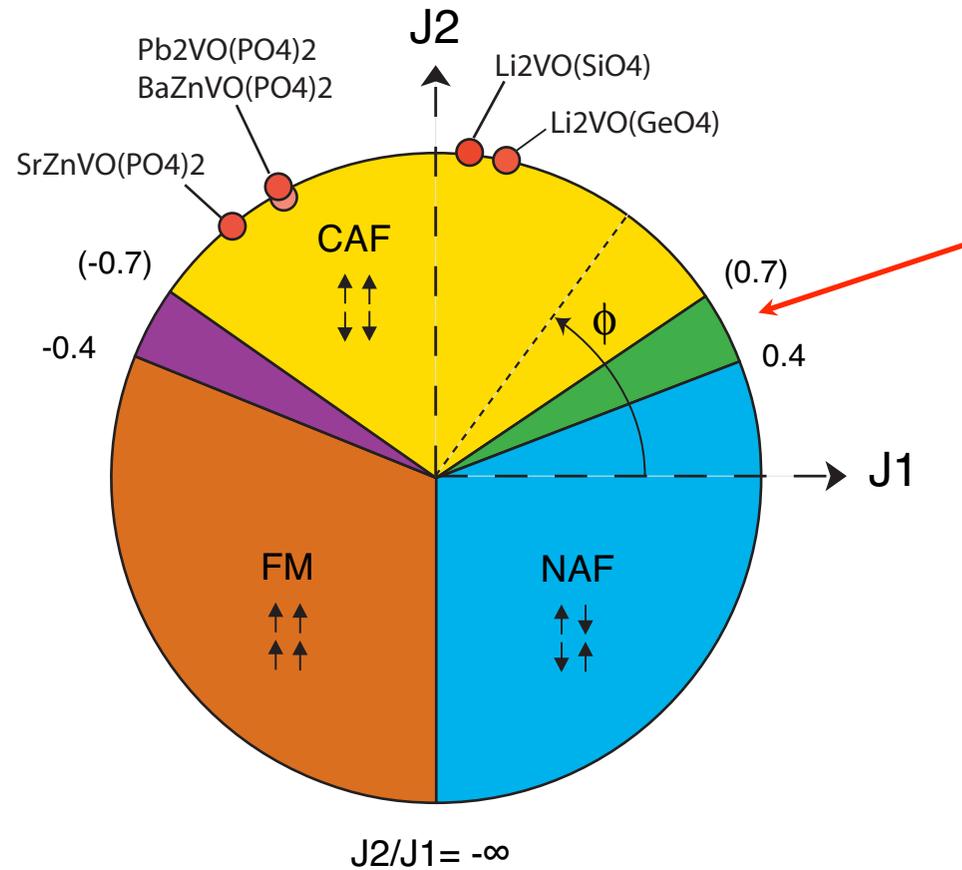
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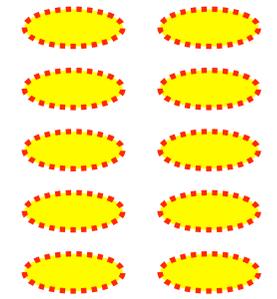


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***gapped,
staggered
dimer phase***
[P. Sindzingre,
Phys. Rev. B **69**,
094418 (2004),
and references
therein]



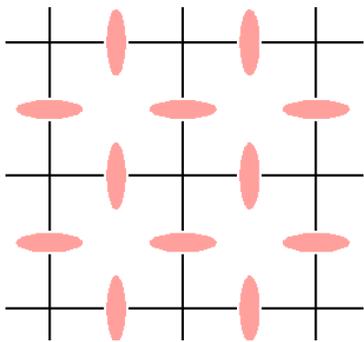
○ singlet

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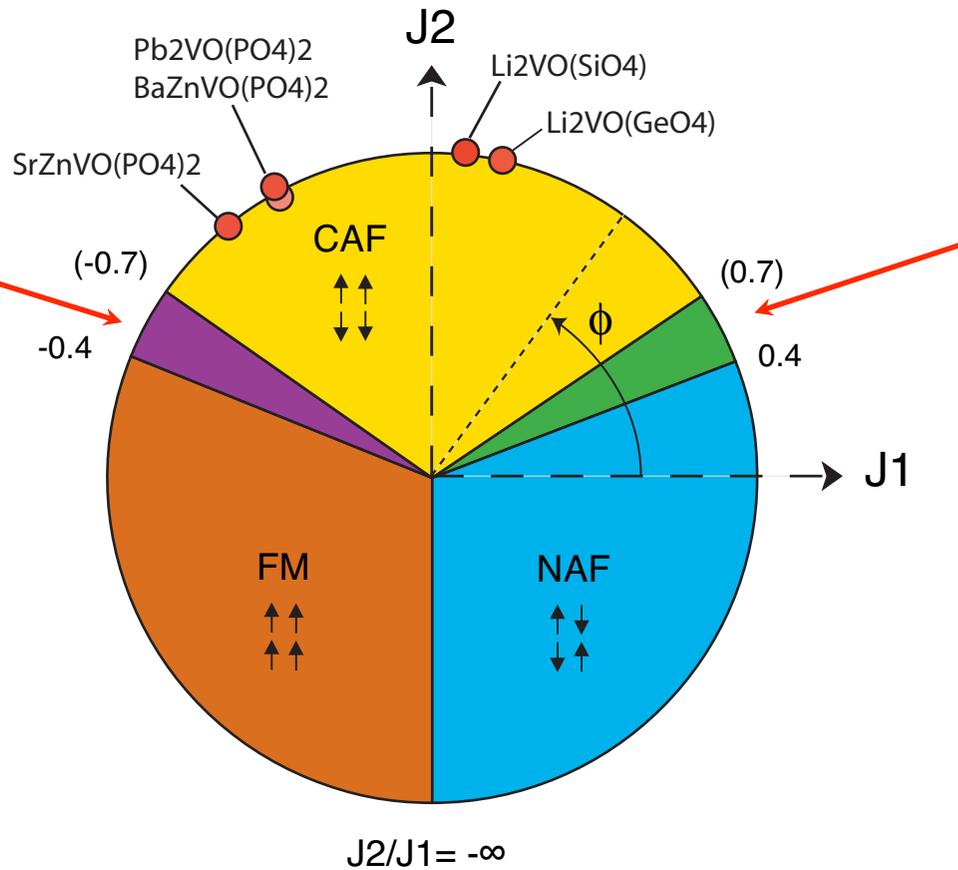
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gapless nematic phase

[N. Shannon *et al.*,
Phys. Rev. Lett. 96, 027213
(2006)]

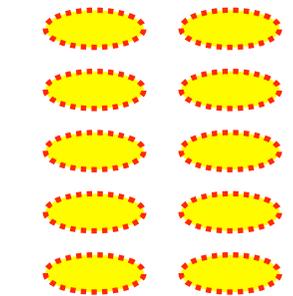


director



gapped, staggered dimer phase

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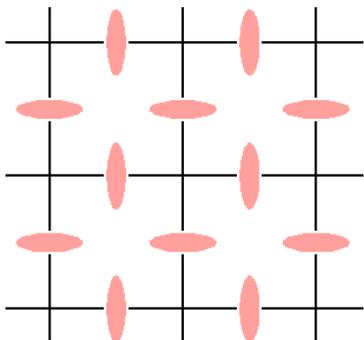
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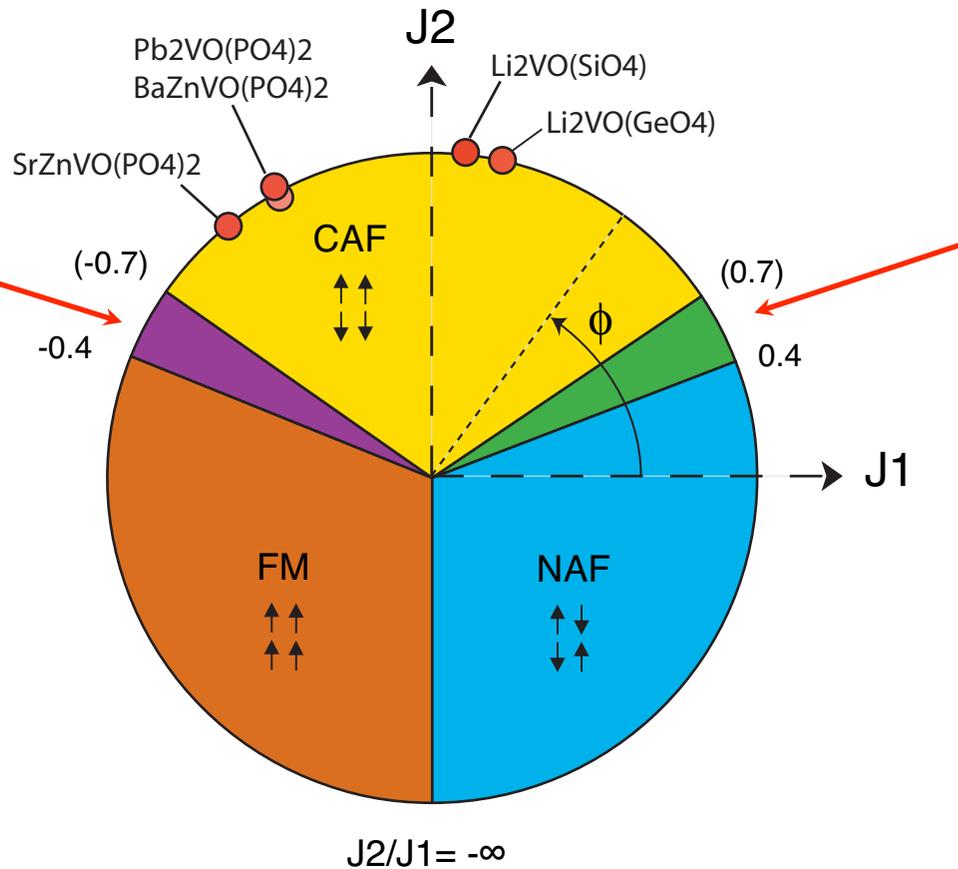
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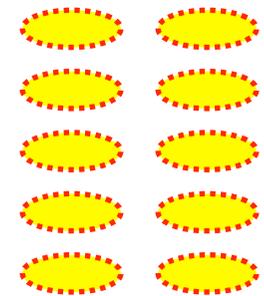


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singlet

new compound CaZnVO(PO4)2 looks promising...

so what happens on a triangular lattice ?

- modeling solid 2D films of He III -

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minimal model :

$$\mathcal{H} = 2J_1 \sum_{\langle ij \rangle} \mathbf{S}_i \mathbf{S}_j + K \sum_{\langle 1234 \rangle} P_{1234} + P_{1234}^{-1}$$

FM $J_1 < 0$
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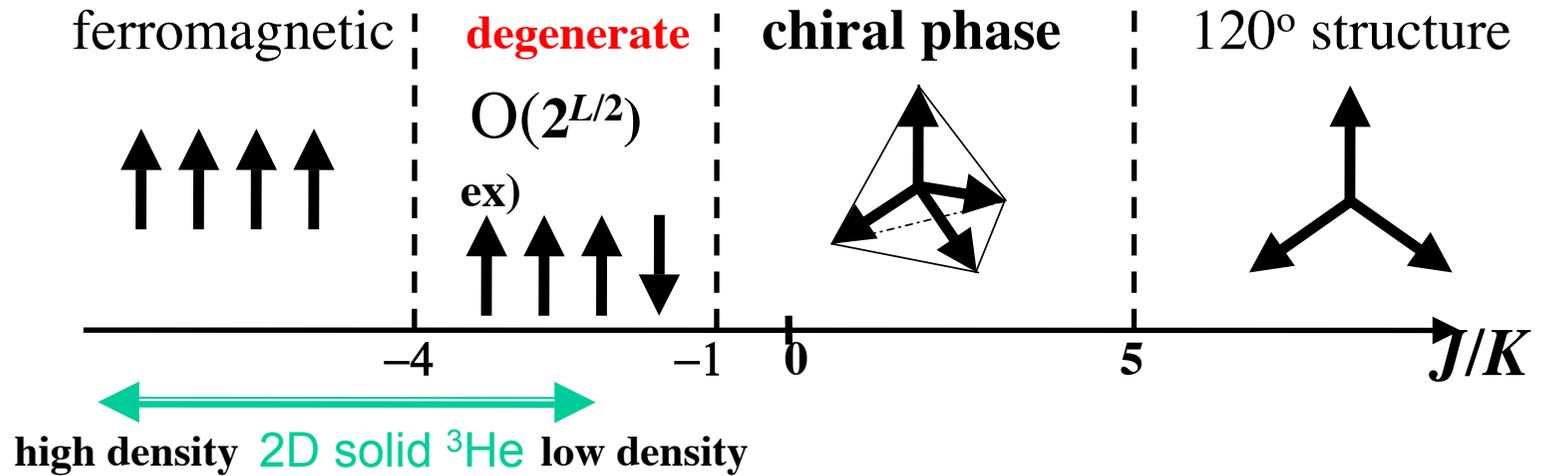
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[T. Momoi and
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PRB **65**,
052403(2002)]

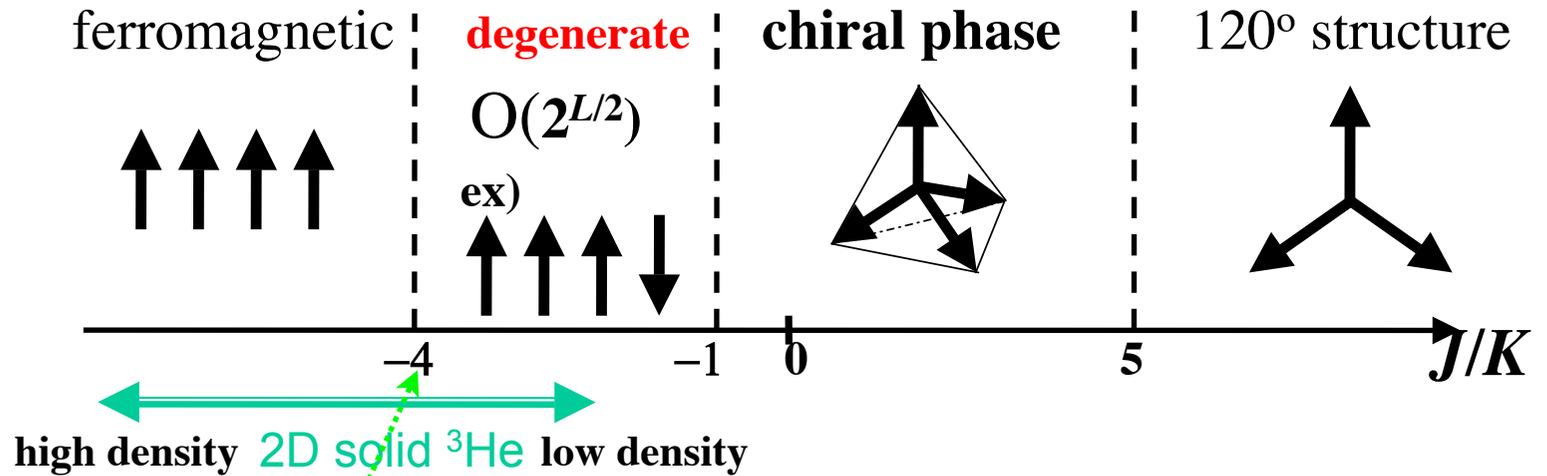
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 at which entire spin wave
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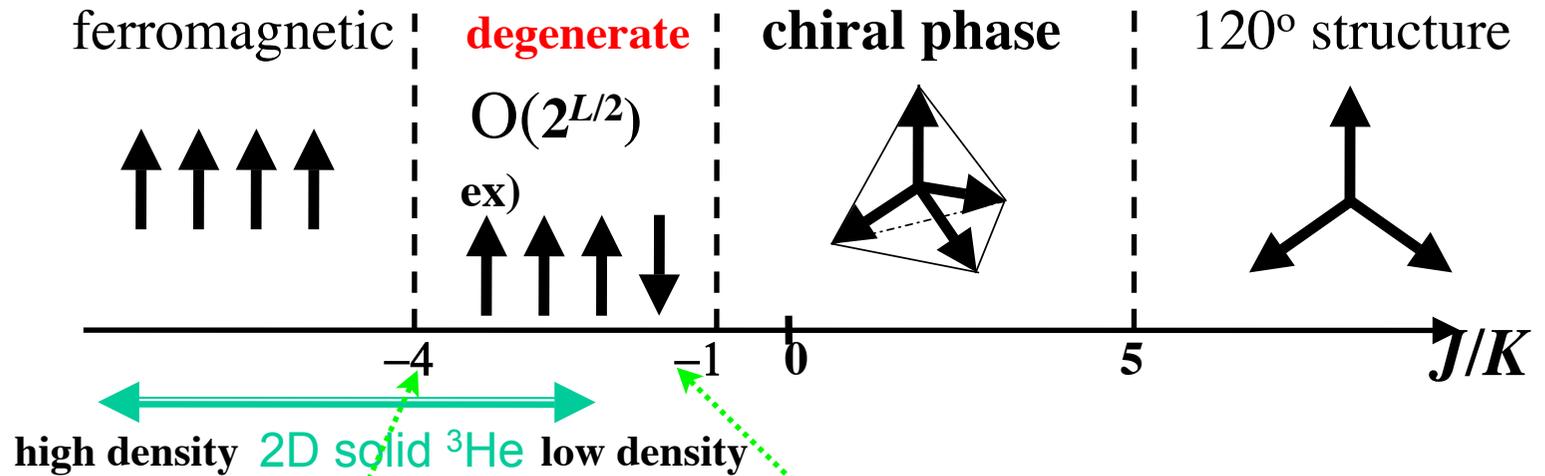
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[Misguich et al., PRL **81**, 1098 (1998)]

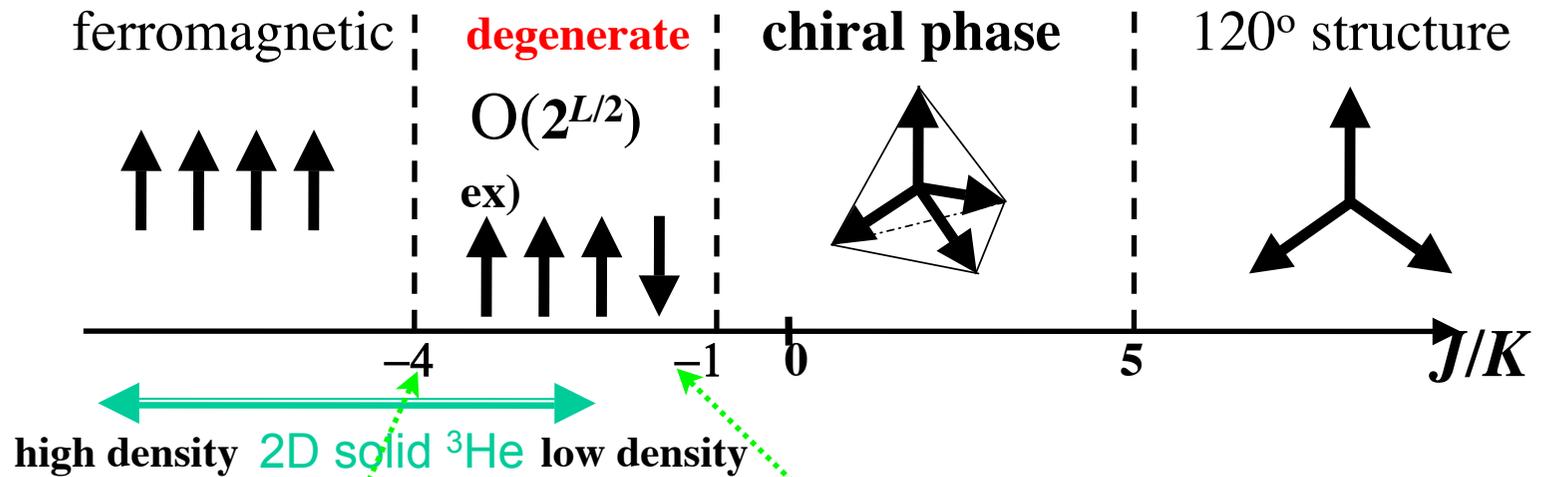
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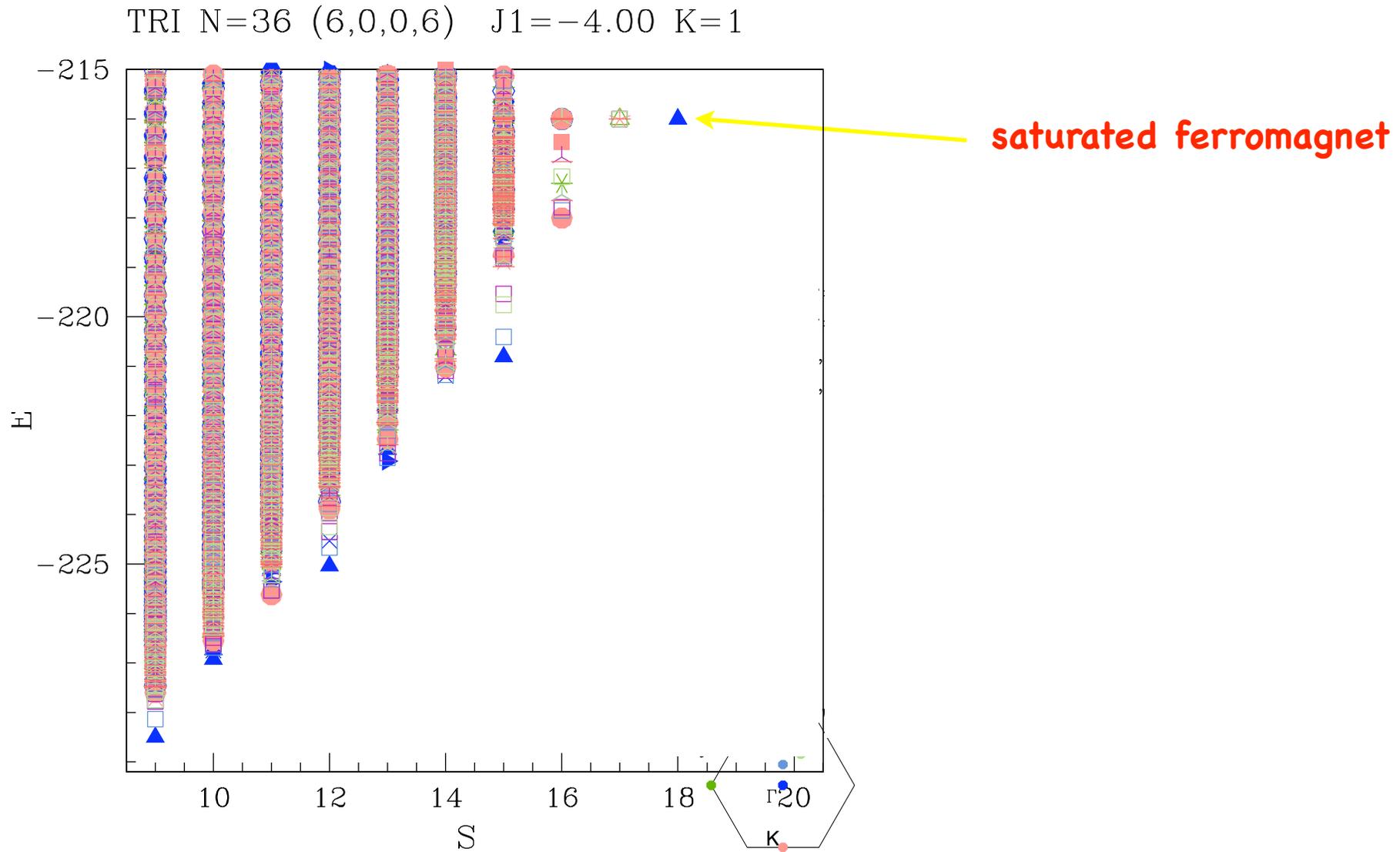
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experiment sees a **gapless** spin liquid bordering on the FM
- could this be another **nematic** state ?

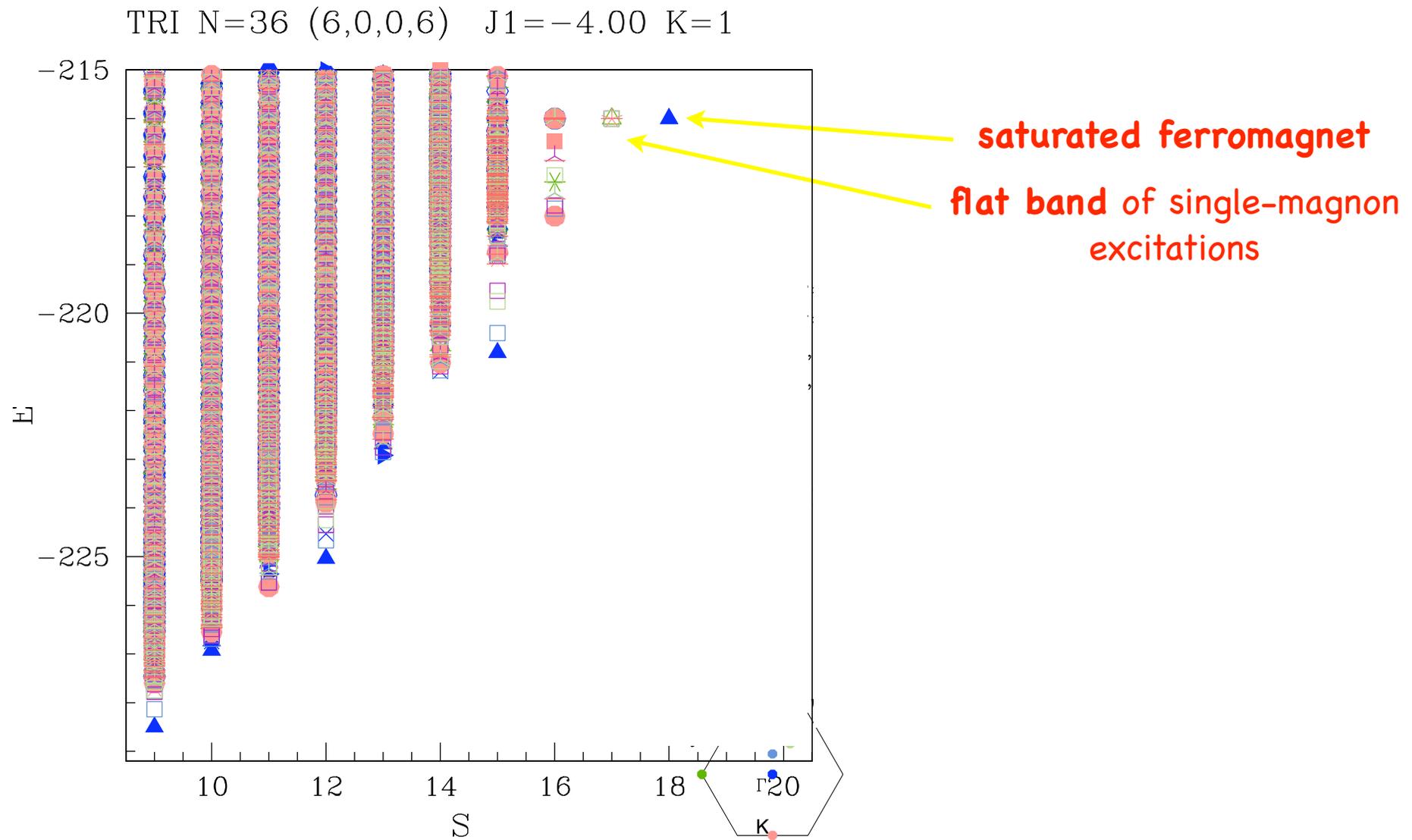
who ordered that ?

- going beyond nematic structure -



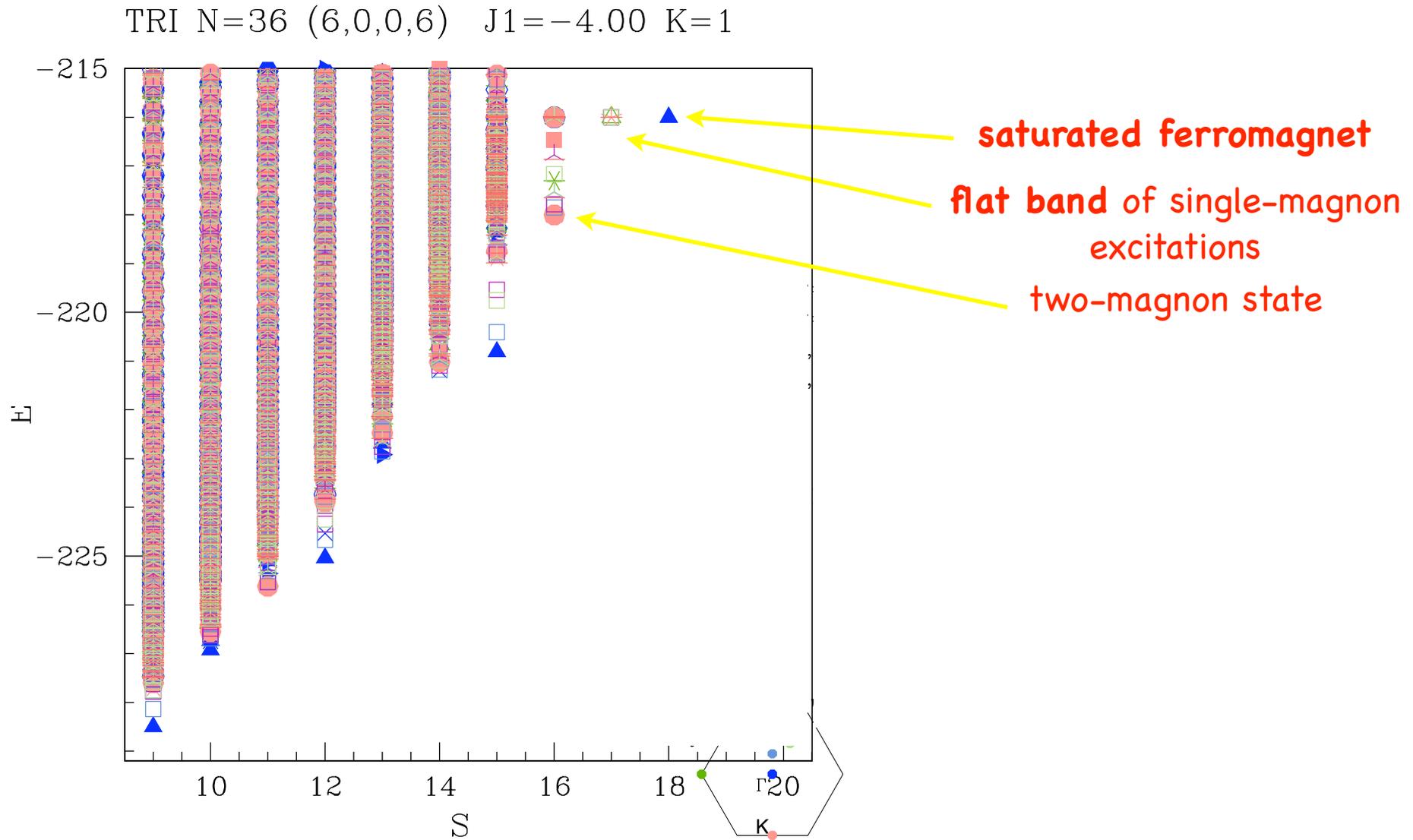
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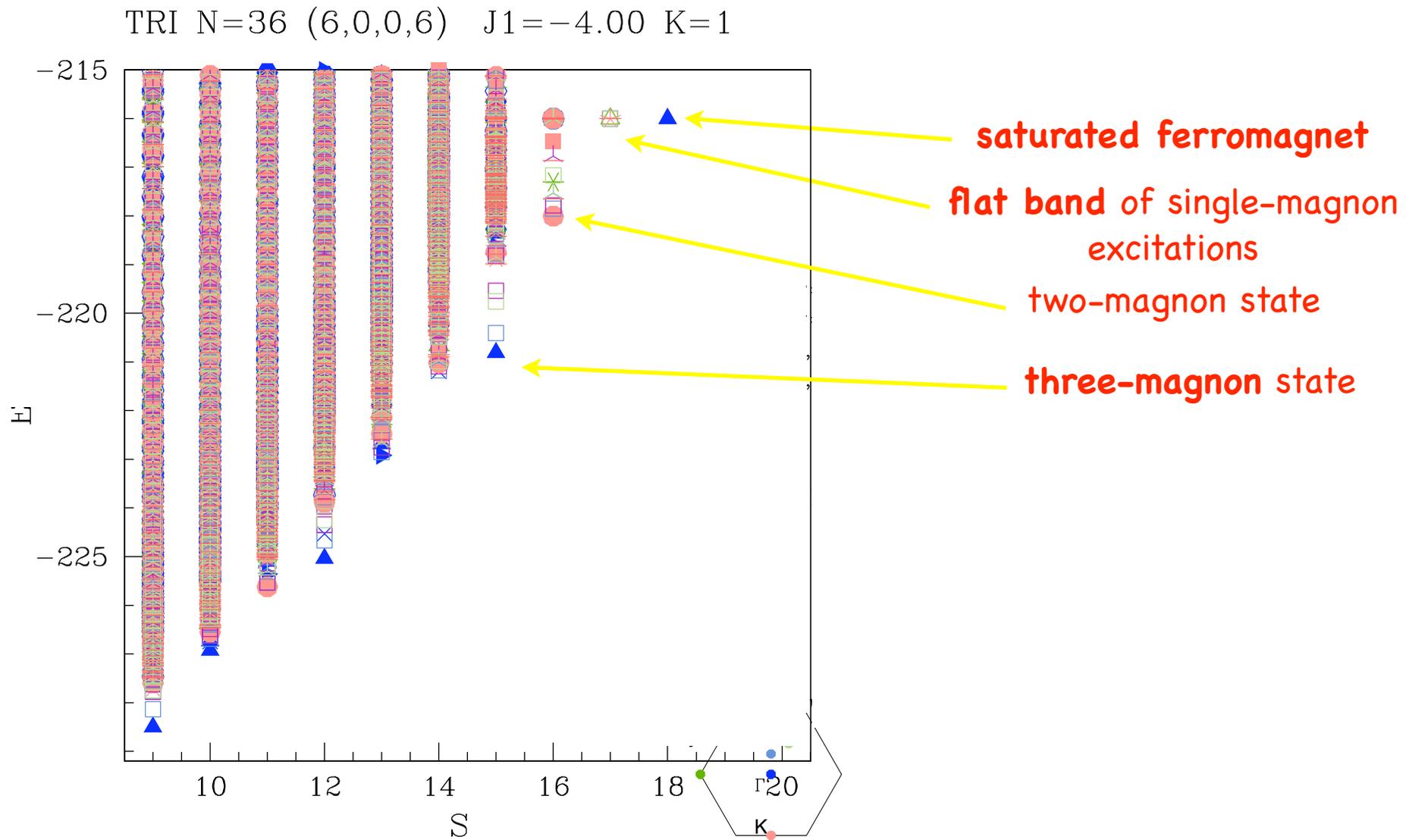
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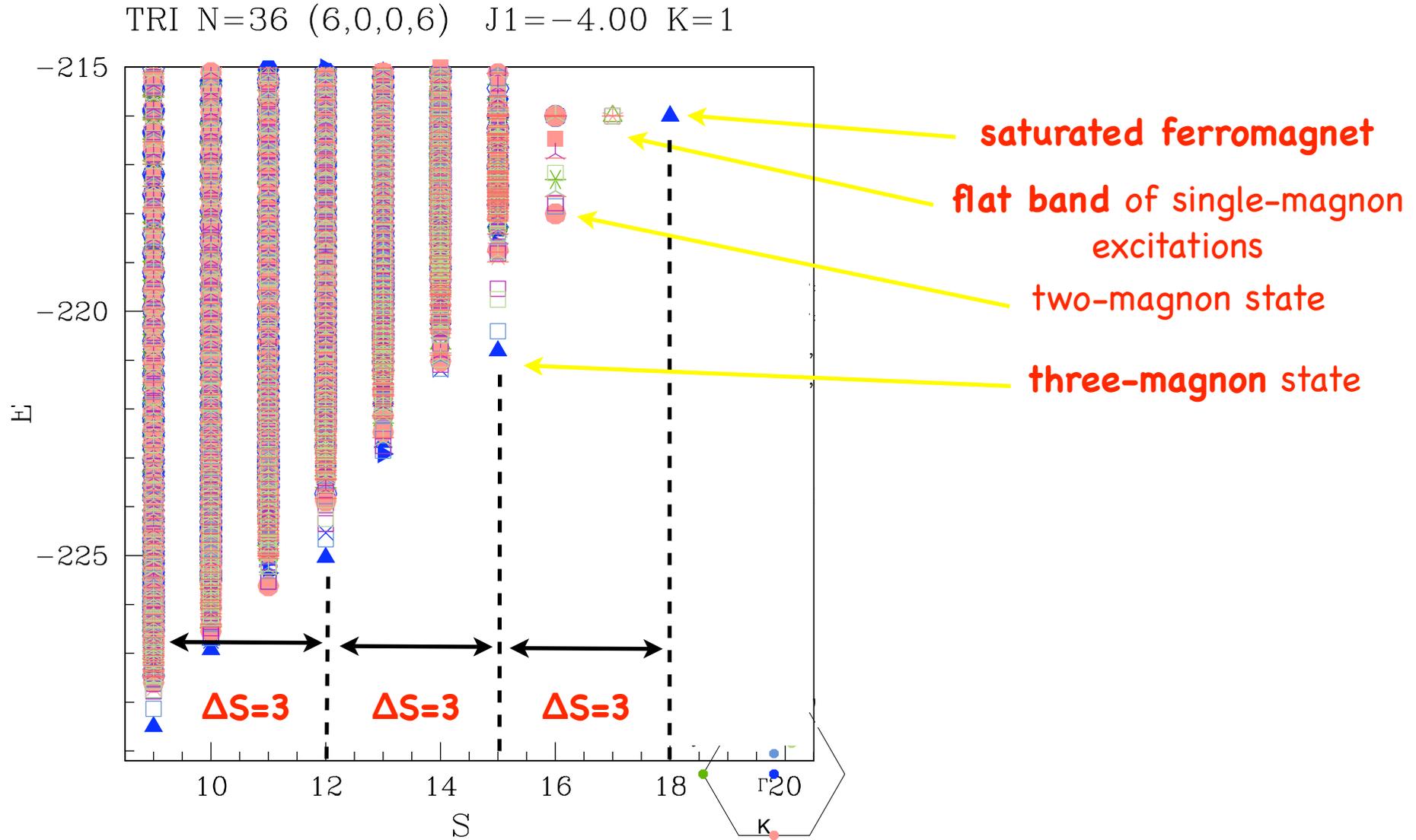
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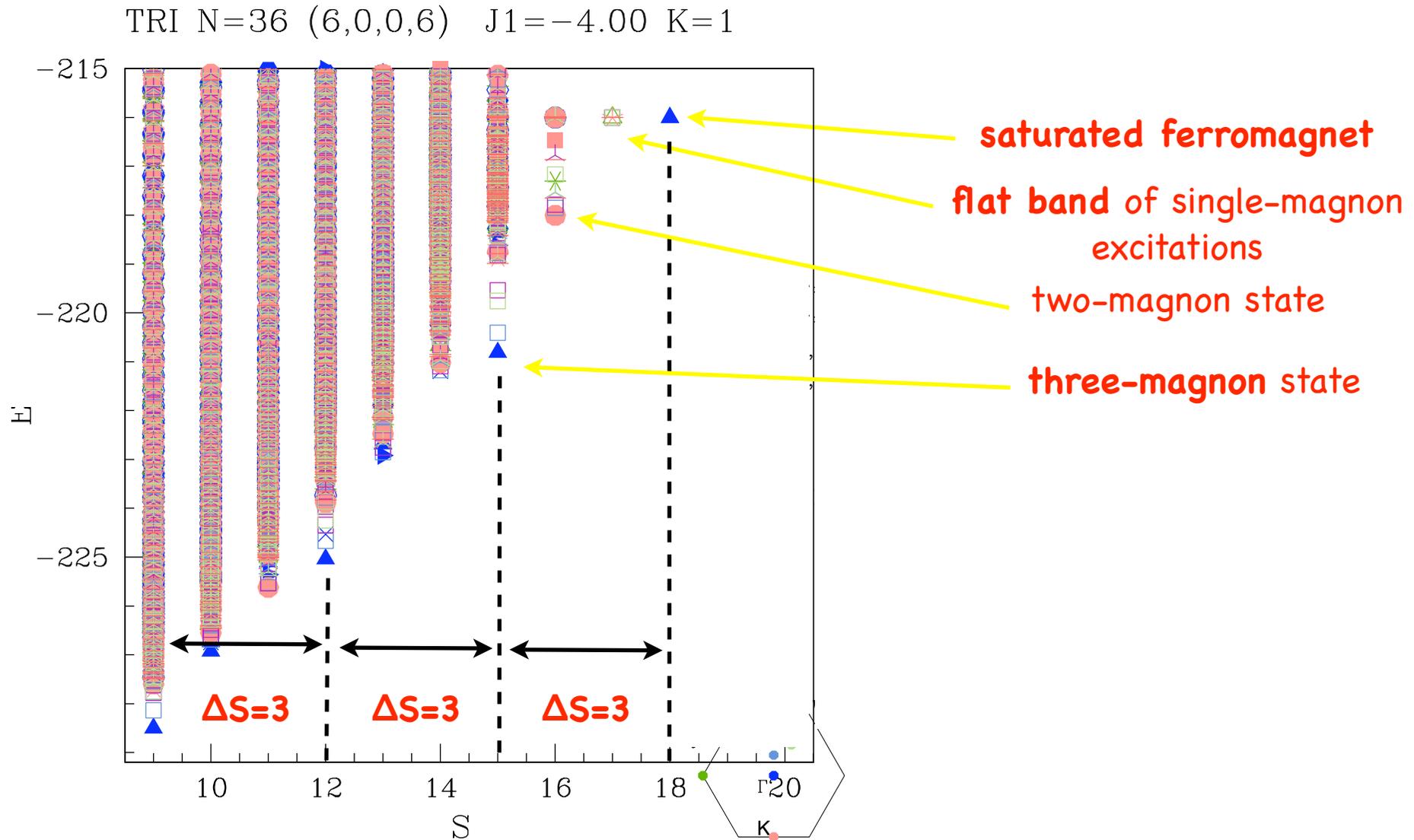
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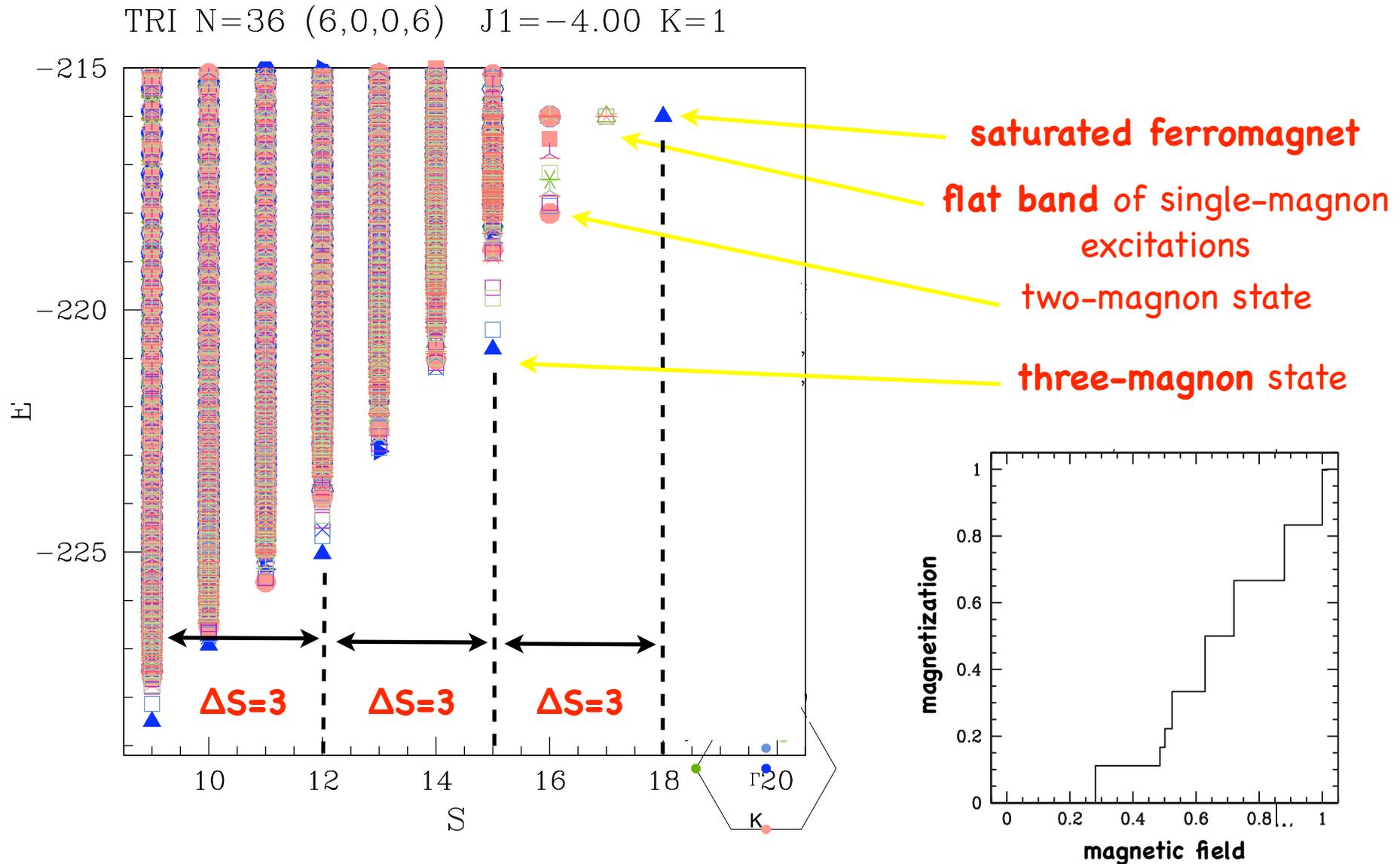
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- structure is period-3 not period-2 -

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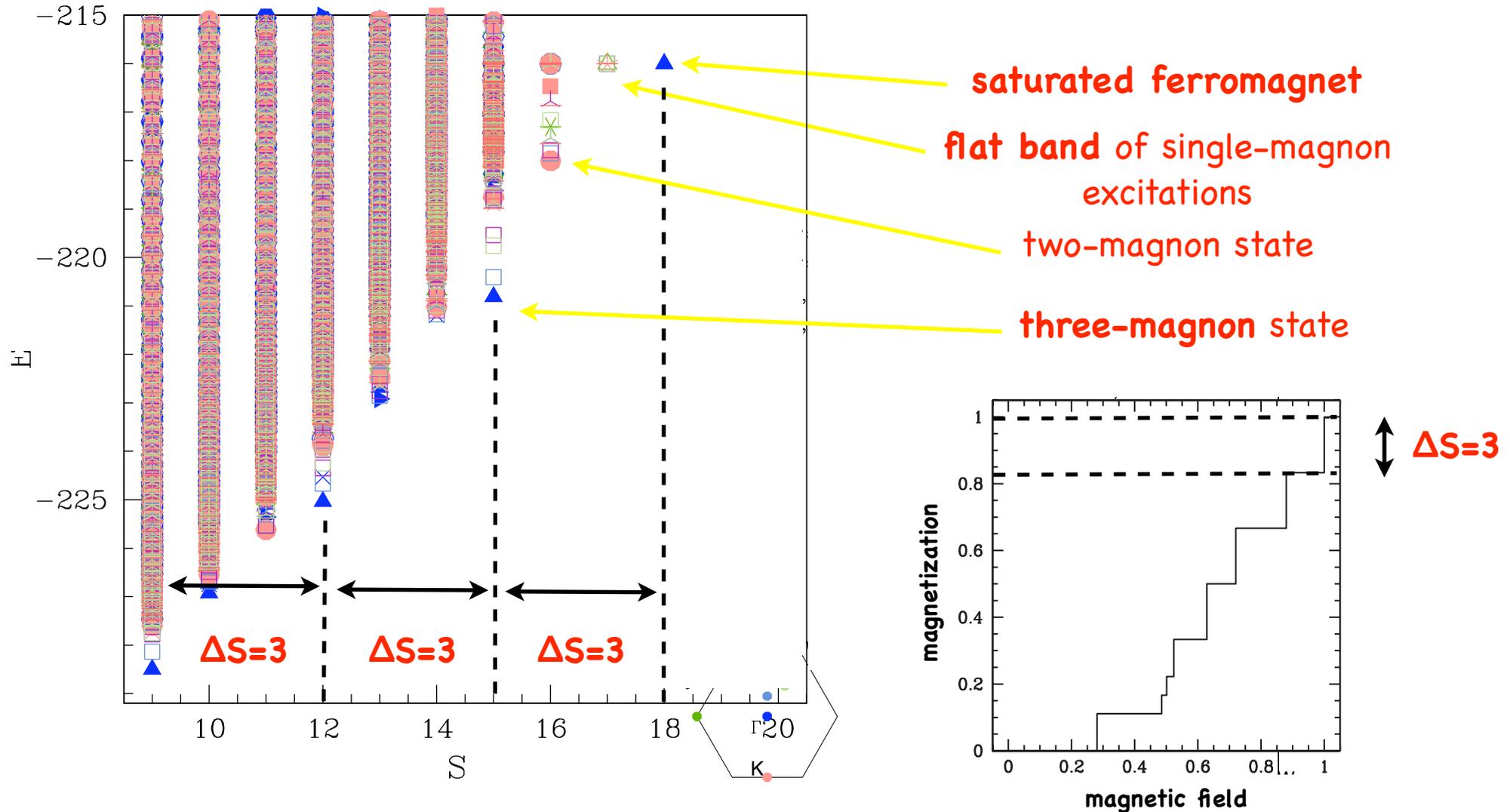


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TRI N=36 (6,0,0,6) J1=-4.00 K=1



- structure is period-3 not period-2 -

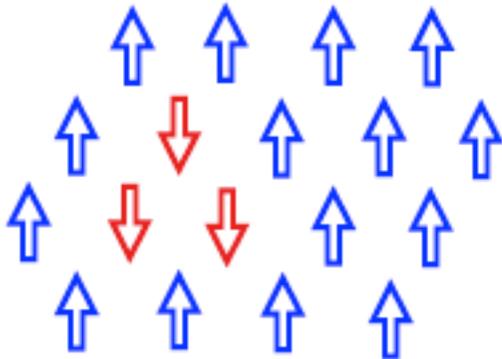
another new quantum phase !

- three-spin bound states at high magnetic field -

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"wiggle-on"

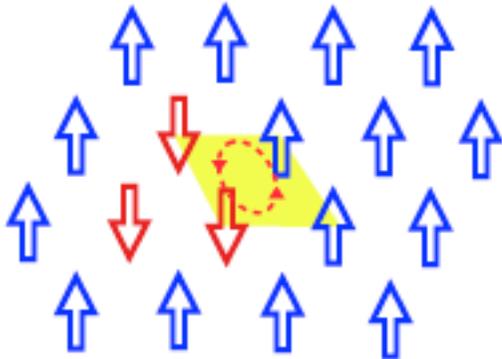


three spins propagate
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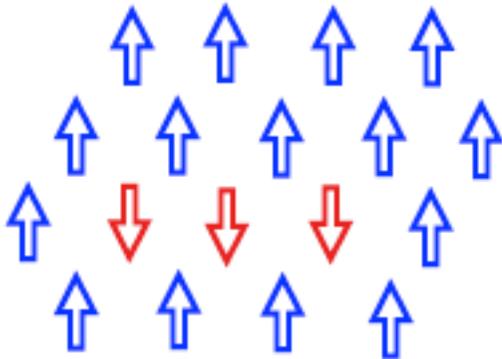


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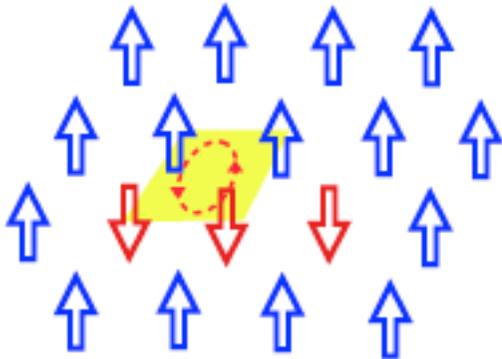


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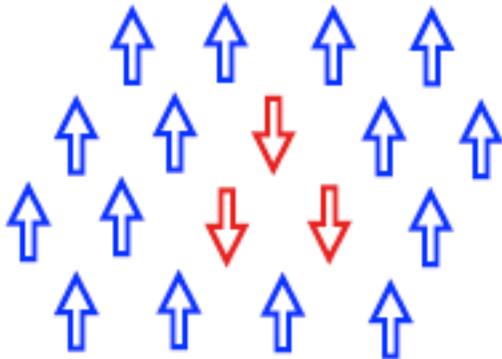


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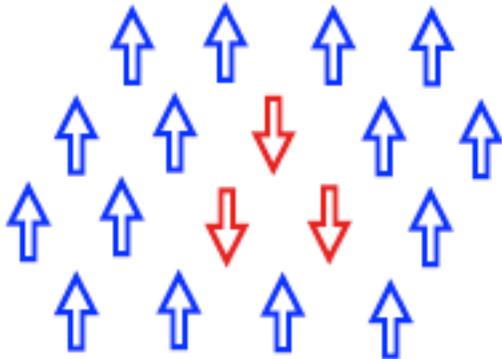


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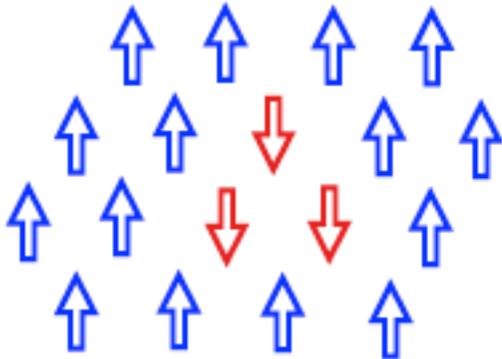


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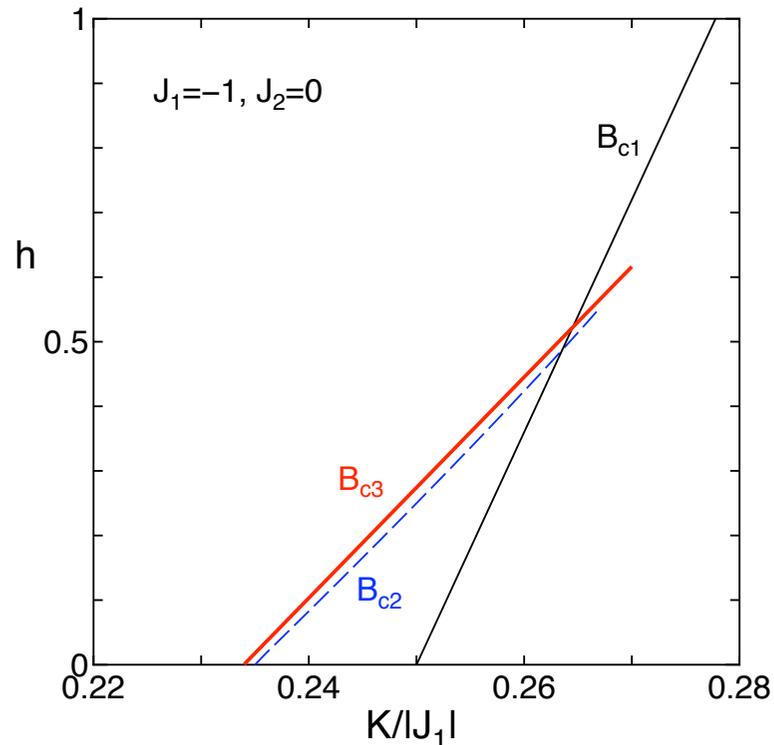
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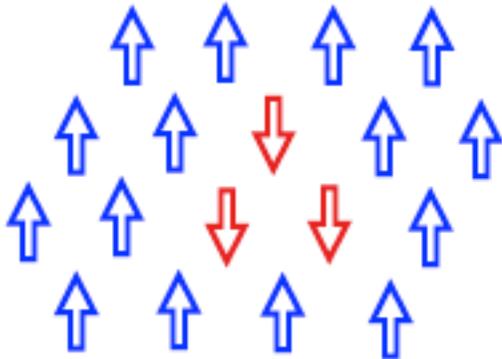
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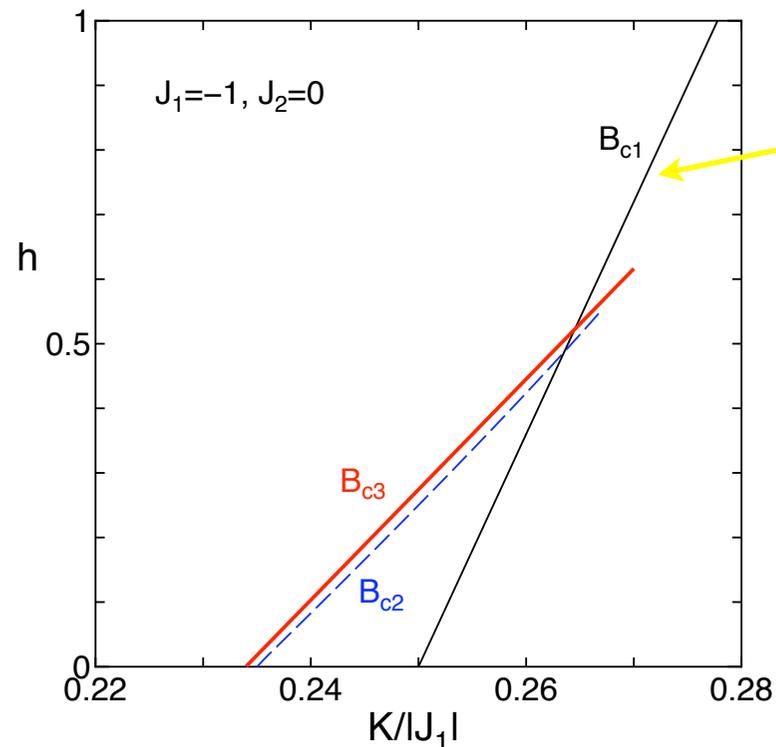
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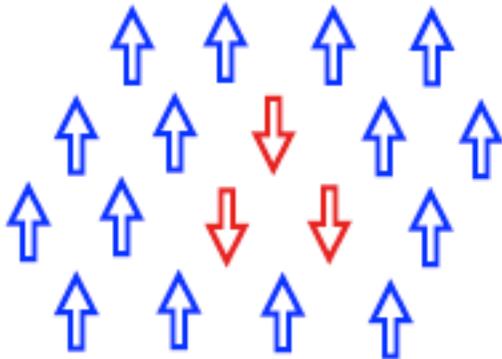


1-magnon
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i.e. **canted Néel**
state (exact)

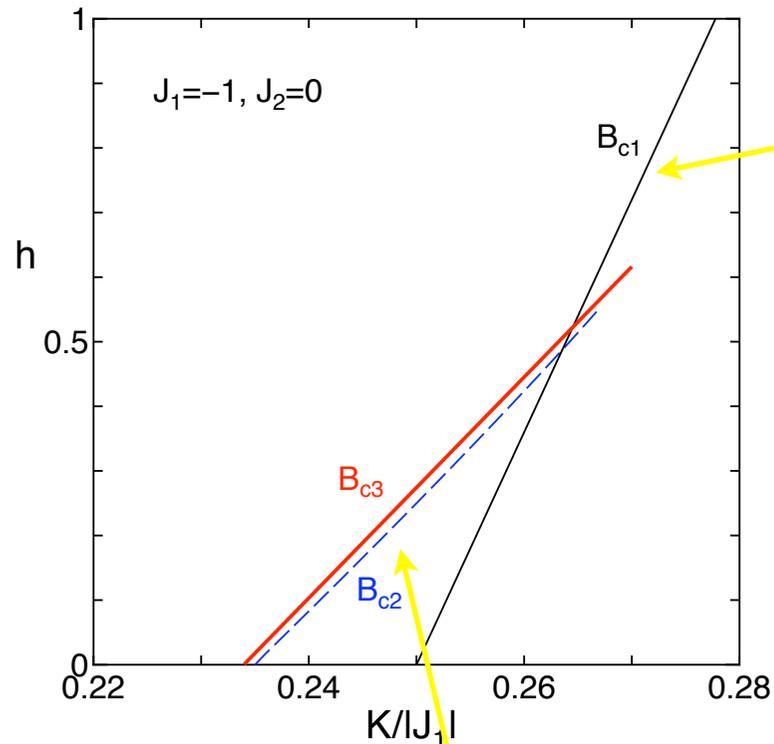
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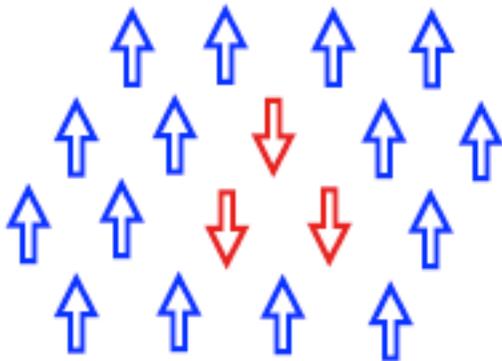
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2-magnon instability
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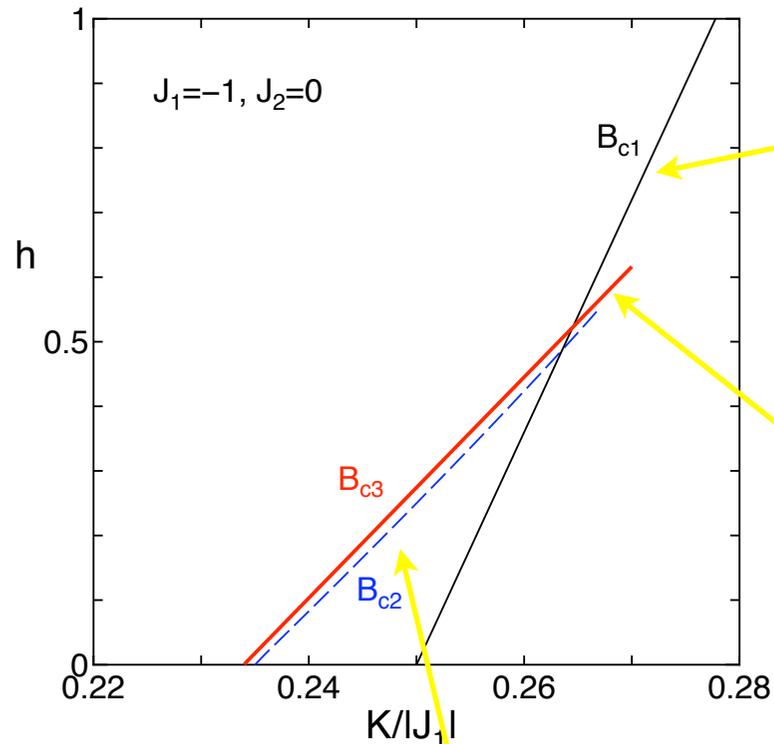
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"wobble-on"



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1-magnon instability
i.e. canted Néel state (exact)

3-magnon instability
i.e. "triatric" state
(variational lower bound)

2-magnon instability
i.e. nematic state (exact)

symmetries of the triatic phase

- three-spin bound states at high magnetic field -

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- three-spin bound states at high magnetic field -

order parameter is
rank 3 tensor :

$$\mathcal{R}e \{ S_i^- S_j^- S_k^- \} = 2S_i^x S_j^x S_k^x - 2S_i^x S_j^y S_k^y - 2S_i^y S_j^x S_k^y - 2S_i^y S_j^y S_k^x$$

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matrix elements of order parameter
are linear combinations of fully
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triangular plaquette

symmetries of the triatic phase

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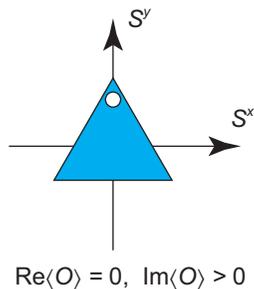
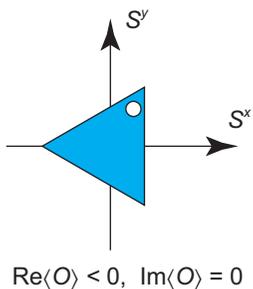
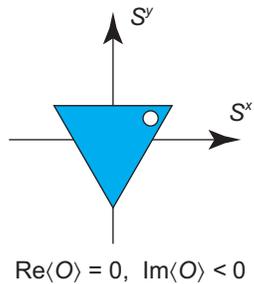
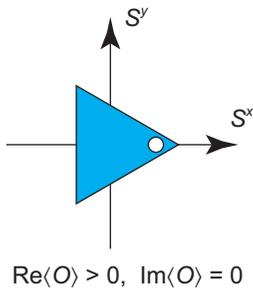
order parameter is
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$$\text{Im} \{ S_i^- S_j^- S_k^- \} = 2S_i^x S_j^x S_k^y + 2S_i^x S_j^y S_k^x + 2S_i^y S_j^x S_k^x - 2S_i^y S_j^y S_k^y$$

matrix elements of order parameter
are linear combinations of fully
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in applied magnetic field, order
parameter is planar and maps onto
itself under rotations through $2\pi/3$



symmetries of the triatic phase

- three-spin bound states at high magnetic field -

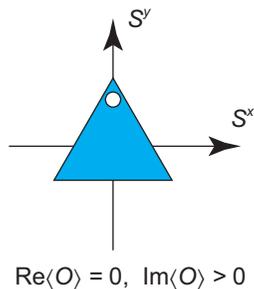
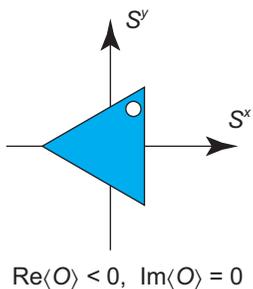
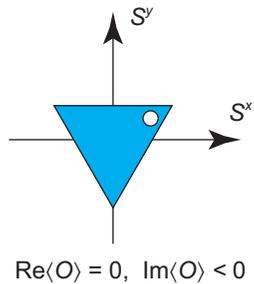
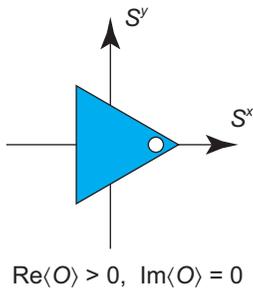
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rank 3 tensor :

$$\text{Re} \{ S_i^- S_j^- S_k^- \} = 2S_i^x S_j^x S_k^x - 2S_i^x S_j^y S_k^y - 2S_i^y S_j^x S_k^y - 2S_i^y S_j^y S_k^x$$

$$\text{Im} \{ S_i^- S_j^- S_k^- \} = 2S_i^x S_j^x S_k^y + 2S_i^x S_j^y S_k^x + 2S_i^y S_j^x S_k^x - 2S_i^y S_j^y S_k^y$$

matrix elements of order parameter
are linear combinations of fully
symmetrized spin operators on
triangular plaquette

in applied magnetic field, order
parameter is planar and maps onto
itself under rotations through $2\pi/3$



FM octopolar order
naively has k^2
dispersion

$$\Rightarrow cV \propto T \text{ in 2D}$$

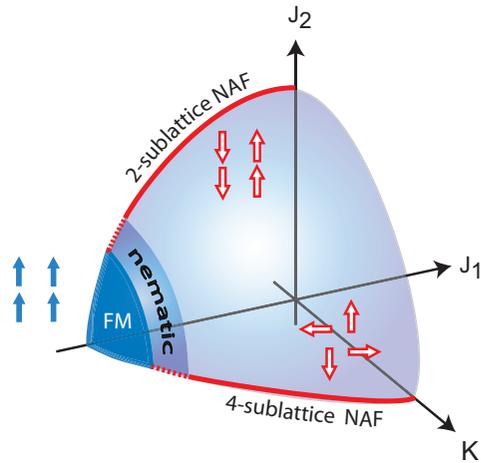
c.f HeIII on graphite

conclusions

- new quantum phases for all the family -

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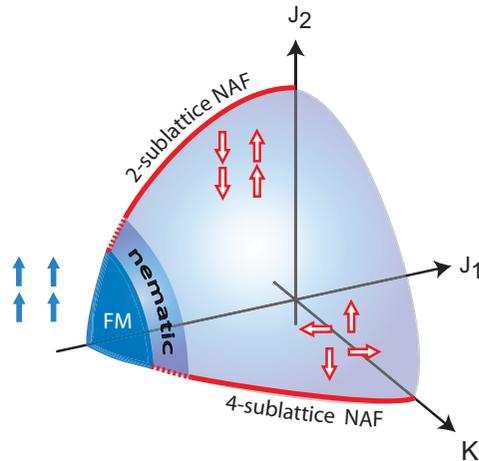
new **nematic** phase in
square lattice
frustrated ferromagnets
(c.f. quasi-2D vanadates)

Shannon, Momoi + Sindzingre,
PRL 2006

Shannon et al., EPJB 2004

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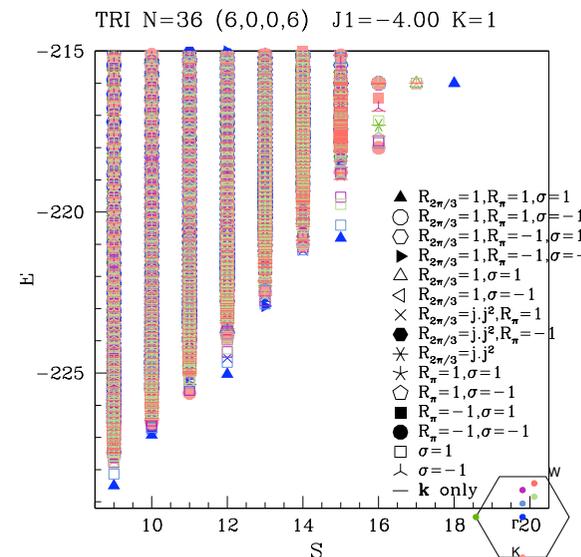
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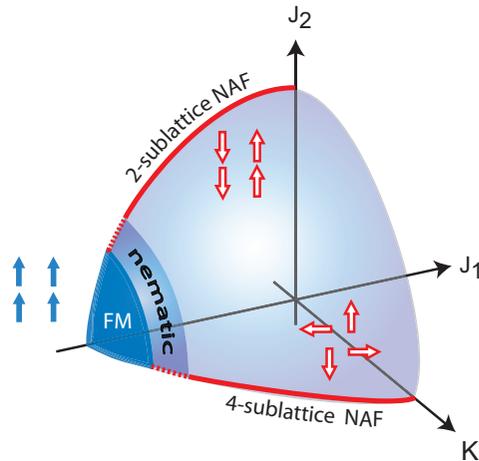
Momoi, Sindzingre + Shannon,
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Momoi + Shannon, PTP 2005



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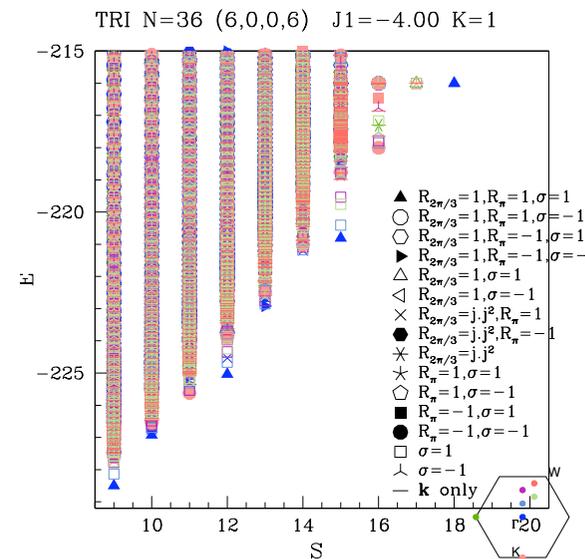
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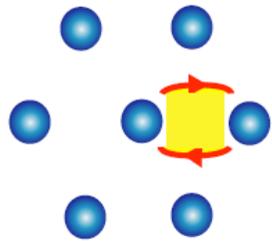


frustrated ferromagnets are fun and there's lots still to learn

that's all folks...

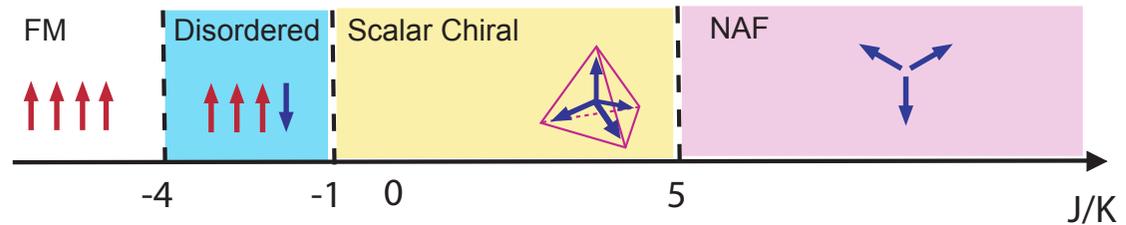
multiple spin exchange on the triangular lattice

$$\mathcal{H} = J \sum_{\langle ij \rangle} P_{ij} + K \sum_{\langle ijkl \rangle} [P_{ijkl} + P_{ijkl}]$$



Classical limit $S = \infty$

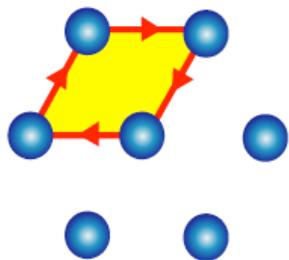
$K > 0$



...Momoi et al.

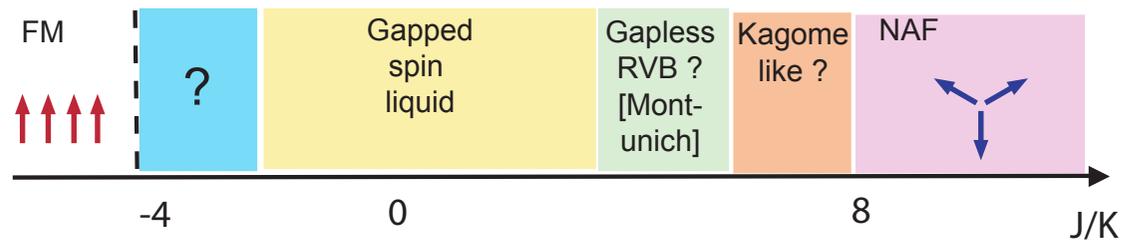
$$P_{ij} = \left[\frac{1 + \vec{\sigma}_i \cdot \vec{\sigma}_j}{2} \right]$$

$$= 2\vec{S}_i \cdot \vec{S}_j + \dots$$



Quantum limit $S = 1/2$

$K > 0$



...Lhullier et al.

He III

$$P_{ijkl} + P_{ijkl}^{-1} = \vec{S}_i \cdot \vec{S}_j + \dots + 4 \left(\vec{S}_i \cdot \vec{S}_j \right) \left(\vec{S}_k \cdot \vec{S}_l \right) + \dots$$